

高等教育（研究生）国家级教学成果奖申请书附件

教学成果应用及效果证明材料

成果名称： 创新引领+内生协同+国际赋能：物理学科基础拔尖
人才培养模式构建与实践

成果完成人： 田浩、张宇、赵纪军、马晶、任延宇、孙秀冬、钱
冬、袁承勋、高峻峰、李立、李淑凤、李德友、谭
鹏、蒋雪、王宇

成果完成单位： 哈尔滨工业大学、大连理工大学、上海交通大学

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(1) 教学成果获奖情况

| 序号 | 成果名称/获奖名称 | 级别 | 获奖等级 | 教师姓名 | 年份 |
|----|------------------|-----|------|------|------|
| 1. | 黑龙江省高等教育教学成果奖 | 省级 | 一等奖 | 田浩 | 2022 |
| 2. | 辽宁省教学成果奖 | 省级 | 一等奖 | 赵纪军 | 2018 |
| 3. | 辽宁省教学成果奖 | 省级 | 一等奖 | 李淑凤 | 2018 |
| 4. | 黑龙江省高等教育教学成果奖 | 省级 | 二等奖 | 张宇 | 2022 |
| 5. | 辽宁省教学成果奖 | 省级 | 二等奖 | 李淑凤 | 2022 |
| 6. | 辽宁省教学成果奖 | 省级 | 二等奖 | 赵纪军 | 2021 |
| 7. | 工业和信息化部研究型教学创新团队 | 省部级 | | 张宇 | 2017 |
| 8. | 哈尔滨工业大学研究生教育成果奖 | 校级 | 一等奖 | 田浩 | 2022 |
| 9. | 哈尔滨工业大学教学成果奖 | 校级 | 一等奖 | 张宇 | 2022 |

(1) 黑龙江省高等教育教学成果奖一等奖获奖证明



(2) 辽宁省教学成果奖一等奖获奖证明



(3) 辽宁省教学成果奖一等奖获奖证明



(4) 黑龙江省高等教育教学成果奖二等奖获奖证明



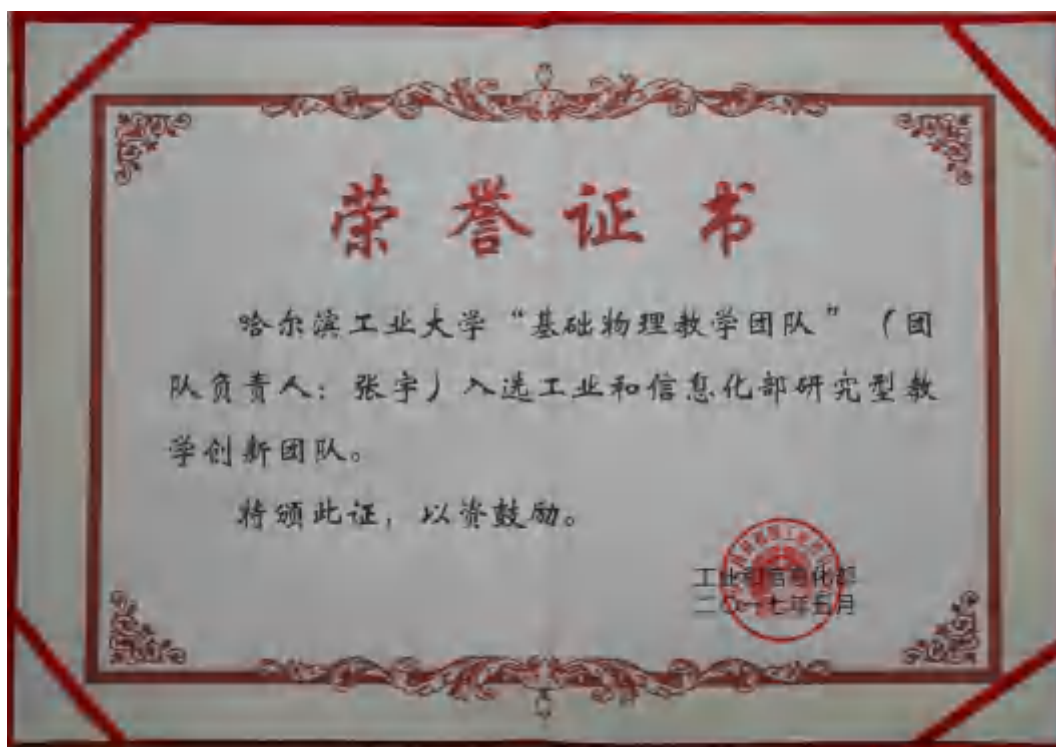
(5) 辽宁省教学成果奖二等奖获奖证明



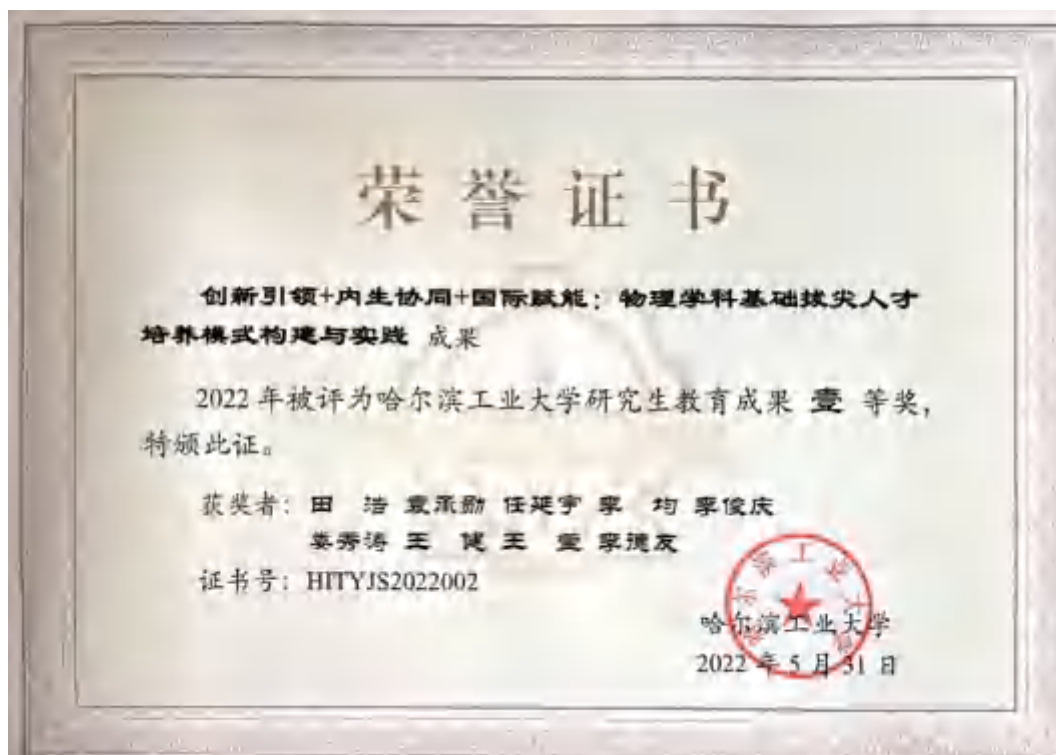
(6) 辽宁省教学成果奖二等奖获奖证明



(7) 工业和信息化部研究型教学创新团队



(8) 哈尔滨工业大学研究生教育成果奖




(9) 哈尔滨工业大学教学成果奖



(2) 自制科教仪器推广应用佐证材料

| 推广类别 | 单位名称（总套数） |
|----------------|---|
| 国外高校 (1所) | 澳大利亚卧龙岗大学(1) |
| 985高校 (17所) | 清华大学(9)、中国科技大学(1)、南京大学(1)、西安交通大学(4)、四川大学(8)、华中科技大学(1)、哈尔滨工业大学(8)、天津大学(1)、中南大学(1)、中山大学(22)、电子科技大学(2)、兰州大学(11)、吉林大学(14)、厦门大学(5)、大连理工大学(9)、重庆大学(3)、中国海洋大学(6) |
| 211高校 (9所) | 南京航空航天大学(1)、南京理工大学(3)、西南大学(3)、陕西师范大学(5)、武汉理工大学(1)、东华大学(1)、东华理工大学(1)、西南交通大学(10)、北京工业大学(4) |
| 其他高校 (23所) | 东北大学秦皇岛分校(2)、陕西科技大学(2)、西安工业大学(3)、江苏大学(2)、河南大学(2)、河南师范大学(2)、合肥学院(2)、西安建筑科技大学(2)、桂林理工大学(4)、桂林电子科技大学(11)、西南民族学院(2)、沈阳工业大学(1)、盐城工学院(2)、重庆科技学院(2)、信阳师范学院(1)、太原科技大学(2)、成都大学(1)、齐鲁工业大学(4)、西安外事学院(5)、河北大学(1)、安阳工学院(1)、昆明理工大学(16)、南京工业大学(14) |
| 研究院所 (10所) | 中科院成都光电所(2)、中科院物理研究所(2)、山东科学院新材料研究所(1)、中电46所(1)、佛山新材料研究所(4)、中国船舶重工集团公司第715所(3)、中国散裂中子源研究中心(1)、北京中芯云科仪器有限公司(1)、乌镇实验室(1)、深圳先进电子材料国际创新研究院(2) |


应用证明

| | |
|------------|---|
| 项目名称: | 高性能低温燃料电池(型号:ZDM-200-100,1套) |
| 使用单位: | 北京航空航天大学材料科学与工程学院 |
| 通讯地址: | 北京市海淀区学院路37号 |
| 应用起止时间: | 2015年至今 |
| 应用情况及社会效益: | <p>由哈尔滨工业大学物理学院教授研究开发的低温质子交换膜燃料电池(型号:ZDM-200-100,1套)具有体积小、功率密度高、启动时间短等优点,在航空航天领域的飞行器和高空飞行器作为能源等应用领域得到了应用。该设备应用于飞行器电源、小卫星、微纳卫星、深空探测器和特种探测器的性能表现:解决了传统低温质子交换膜燃料电池启动慢、功率密度低、低温启动困难等问题,实现了高功率密度、体积小、启动快的性能指标,为特种飞行器电源研究、应用了有效的社会效益。</p> <p>同时,该设备还用于“我国首个极地飞行试验器”极区高空探测飞行器的研制和测试,解决了极地高空探测、航空比翼型飞行器飞行控制等问题,受到中科院60余人、航空工业6人、解放军10人、俄罗斯经济代表团、俄罗斯商、政府官员、宇航员、航天员等的肯定,应用效果良好。</p> |
| 联系人: |  北京航空航天大学材料学院 2015年12月15日 |

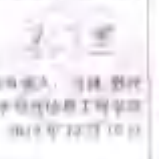
应用证明

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|------------|---|
| 项目名称: | 低温质子交换膜燃料电池(型号:ZDM-200-100,1套) 高温质子交换膜燃料电池(型号:ZDM-200-1000,1套) 低温质子交换膜燃料电池(型号:ZDM-200-100,1套) |
| 使用单位: | 清华大学新能源与燃料电池工艺研究中心实验室 |
| 通讯地址: | 北京市海淀区清华大学材料学院楼 8513 室 |
| 应用起止时间: | 2015 年至今 |
| 应用情况及社会效益: | <p>清华大学能源与动力工程系开发的低温质子交换膜燃料电池在航天领域、航空领域和军事领域得到了广泛的应用,解决了传统低温质子交换膜燃料电池启动慢、功率密度低、低温启动困难等问题,实现了高功率密度、体积小、启动快的性能指标,为特种飞行器电源研究、应用了有效的社会效益。</p> <p>同时,该设备还用于“我国首个极地飞行试验器”极区高空探测飞行器的研制和测试,解决了极地高空探测、航空比翼型飞行器飞行控制等问题,受到中科院60余人、航空工业6人、解放军10人、俄罗斯经济代表团、俄罗斯商、政府官员、宇航员、航天员等的肯定,应用效果良好。</p> |
| 联系人: |  清华大学能源与动力工程系 2015年12月15日 |

应用证明

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|------------|--|
| 项目名称: | 高性能低温燃料电池(型号:ZDM-200-100,1套) 高温质子交换膜燃料电池(型号:ZDM-200-1000,1套) |
| 使用单位: | 电子科技大学材料与工程学院 |
| 通讯地址: | 四川省成都市高新区二环路中段11号 |
| 应用起止时间: | 2015 年至今 |
| 应用情况及社会效益: | <p>由哈尔滨工业大学物理学院教授研究开发的低温质子交换膜燃料电池(型号:ZDM-200-100,1套)、高温质子交换膜燃料电池(型号:ZDM-200-1000,1套)在航天领域、航空领域和军事领域得到了广泛的应用,解决了传统低温质子交换膜燃料电池启动慢、功率密度低、低温启动困难等问题,实现了高功率密度、体积小、启动快的性能指标,为特种飞行器电源研究、应用了有效的社会效益。</p> <p>同时,该设备还用于“我国首个极地飞行试验器”极区高空探测飞行器的研制和测试,解决了极地高空探测、航空比翼型飞行器飞行控制等问题,受到中科院60余人、航空工业6人、解放军10人、俄罗斯经济代表团、俄罗斯商、政府官员、宇航员、航天员等的肯定,应用效果良好。</p> |
| 联系人: |  电子科技大学材料与工程学院 2015年12月15日 |

应用证明

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|------------|---|
| 项目名称: | 高性能低温燃料电池(型号:ZDM-200-100,1套) 高温质子交换膜燃料电池(型号:ZDM-200-1000,1套) |
| 使用单位: | 清华大学能源与动力工程系实验室 |
| 通讯地址: | 北京市海淀区清华大学材料学院楼 8513 室 |
| 应用起止时间: | 2015 年至今 |
| 应用情况及社会效益: | <p>清华大学能源与动力工程系开发的低温质子交换膜燃料电池在航天领域、航空领域和军事领域得到了广泛的应用,解决了传统低温质子交换膜燃料电池启动慢、功率密度低、低温启动困难等问题,实现了高功率密度、体积小、启动快的性能指标,为特种飞行器电源研究、应用了有效的社会效益。</p> <p>同时,该设备还用于“我国首个极地飞行试验器”极区高空探测飞行器的研制和测试,解决了极地高空探测、航空比翼型飞行器飞行控制等问题,受到中科院60余人、航空工业6人、解放军10人、俄罗斯经济代表团、俄罗斯商、政府官员、宇航员、航天员等的肯定,应用效果良好。</p> |
| 联系人: |  清华大学能源与动力工程系 2015年12月15日 |

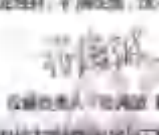
应用证明

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| 项目名称: | 高温介电温湿测量仪(型号 TDM-RT-1000, 1套) |
| 应用单位: | 西安交通大学材料科学与工程学院 |
| 通讯地址: | 陕西省西安市雁塔区雁塔路9号 |
| 应用起止时间: | 2019年至今 |
| 应用情况及社会效益: | <p>西安交通大学材料科学与工程学院开发的“高温介电温湿测量仪”(型号 TDM-RT-1000, 1套)在我单位使用,并在南京信息工程大学完成了项目,西安交通大学材料科学与工程学院支持该设备对我单位的课题中起到了应用,该设备应用于开展新型材料的高温介电温湿性能研究,解决了传统测试方法的问题,提高了测试材料的效率和准确性,带来了可观的社会效益。</p> <p>该系列测试精度高,效率高,稳定性好,扩展性强,技术化的明显,应用效果好。</p> |
| 特此证明! | |


 项目负责人: 陶勇
 西安交通大学材料科学与工程学院
 2019年12月18日

应用证明

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|------------|---|
| 项目名称: | 高温介电温湿测量仪(型号 TDM-RT-1000, 1套) |
| 应用单位: | 华中科技大学材料科学与工程学院 |
| 通讯地址: | 武汉市洪山区珞喻路1037号 |
| 应用起止时间: | 2015年至今 |
| 应用情况及社会效益: | <p>西安交通大学材料科学与工程学院开发的“高温介电温湿测量仪”(型号 TDM-RT-1000, 1套)在我单位使用,并在湖北为“量子阱结构的导电材料”项目提供高温介电温湿测试材料研究用(1732147)和“高温介电温湿测量仪”项目提供材料研究“高性能材料与电路集成” (11672042)项目提供材料研究,该设备在我单位得到了应用,用于开展大电流密度下的高温介电材料研究,解决了高温介电材料测试的问题,带来了可观的社会效益。</p> <p>该系列测试精度高,效率高,稳定性好,扩展性强,技术化的明显,应用效果好。</p> |
| 特此证明! | |


 项目负责人: 任其波
 华中科技大学材料科学与工程学院
 2019年12月18日


应用证明

| | |
|------------|--|
| 项目名称: | 高温介电温湿测量仪(型号 TDM-RT-1000, 1套) |
| 应用单位: | 南京工业大学现代工程与材料学院 |
| 通讯地址: | 南京市鼓楼区江浦路21号 |
| 应用起止时间: | 2016年至今 |
| 应用情况及社会效益: | <p>西安交通大学材料科学与工程学院开发的“高温介电温湿测量仪”(型号 TDM-RT-1000, 1套)在我单位使用,并在南京信息工程大学完成了项目,该设备应用于开展新型材料的高温介电温湿性能研究,解决了传统测试方法的问题,提高了测试材料的效率和准确性,带来了可观的社会效益。</p> <p>该系列测试精度高,效率高,稳定性好,扩展性强,技术化的明显,应用效果好。</p> |
| 特此证明! | |


 项目负责人: 陶勇
 南京工业大学现代工程与材料学院
 2019年12月18日

应用证明

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|------------|---|
| 项目名称: | 高温介电温湿测量仪(型号 TDM-RT-1000, 1套) |
| 应用单位: | 桐乡经纬科技股份有限公司 |
| 通讯地址: | 浙江省嘉兴市南湖新区南湖205号2楼24室 |
| 应用起止时间: | 2018年至今 |
| 应用情况及社会效益: | <p>西安交通大学材料科学与工程学院开发的“高温介电温湿测量仪”(型号 TDM-RT-1000, 1套)在我单位使用,并在浙江经纬科技股份有限公司,用于开展新型材料的高温介电温湿性能研究,解决了传统测试方法的问题,提高了测试材料的效率和准确性,带来了可观的社会效益。</p> <p>该系列测试精度高,效率高,稳定性好,扩展性强,技术化的明显,应用效果好。</p> |
| 特此证明! | |


 负责人: 陶勇
 桐乡经纬科技股份有限公司
 2019年12月18日

应用证明


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|---------|-------------------------------|
| 项目名称: | 高温介电陶瓷量仪(型号:Z208-07-1000, I型) |
| 应用单位: | 西安交通大学电子陶瓷与器件教育部重点实验室 |
| 通讯地址: | 陕西省西安市雁塔区西安交通大学曲江校区西六教 |
| 应用起止时间: | 2014-到现在 |

应用情况及社会效益:

西安理工大学物理学院冯教授研发的高温介电陶瓷量仪(型号:Z208-07-1000, I型)在投产后使用,并在“国家高技术研究发展计划(863计划)陶瓷与介电材料”项目支持下进行了实验并得到了应用。该项目解决了传统陶瓷电容元件在高温下性能下降的问题,提高了产品的可靠性和使用寿命,为国内陶瓷电容元件的研究、生产和应用提供了重要的技术支持。

该产品性能稳定、精度高、使用寿命长、技术先进、应用效果良好。

特此证明!


 西安交通大学电子陶瓷与器件教育部重点实验室
 2014年11月10日

应用证明


| | |
|---------|---|
| 项目名称: | 高温介电陶瓷量仪(型号:Z208-07-1000, I型) 高温介电陶瓷量仪(型号:Z208-07-1000, II型) |
| 应用单位: | 西安理工大学材料学院 |
| 通讯地址: | 陕西省西安市雁塔区大雁塔路2号 |
| 应用起止时间: | 2014年至今 |

应用情况及社会效益:

西安理工大学物理学院冯教授研发的高温介电陶瓷量仪(型号:Z208-07-1000, I型)、II型)在投产后使用,并在“国家高技术研究发展计划(863计划)陶瓷与介电材料”项目支持下进行了实验并得到了应用。该项目解决了传统陶瓷电容元件在高温下性能下降的问题,提高了产品的可靠性和使用寿命,为国内陶瓷电容元件的研究、生产和应用提供了重要的技术支持。

该产品性能稳定、精度高、使用寿命长、技术先进、应用效果良好。

特此证明!


 西安理工大学材料学院
 2014年12月10日

应用证明


| | |
|---------|------------------------------|
| 项目名称: | 高温介电陶瓷量仪(型号:Z208-07-800, I型) |
| 应用单位: | 浙江嘉德电子股份有限公司 |
| 通讯地址: | 浙江省嘉兴市南湖机场1188号 |
| 应用起止时间: | 2017年至今 |

应用情况及社会效益:

西安理工大学物理学院冯教授研发的高温介电陶瓷量仪(型号:Z208-07-800, I型)在我公司使用,并在高温陶瓷材料、氧化铝陶瓷材料等研究领域进行了实验和产业化生产得到了应用。该项目解决了传统陶瓷电容元件在高温下性能下降的问题,提高了产品的可靠性和使用寿命,为国内陶瓷电容元件的研究、生产和应用提供了重要的技术支持。

该产品性能稳定、精度高、使用寿命长、技术先进、应用效果良好。

特此证明!


 浙江嘉德电子股份有限公司
 2018年11月10日

应用证明


| | |
|---------|-------------------------------|
| 项目名称: | 高温介电陶瓷量仪(型号:Z208-07-1000, I型) |
| 应用单位: | 中南大学粉末冶金研究院 |
| 通讯地址: | 湖南省长沙市岳麓区麓山南路933号 |
| 应用起止时间: | 2019年12月至今 |

应用情况及社会效益:

西安理工大学物理学院冯教授研发的高温介电陶瓷量仪(型号:Z208-07-1000, I型)在投产后使用,并在国家自然科学面上项目及科技项目的资助下得到了应用。该项目解决了传统陶瓷电容元件在高温下性能下降的问题,提高了产品的可靠性和使用寿命,为国内陶瓷电容元件的研究、生产和应用提供了重要的技术支持。

该产品性能稳定、精度高、使用寿命长、技术先进、应用效果良好。


特此证明!


 中南大学粉末冶金研究院
 2019年12月10日


应用证明

| | |
|------------|--|
| 项目名称: | 高压介电谱测量仪(型号 TZDM-RT-1000, 1套) |
| 应用单位: | 西安工业大学机电工程学院 |
| 通讯地址: | 西安市未央区纺西街2号 |
| 应用起止时间: | 2018年至今 |
| 应用情况及社会效益: | <p>西安工业大学机电工程学院副教授张晋开发的超高压介电谱测量仪(型号 TZDM-RT-1000, 1套)在各单位使用, 在应用过程中, 为提升材料性能的实验条件起到了作用, 为改善仪器于常规测量设备相比的稳定性及便捷性研究, 提升了粉末涂层材料电学性能测试的测试精度, 促进了光电互容材料的研究, 带来了可观的社会效益。</p> <p>同时, 该仪器开发了“介电材料温度湿度测量”, 解决了粉末涂层材料、气态薄膜材料的测试问题, 受到本单位师生、研究人员好评。</p> <p>该仪器测试精度高, 效率高, 稳定性好, 防护性强, 技术先进可靠, 应用效果好。</p> |
| 特此证明! | |
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应用证明

| | |
|------------|---|
| 项目名称: | 高压介电谱测量仪(型号 TZDM-RT-1000, 1套) |
| 应用单位: | 西安理工大学材料科学与工程学院 |
| 通讯地址: | 陕西省西安市雁塔区雁塔西路28号 |
| 应用起止时间: | 2017年至今 |
| 应用情况及社会效益: | <p>西安理工大学材料科学与工程学院副教授张晋开发的超高压介电谱测量仪(型号 TZDM-RT-1000, 1套)在各单位使用, 在应用过程中, 为提升材料性能的实验条件起到了作用, 为改善仪器于常规测量设备相比的稳定性及便捷性研究, 提升了粉末涂层材料电学性能测试的测试精度, 促进了光电互容材料的研究, 带来了可观的社会效益。</p> <p>同时, 该仪器开发了“介电材料温度湿度测量”, 解决了粉末涂层材料、气态薄膜材料的测试问题, 受到本单位师生、研究人员好评。</p> <p>该仪器测试精度高, 效率高, 稳定性好, 防护性强, 技术先进可靠, 应用效果好。</p> |
| 特此证明! | |
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应用证明

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|------------|--|
| 项目名称: | 高压介电谱测量仪(型号 TZDM-RT-1000, 1套) |
| 应用单位: | 西安大学材料与电子学院 |
| 通讯地址: | 西安市金明大道1号西安大学新校区15栋东 |
| 应用起止时间: | 2019年7月-至今 |
| 应用情况及社会效益: | <p>西安理工大学教授张晋副教授开发的超高压介电谱测量仪(型号 TZDM-RT-1000, 1套)在各单位使用, 在国家自然基金面上项目“基于高性能太赫兹电介质的多功能复合电介质材料及其制备工艺研究”科研项目支持下, 为提升材料性能的实验条件起到了作用, 为改善仪器于常规测量设备相比的稳定性及便捷性研究, 提升了粉末涂层材料电学性能测试的测试精度, 促进了光电互容材料的研究, 带来了可观的社会效益。</p> <p>同时, 该仪器开发了“介电材料温度湿度测量”, 解决了粉末涂层材料、气态薄膜材料的测试问题, 受到本单位师生、研究人员好评。</p> <p>该仪器测试精度高, 效率高, 稳定性好, 防护性强, 技术先进可靠, 应用效果好。</p> |
| 特此证明! | |
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应用证明

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|------------|---|
| 项目名称: | 高压介电谱测量仪(型号 TZDM-RT-1000, 1套) |
| 应用单位: | 西安理工大学材料科学与工程学院 |
| 通讯地址: | 陕西省西安市雁塔区雁塔西路28号 |
| 应用起止时间: | 2018年至今 |
| 应用情况及社会效益: | <p>西安理工大学材料科学与工程学院副教授张晋开发的超高压介电谱测量仪(型号 TZDM-RT-1000, 1套)在各单位使用, 在应用过程中, 为提升材料性能的实验条件起到了作用, 为改善仪器于常规测量设备相比的稳定性及便捷性研究, 提升了粉末涂层材料电学性能测试的测试精度, 促进了光电互容材料的研究, 带来了可观的社会效益。</p> <p>同时, 该仪器开发了“介电材料温度湿度测量”, 解决了粉末涂层材料、气态薄膜材料的测试问题, 受到本单位师生、研究人员好评。</p> <p>该仪器测试精度高, 效率高, 稳定性好, 防护性强, 技术先进可靠, 应用效果好。</p> |
| 特此证明! | |
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应用证明

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| 项目名称: | 中国中车磁浮列车(型号:ZDM-01-1000,2套); 南京中车磁浮列车(型号:ZDM-01-200,2套); 中国中车磁浮列车(型号:ZDM-01-300,1套); 中国中车磁浮列车(型号:ZDM-01-400,1套); 中国中车磁浮列车(型号:ZDM-01-500,1套); |
| 应用单位: | 中国中车股份有限公司 |
| 应用地址: | 四川省成都市高新区天府大道中段138号中国中车集团总部 |
| 应用起止时间: | 2015年至今 |
| 应用情况及社会效益: | <p>中国中车集团是中国轨道交通装备制造业的龙头企业，其产品广泛应用于城市轨道交通、城际铁路、高速铁路、城市轨道交通等领域。中国中车集团的产品在国内外市场上享有良好的声誉，并得到了广泛的应用。</p> <p>中国中车集团的产品在国内外市场上享有良好的声誉，并得到了广泛的应用。中国中车集团的产品在国内外市场上享有良好的声誉，并得到了广泛的应用。</p> |
| 特此证明: |  西南科技大学 2015年12月15日 |

应用证明

| | |
|------------|--|
| 项目名称: | 南京中车磁浮列车(型号:ZDM-01-1000,2套); 南京中车磁浮列车(型号:ZDM-01-200,2套); 南京中车磁浮列车(型号:ZDM-01-300,1套); 南京中车磁浮列车(型号:ZDM-01-400,1套); 南京中车磁浮列车(型号:ZDM-01-500,1套); |
| 应用单位: | 西南科技大学材料科学与工程学院 |
| 应用地址: | 四川省绵阳市江油市江油市材料学院 |
| 应用起止时间: | 2015年至今 |
| 应用情况及社会效益: | <p>西南科技大学材料科学与工程学院是国内知名的材料科学研究和人才培养基地。该学院在材料科学领域具有雄厚的科研实力和丰富的教学经验。</p> <p>西南科技大学材料科学与工程学院在材料科学领域具有雄厚的科研实力和丰富的教学经验。西南科技大学材料科学与工程学院在材料科学领域具有雄厚的科研实力和丰富的教学经验。</p> |
| 特此证明: |  西南科技大学 2015年12月15日 |

推广应用证明

| | |
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| 项目名称: | 准静态法比热、热传导系数测定仪(型号:KY-8008,46套) |
| 推广单位: | 四川世纪中科光电技术有限公司 |
| 应用地址: | 四川省成都市经济技术开发区南二环路5号 |
| 推广应用起止时间: | 2012年至今 |
| 推广应用情况及社会效益: | <p>准静态法比热、热传导系数测定仪(型号:KY-8008,46套)经推广应用于国内多所高校、科研院所、企事业单位。该产品在材料科学、物理、化学、生物、医学等领域具有广泛的应用。</p> <p>准静态法比热、热传导系数测定仪(型号:KY-8008,46套)经推广应用于国内多所高校、科研院所、企事业单位。该产品在材料科学、物理、化学、生物、医学等领域具有广泛的应用。</p> |
| 特此证明: |  四川世纪中科光电技术有限公司 2015年12月15日 |

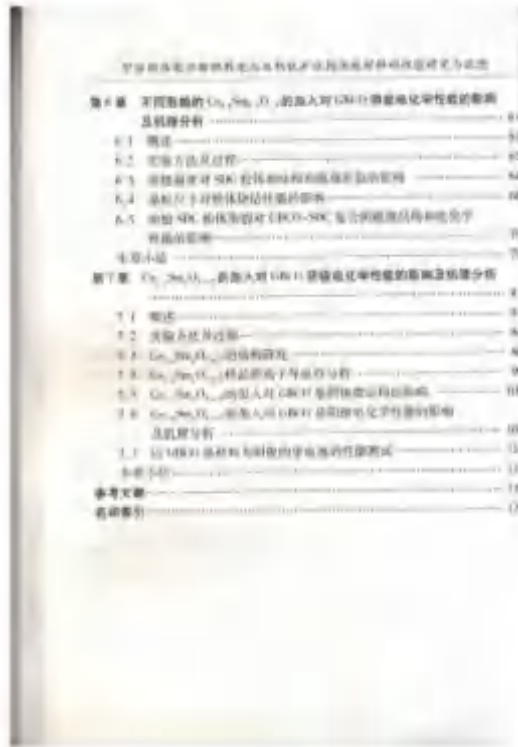
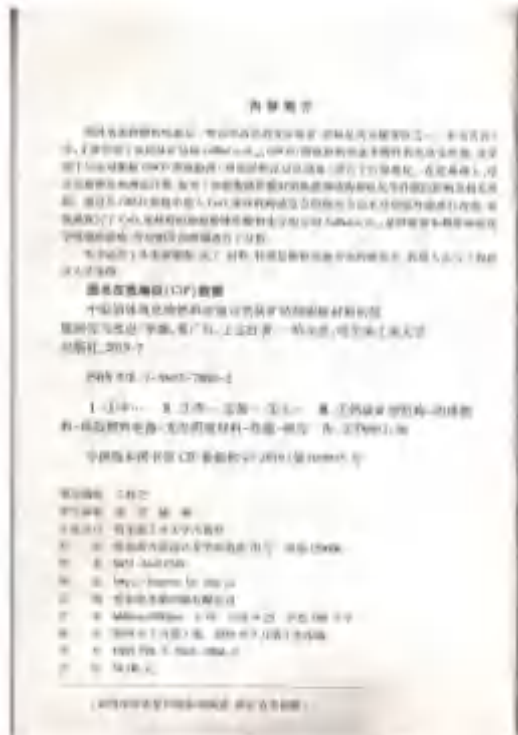
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| 序号 | 教材名称 | 出版社 | 出版年份 |
|----|---|---------------------------|------|
| 1 | 核电子学基础 | 新华大学出版社 | 2021 |
| 2 | 中温固体氧化物燃料电池双钙钛矿结构阴极材料的性能研究与改进 | 哈尔滨工业大学出版社 | 2019 |
| 3 | Introduction to Kinetics of Glow Discharges | Morgan&Claypool Publisher | 2018 |
| 4 | 钽铌酸钾功能晶体 | 科学出版社 | 2018 |
| 5 | In:Fe:Cu:LiNbO ₃ 晶体蓝光光折变与非挥发全息存储 | 哈尔滨工业大学出版社 | 2018 |
| 6 | 分数阶变换与光学图像加密 | 科学出版社 | 2017 |
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| 8 | 光学非线性测量新技术—4f 相位成像技术 | 国防工业出版社 | 2016 |
| 9 | 宇宙中的生命 | 机械工业出版社 | 2016 |
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| 12 | 碳基和类碳超硬材料的第一性原理研究 | 大连理工大学出版社 | 2018 |
| 13 | Graphene Oxide: Physics and Applications | Springer | 2015 |

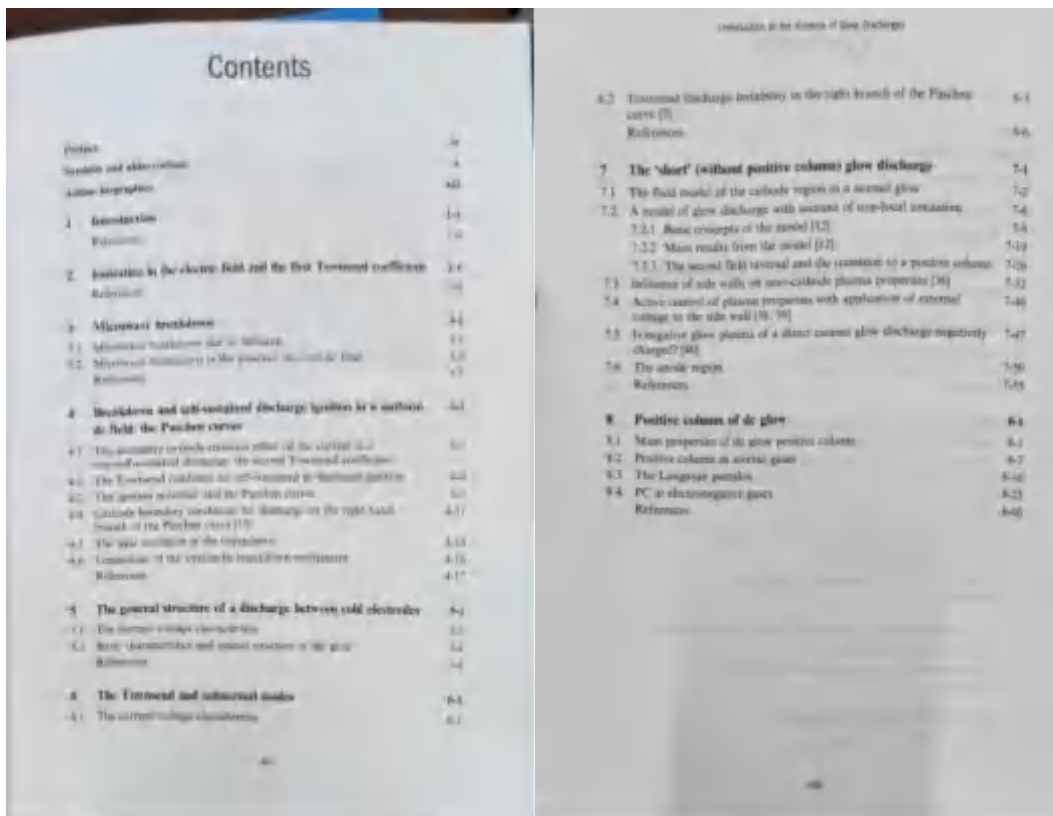
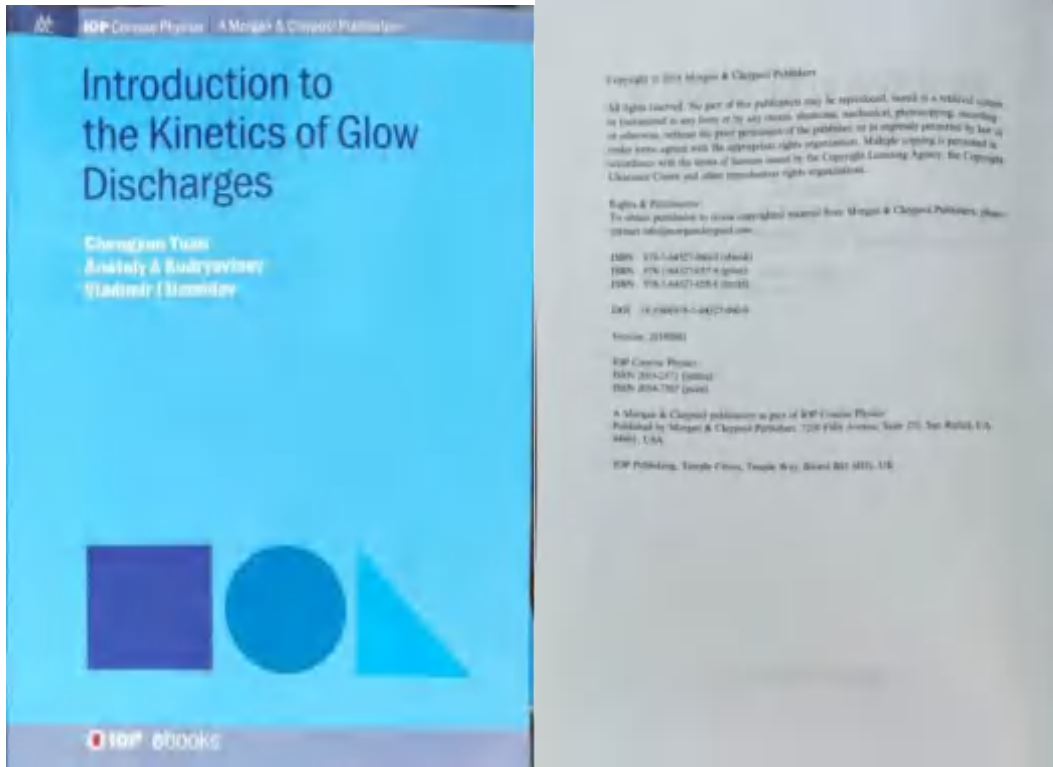
(1) 核电子学基础



(2) 中温固体氧化物燃料电池双钙钛矿结构阴极材料的性能研究与改进



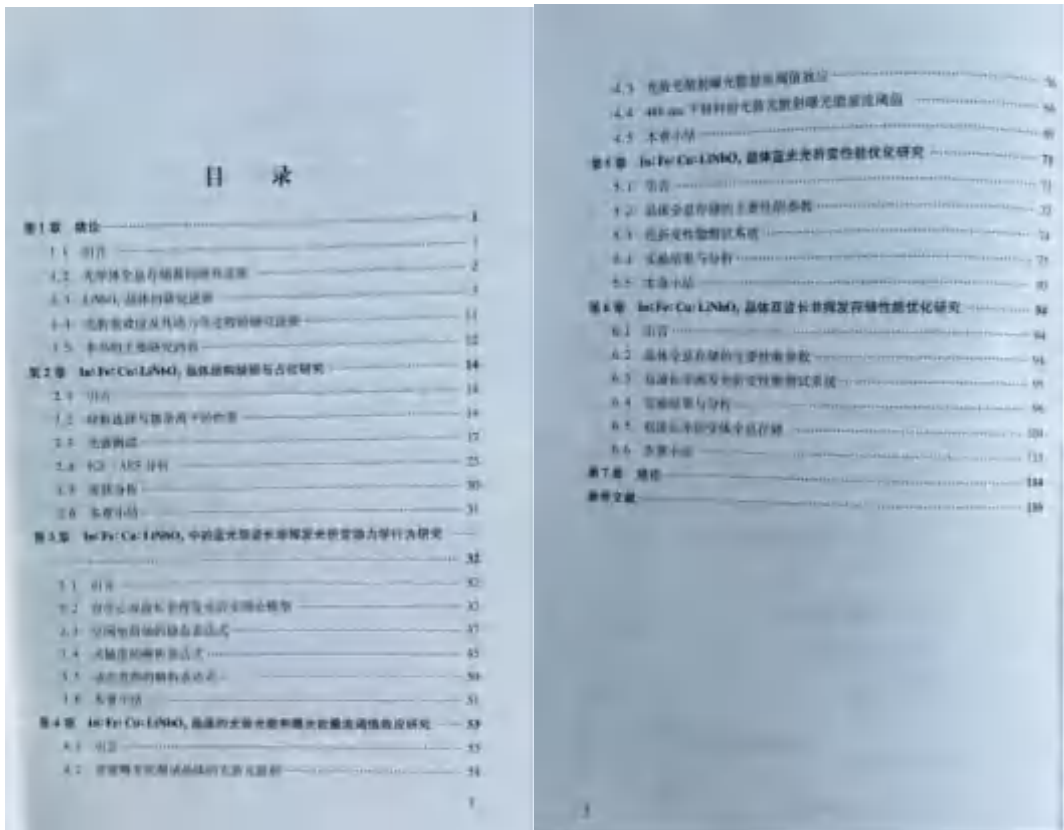
(3) Introduction to the Kinetics of Glow Discharges



(4) 钽铌酸钾功能晶体



(5) In:Fe:Cu:LiNbO₃ 晶体蓝光光折变与非挥发全息存储



(6) 分数阶变换与光学图像加密



内 容 简 介

本书主要介绍分数阶傅里叶变换与分数阶傅里叶变换的分数阶傅里叶变换及其在图像加密和图像恢复中的应用。全书共分4章。第1章介绍分数阶傅里叶变换和分数阶傅里叶变换的傅里叶变换。第2章介绍分数阶傅里叶变换与分数阶傅里叶变换的傅里叶变换。第3章介绍分数阶傅里叶变换与分数阶傅里叶变换的傅里叶变换。第4章介绍分数阶傅里叶变换与分数阶傅里叶变换的傅里叶变换。本书可作为高等院校数学和光学专业师生的教材,也可供从事该领域的研究人员参考。

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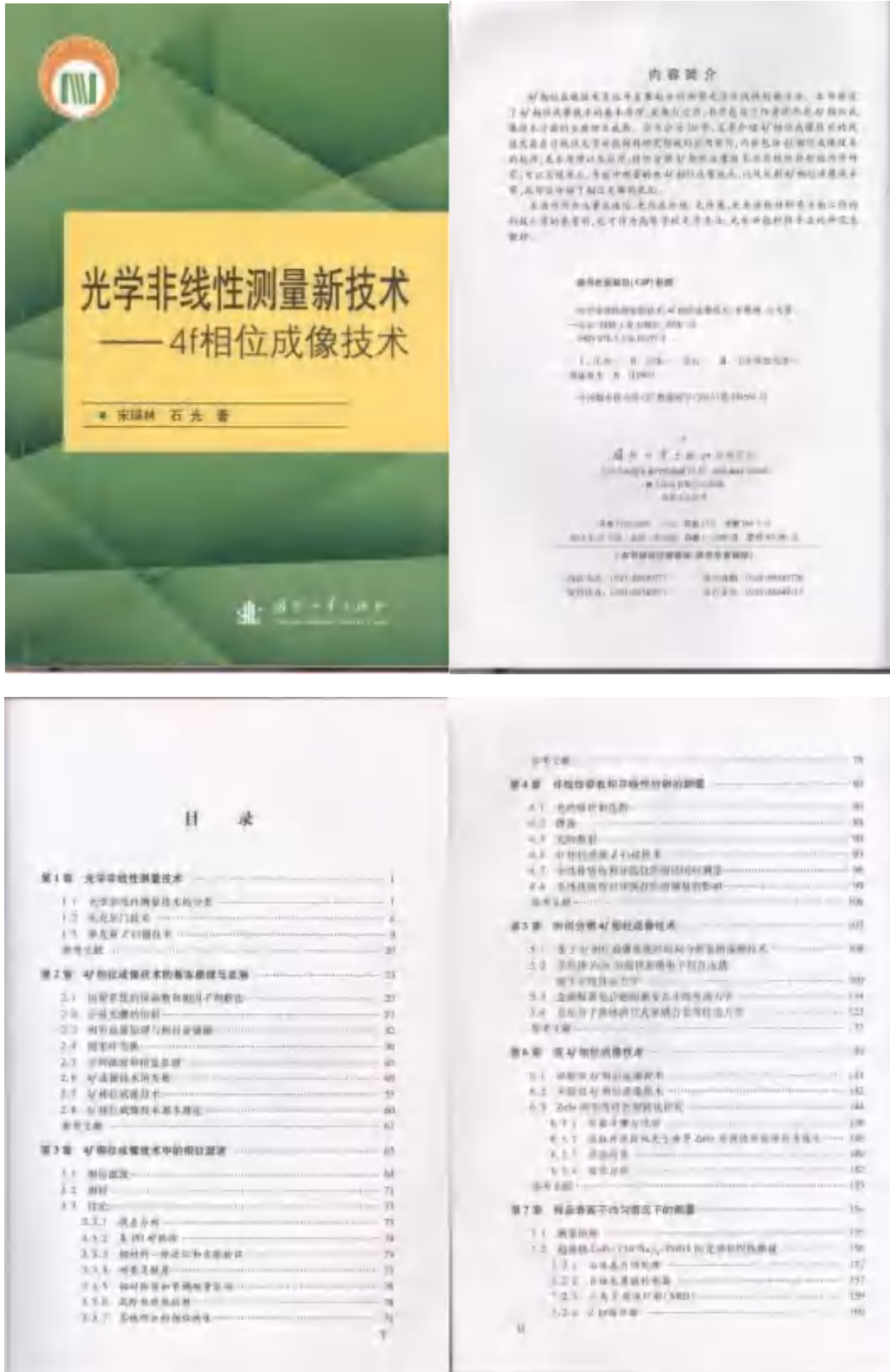
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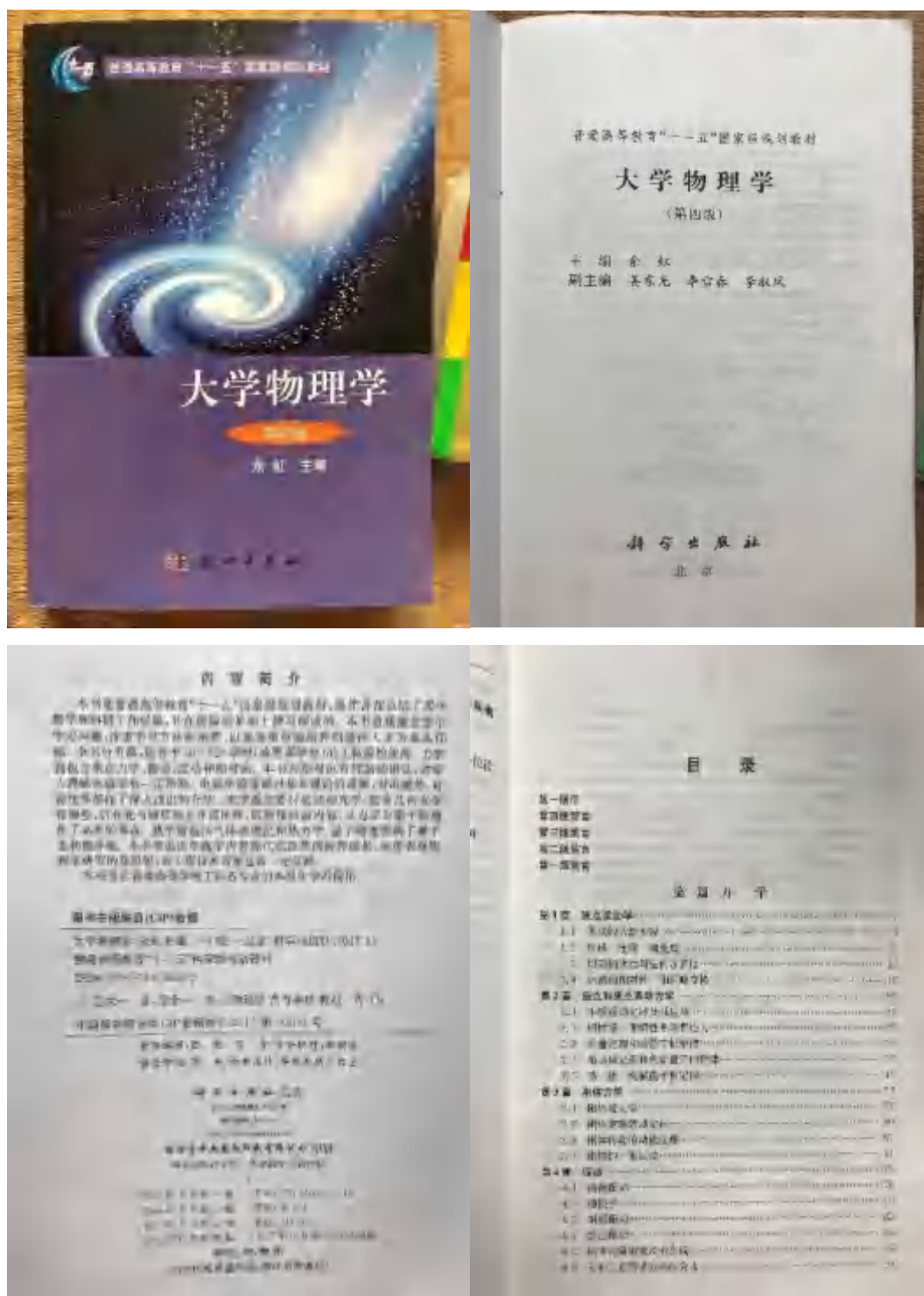
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(1) Ferroelectric crystals with giant electro-optic property enabling ultracompact Q-switches, Science

RESEARCH ARTICLE

RESEARCH ARTICLE

ELECTROOPTICS

Ferroelectric crystals with giant electro-optic property enabling ultracompact Q-switches

Xin Liu¹, Peng Fan^{2,3}, Xie Ma⁴, Dongqiang Wang⁵, Xinyu Jia⁶, Yue Liu⁷, Bin Ku⁸, Lian Qian⁹, Chaoxi Dai¹⁰, Shi Wang¹¹, Wujiao Zhao¹², Chaojie Wu¹³, Xianli Song¹⁴, Hailiang Guo¹⁵, Xudong Li¹⁶, Xiangyang Wu¹⁷, Long Qing Chen¹⁸, Zhen Gu¹⁹, Fei Li²⁰, Bin Han²¹, Shiguo Zhang²²

Recent fast-growth (FTG) crystals, which exhibit extremely high piezoelectricity, are believed to possess high electro-optic (EO) coefficients. However, the optical transparency of related FTG crystals is severely reduced as a result of light scattering and absorption by domain walls, limiting electro-optic applications. Through systematic design of a ferroelectric phase, crystal orientation and poling technique, we successfully removed all light scattering domains with well-aligned or randomly high concentrations of 95.0% α antiferroelectric-like crystal regions, with an ultrahigh EO coefficient r_{33} of 900 picometers per volt (pm V⁻¹), ~30 times as high as that of conventionally used EO crystals. These new crystals exhibit ultracompact EO Q-switches that require very low driving voltages, with superior performance to that of commercial Q-switches. Development of these crystals is important for the portability and low driving voltage EO devices.

Precise control of light propagation and intensity is essential to numerous photonic devices ranging from lasers to optical amplifiers and modulators (1, 2). The rapid and efficient control of optical signals through electrical means is a major challenge. FTG crystals could meet this challenge in their refractive indices in response to an applied electric field (3), via the Pockels effect (4, 5). The primary benefit of the Pockels effect is that it allows a linear correlation between the change in refractive indices and the applied electric field (6), and is a strong effect in crystal (7, 8), of which are critical to a wide range of photonic applications.

Ferroelectric crystals represented by LiNbO₃ (9, 10) and BiFeO₃ (11, 12) are an important component of existing EO devices because of their availability in large size and good temperature stability (3, 13). BiFeO₃ has very high EO that surpasses other EO crystals (14, 15).

Supporting Information for this article is available at www.sciencemag.org. See the article for details. **REFERENCES AND NOTES** 1. M. J. Collins, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 2. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 3. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 4. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 5. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 6. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 7. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 8. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 9. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 10. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 11. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 12. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 13. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 14. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009). 15. S. Wang, *Optical Modulators*, 2nd ed. (Springer, New York, 2009).

Q-switches are (16) believed to be important for laser systems because of their compact size and high efficiency (17). However, the extremely low efficiency EO Q-switch devices (18, 19) and BiFeO₃ crystals (20) are at pm V⁻¹, respectively, under the use of high voltage and/or thick crystal (21, 22) devices, leading to high auxiliary and high voltage device supply and difficulty in miniaturization (23). This issue has become a key obstacle for developing device applications. Therefore, the discovery and development of alternative high EO coefficient ferroelectric crystals to miniaturize the crystal volume and size of EO devices is highly desirable.

Many reported ferroelectric crystals exhibit ultrahigh EO coefficients on the order of 10³ pm V⁻¹, such as BiFeO₃ (24, 25), and Bi₂Fe₂Mo₂O₁₂ (26). PFM-FTG—newly developed ferroelectric phase (17, 27). Despite the promising EO properties, optical clarity of ferroelectric ferroelectric is a long-standing challenge due to the presence of light scattering and absorption of domain walls arising from the difference in refractive indices of the adjacent domains with different orientations, resulting high optical loss or even opaque to the light in some cases (28, 29). In this regard, low driving voltage EO devices have been made to overcome the high-driving-voltage with a thin ferroelectric film with very limited systems (30–32). For example, optical Q-switches based on ferroelectric polymer devices are widely used in many cases without

domain walls. However, the EO coefficients of single-domain crystals are much smaller than those of the crystals poling along a specific crystallographic (111) direction (33). Most recent efforts to manipulate the domain structures of PFM-FTG crystals through an electric field poling, has demonstrated limited approach to largely reduce the light scattering domains with a domain-structured PFM-FTG crystals (34). Although the optical transparency along the crystal direction is greatly improved with a maximum high EO coefficient (1.2–2.0) pm V⁻¹, can be achieved on the basis of this method, the crystal remains opaque along the other orthogonal directions (35, 36). Thus, designing EO devices on the basis of PFM-FTG crystals remains difficult because of the challenges in obtaining ultracompact device implementation with great EO performance in the crystals.

We developed a specific high EO coefficient poling crystals by (111) oriented PFM-FTG (Bi₂Fe₂Mo₂O₁₂, Bi₂Fe₂Mo₂O₁₂, PFM-FTG) crystals. Ultracompact crystals in (111) orientation is naturally ultracompact devices (37) because of anisotropic domain walls, and to achieve high EO coefficient devices the ultracompact crystals from turbulent poling-domain structure. The high EO coefficient PFM-FTG crystals exhibit an ultrahigh EO coefficient r_{33} in the range of 900 to 1000 pm V⁻¹, with a temperature range of 40 to 100°C and a frequency range of 10 to 10³ Hz. We used such poling PFM-FTG crystals to construct an ultracompact ferroelectric EO Q-switch and demonstrated its stability and effectiveness in miniaturization and driving voltage requirements consistent with the size of other EO devices.

Selection of relaxed ferroelectric crystals

We reported the development of PFM-FTG (Bi₂Fe₂Mo₂O₁₂, Bi₂Fe₂Mo₂O₁₂) crystals as a suitable candidate, as they possess comparable high temperature electrical properties (e.g., dielectric and piezoelectric properties) but have much lower residual stress and domain wall densities compared with those of the previously reported relaxed ferroelectric crystals PFM-FTG (38). In addition, high transparency along both the poling and nonpoling directions in PFM-FTG crystals, we selected the PFM-FTG crystal with three perpendicular PFM-FTG crystal with three perpendicular along the crystallographic (111), (111), and (111) directions. We chose (111) field oriented domain structure for the ultracompact crystal poling along the (111) direction (Fig. 1A to 1C) that can reduce volume with poling along the (111) and (111) directions relative to the crystal. Strong EO devices with the ultracompact EO devices with arbitrary of the size of the (111) crystals, whereas they are almost perfectly aligned (parallel along the [111] direction) in the view of the (111)

(2) Light-induced dimension crossover dictated by excitonic correlations, Nature Communications



ARTICLE

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Light-induced dimension crossover dictated by excitonic correlations

Yun Cheng^{1,2,3,4}, Afral Zory^{5,6,7,8}, Jun Li⁴, Wei Xia^{6,7}, Shanfeng Ebran⁹, Wenxuan Zhao⁹, Yellian Li⁸, Fengfeng Qi^{1,2}, Jun Wu^{1,2}, Linqiang Zhao^{1,2}, Pengfei Zhu^{1,2}, Xiao Zou^{1,2}, Tao Jiang^{1,2}, Yanfeng Guo^{1,5}, Lexian Yang⁹, Dong Qian⁶, Wentao Zhang⁶, Anshul Kogar^{10,11}, Michael W. Zuerch^{12,13}, Dao Xiang^{1,2,14,15} & Jie Zhang^{1,2,11,16}

In low-dimensional systems with strong electronic correlations, the application of an ultrashort laser pulse often yields novel phases that are otherwise inaccessible. The central challenge in understanding such phenomena is to determine how dimensionality and many-body correlations together govern the pathway of a non-adiabatic transition. To this end, we examine a layered compound, LiTlSe_2 , whose three-dimensional charge-density waves (3D CDW) state also features exciton condensation due to strong electron-hole interactions. We find that photoexcitation suppresses the equilibrium 3D CDW while creating a nonequilibrium 2D CDW. Remarkably, the dimension reduction does not occur unless bound electron-hole pairs are broken. This relation suggests that excitonic correlations maintain the out-of-plane CDW coherence, settling a long-standing debate over their role in the CDW transition. Our findings demonstrate how optical manipulation of electronic interaction enables one to control the dimensionality of a broken-symmetry order, paving the way for realizing other emergent states in strongly correlated systems.

Key findings In low-dimensional systems with strong electronic correlations, the application of an ultrashort laser pulse often yields novel phases that are otherwise inaccessible. The central challenge in understanding such phenomena is to determine how dimensionality and many-body correlations together govern the pathway of a non-adiabatic transition. To this end, we examine a layered compound, LiTlSe_2 , whose three-dimensional charge-density waves (3D CDW) state also features exciton condensation due to strong electron-hole interactions. We find that photoexcitation suppresses the equilibrium 3D CDW while creating a nonequilibrium 2D CDW. Remarkably, the dimension reduction does not occur unless bound electron-hole pairs are broken. This relation suggests that excitonic correlations maintain the out-of-plane CDW coherence, settling a long-standing debate over their role in the CDW transition. Our findings demonstrate how optical manipulation of electronic interaction enables one to control the dimensionality of a broken-symmetry order, paving the way for realizing other emergent states in strongly correlated systems.

Supplementary Information Supplementary Information is available for this article. For more information on this article, please refer to the article text.

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(3) Multidimensional phase singularities in nanophotonics, Science

RESEARCH

REVIEW SUMMARY

NANOPHOTONICS

Multidimensional phase singularities in nanophotonics

Jiefeng Ni¹, Guo Huang¹, Lei Ming², Zhen-Min Liu¹, Qinghai Tang¹, Yuli Kravtsov¹, Cheng-Wei Qiu¹

BACKGROUND: Optical vortices provide one of the promising solutions for enhancing the capacity of data multiplexing, to meet an unprecedented growth in big data and internet traffic information. The optical vortex possesses phase singularities—an undivided phase in the beam center—used for advanced optical detection by carrying the orbital angular momentum (OAM) of light. The phase (OAM) microscopy has improved resolution for light, which can be used in optical manipulation, quantum-secure communication, microscopic imaging, and optical communication. In particular, the unlimited OAM states can potentially provide more channels for data transmission in optical communication and quantum information processing. Although the beam-like optical systems composed of bulky optical components can stably realize

the generation and detection of vortex beams, many practical applications of photonic OAM require all-optical systems for on-chip devices for matching the scale of optical fibers. Toward this aim, the phase singularities of light has raised considerable attention in multiple dimensions, with the vortex generator and detector being desired in the micro- and nanoscale for integration with compact nanophotonic devices.

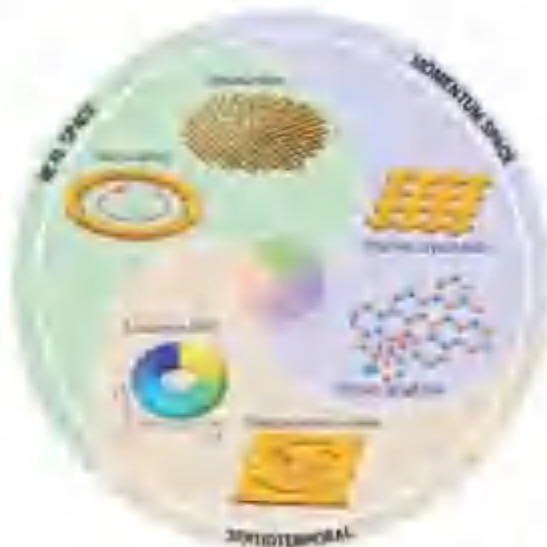
ADVANCES: Here, we focus optical vortices in subwavelength photonic structures spanned in multiple dimensions such as real space, momentum space, and the spatiotemporal domain. In real space, stable OAM beams can be generated by the resonant device such as microcavity and phase photonic metamaterials, and planar optical nanostructures with a helical

phase front. The vortex generation in the form of whispering-gallery mode microresonator can also realize on-chip generation OAM modes from incident light with phase-matching locking. In momentum space, phase singularities also appear on photonic crystal slabs near the condition of band gaps in the momentum space induced by the Bragg-scattering Berry phase. Such subwavelength nanostructures can flexibly generate vortex beams without nanoscale alignment problems. Vortex beams generated by multiphotonic structures can still exhibit features identical to those of conventional bulk devices, including phase modulation and an efficient optical vortex. On the receiving end, nanophotonic structures can well distinguish the OAM superstates in light beams through interference patterns or the photonic crystal effect. Particularly, these microscopic vortex generators and detectors allow staggered multiplexing with optical fibers of on-chip photonic devices for optical communication and quantum information processing. Furthermore, the multidimensional vortex generators with nonlinear, refractive, transparent, and transparent OAM make optical communication possible with higher capacity and improved security.

OUTLOOK: Optical vortices feature fundamentally nonlinear with superlattice structure in nature. The smallest ground state systems has undergone a revolution in the past three decades. The study of vortex beams generated by nano- and microscale structures has been established as a new research direction in optical and integrated nanophotonics. Meanwhile, the research on optical vortices is extended from the conventional real space dimension to the momentum space and time- and frequency-varying structure dimensions. The generation of distinguishable vortex carries new physics and design principles in nanophotonic vortex devices. In addition, on-chip vortex generation, transmission, and detection demonstrate unprecedented functionalities for broader potential applications, especially in optical communication and quantum information processing. From the maturity of conventional fiber technology for carrying optical and electrical information via a stable chip, we expect to see very soon many practical applications and nanoscale products based on these research results.

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Multidimensional phase singularities in on-chip nanophotonics. Various structures of phase singularities spanning space, momentum space, and the spatiotemporal domain realized in a variety of micro- and nanoscale photonic devices. <https://doi.org/10.1038/s41533-023-00271-9> and <https://www.nature.com/articles/s41533-023-00271-9>

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1 of 1

(4) Photoluminescence mechanism of carbon dots: triggering high-color-purity red fluorescence emission through edge amino protonation, Nature Communications



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Photoluminescence mechanism of carbon dots: triggering high-color-purity red fluorescence emission through edge amino protonation

Qing Zhang^{1,2}, Ruoyu Wang¹, Bowen Feng¹, Xiaoxia Zhong¹ & Kostya (Ken) Ostrikov³

Due to complex structure and surface functionalities, photoluminescence mechanisms of Carbon Dots are unknown, and it is challenging to synthesize Carbon Dots to achieve the desired optical properties. Herein, Carbon Dots simultaneously exhibiting high-color-purity (FWHM-24 nm) long wavelength one-photon fluorescence emission at 620 nm and NIR induced two-photon fluorescence emission at 630 and 680 nm are prepared by edge amino protonation treatment. Systematic analysis reveals that the protonation of 2,3-diaminophenazine changes the molecular state of Carbon Dots, decreases the photon transition band gap, and triggers red fluorescence emission with the dramatically narrowed peak width. As the oxidation products of reactant α -phenylenediamine, the emergence of 2,3-diaminophenazine as a photoluminescence determiner suggests that fluorophore products of precursor conversion are viable determinants of the desired luminescence properties of Carbon Dots. This work shows a new way for predicting and controlling photoluminescence properties of Carbon Dots, and may guide the development of tunable Carbon Dots for a broad range of applications.

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(5) High-efficiency broadband achromatic metalens for near-IR biological imaging window, Nature Communications



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High-efficiency broadband achromatic metalens for near-IR biological imaging window

Yujie Wang^{1,5}, Qinmao Chen^{1,5}, Wenhong Yang^{1,5}, Ziheng Ji^{1,5}, Limin Jin¹, Xing Ma¹, Qinghai Song¹, Alexandra Boltasseva², Jiecai Han³, Vladimir M. Shalaev² & Shumin Xiao^{1,3,4✉}

Over the past years, broadband achromatic metalenses have been intensively studied due to their great potential for applications in consumer and industry products. Even though significant progress has been made, the efficiency of technologically relevant silicon metalenses is limited by the intrinsic material loss above the bandgap. In turn, the recently proposed achromatic metalens utilizing transparent, high-index materials such as titanium dioxide has been restricted by the small thickness and showed relatively low focusing efficiency at longer wavelengths. Consequently, metalens-based optical imaging in the biological transparency window has so far been severely limited. Herein, we experimentally demonstrate a polarization-insensitive, broadband titanium dioxide achromatic metalens for applications in the near-infrared biological imaging. A large-scale fabrication technology has been developed to produce titanium dioxide nanopillars with record-high aspect ratios featuring pillar heights of 1.5 μm and $\sim 90^\circ$ vertical sidewalls. The demonstrated metalens exhibits dramatically increased group delay range, and the spectral range of achromatism is substantially extended to the wavelength range of 650–1000 nm with an average efficiency of 77.1%–88.5% and a numerical aperture of 0.24–0.1. This research paves a solid step towards practical applications of flat photonics.

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(6) Engineering single-atomic ruthenium catalytic sites on defective nickel-iron layered double hydroxide for overall water splitting, Nature Communications



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Engineering single-atomic ruthenium catalytic sites on defective nickel-iron layered double hydroxide for overall water splitting

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Rational design of single atom catalyst is critical for efficient sustainable energy conversion. However, the atomic-level control of active sites is essential for electrocatalytic materials in alkaline electrolyte. Moreover, well-defined surface structures lead to in-depth understanding of catalytic mechanisms. Herein, we report a single-atomic-site ruthenium stabilized on defective nickel-iron layered double hydroxide nanosheets (Ru₁/D-NiFe LDH). Under precise regulation of local coordination environments of catalytically active sites and the existence of the defects, Ru₁/D-NiFe LDH delivers an ultralow overpotential of 18 mV at 10 mA cm⁻² for hydrogen evolution reaction, surpassing the commercial Pt/C catalyst. Density functional theory calculations reveal that Ru₁/D-NiFe LDH optimizes the adsorption energies of intermediates for hydrogen evolution reaction and promotes the O–O coupling at a Ru–O active site for oxygen evolution reaction. The Ru₁/D-NiFe LDH as an ideal model reveals superior water splitting performance with potential for the development of promising water-alkali electrocatalysts.

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(7) Direct observation of chaotic resonances in optical microcavities, Light Science & Applications

ARTICLE

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Direct observation of chaotic resonances in optical microcavities

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Abstract

Optical microcavities play a significant role in the study of classical and quantum chaos. To date, most experimental investigations of time-reversal wave dynamics have focused on the properties of their inputs and outputs, without directly interrogating the dynamics and the associated mode patterns inside. As a result, the key information is rarely retrieved with certainty, which significantly hinders the verification and understanding of the actual chaotic motion. Here, we demonstrate a simple and robust approach to directly and rapidly map the internal mode patterns in chaotic microcavities, by introducing a local index perturbation through a pump laser, via hybrid a spectral response of optical microcavities that is proportional to the internal field distribution. With this technique, chaotic modes with staggered mode spacing can be distinguished. Consequently, a complete chaos-assisted tunneling (CAT) and its time-reversed process are experimentally verified in the optical domain with unprecedented certainty.

Introduction

The detailed understanding of mixed-phase space with both regular and chaotic motions is essential for a wide range of physical systems, ranging from atomic and molecular physics to mesoscopic science and even astro-physics^{1–4}. Among the numerous quantum chaotic phenomena in mixed-phase space, the chaos-assisted tunneling (CAT) between distinct regular phase space regions separated by a chaotic sea is particularly interesting and has been intensively studied in molecular dynamics, cold atom assemblies, and whispering gallery modes (WGM) microcavities^{5–7}. Among these systems, WGM microcavities have a series of advantages to study CAT, including easy control of internal quantum or classical dynamics and multiple manifestations of CAT

beyond the regular level splitting^{8–14}. In past decades, the rapid progress in WGM microcavities^{15–22} has not only enriched fundamental physics such as CAT, resonance-assisted tunneling²³, and turnstile transport²⁴, but also enabled many practical applications from broadband coupling^{25,26} to integrated laser sources^{27–30}, sensors^{30,31} and all-optical switches³².

Despite these achievements, the conventional study of mixed-phase space in optical WGM microcavities focuses on its ancillary signatures in the inputs and the outputs, without directly interrogating the dynamics and the associated mode patterns inside. These techniques include, for example, comparing the experimentally observed free spectral range (FSR) and far-field radiation patterns with theoretical predictions^{33–35,36,37,38}. However, such comparisons often lead to educated conjectures instead of conclusive affirmations, especially when multiple mode groups have similar FSRs and far-field patterns^{39,40}. The association of internal mode patterns with experimentally observed resonances is hence crucial to provide a direct picture of quantum chaos in optical microcavities.

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Suppressing meta-holographic artifacts by laser coherence tuning

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Abstract

A metasurface hologram combines the spatial resolution and large viewing angles with a planar form factor and compact foot. However, it suffers coherent artifacts originating from electro-magnetic cross-talk between closely packed meta-atoms and fabrication defects of nanoscale features. Here, we introduce an efficient method to suppress all artifacts by fine-tuning the spatial coherence of illumination. Our method is implemented with a degenerate laser, which allows a precise and continuous tuning of the spatial coherence over a wide range, with little variation in the emission spectrum and total power. We find the optimal degree of spatial coherence to suppress the coherent artifacts of a meta-hologram while maintaining the image fidelity. This work paves the way to compact and dynamic holographic displays free of coherent defects.

Introduction

Artificial metasurfaces, comprised of a two-dimensional (2D) array of subwavelength scatterers, have shown unprecedented ability in controlling optical waveform and converting conventional bulky optical elements into planar thin films^{1–3}. One prominent example is the metasurface hologram (meta-hologram)^{4–7}. An ideal metasurface is capable of reconstructing a three-dimensional (3D) holographic image with a high spatial resolution and large viewing angles, while suppressing high-order diffraction^{8–12}. Very recently, multi-color, multiplexed, and dynamic meta-holograms have been proposed and demonstrated, illustrating a great potential in information processing, 3D display, high-density data storage, and optical image encoding^{13–16}. Despite of these remarkable advances, the road to practical applications of meta-holograms is hindered by coherent artifacts. Such artifacts originate from near-field interactions of subwavelength scatterers

(meta-atoms), fabrication defects and phase dislocations, causing image distortion and degradation^{17,18}. While coherent artifacts and speckle noise are well-known issues for conventional holography, they are more significant in regard to meta-holography, as close packing of meta-atoms enhances the cross-talk and fabrication of nanoscale features is susceptible to error. Such artifacts cause severe distortions of holographic images, which are extremely difficult to correct. Recently, machine-learning based optimization techniques have been applied to high-performance metasurface design^{19–21}. They require a large amount of training data, which are difficult and expensive to acquire for large-scale meta-holograms. Moreover, the coherent artifacts caused by random fabrication imperfections cannot be removed by optimizing the meta-hologram design. An alternative way of suppressing coherent artifacts is adjusting the spatial and/or temporal coherence of illuminating light²². In general, it is more efficient to suppress speckle noise by reducing spatial coherence than temporal coherence²³. While lowering the temporal coherence with broadband illumination provides spectral compounding, the different wavelengths generate different radial scalings of speckle patterns resulting in a deficient suppression of coherent artifacts²⁴. Previously, lowering the spatial

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(9) Transparent ferroelectric crystals with ultrahigh piezoelectricity, Nature

Article

Transparent ferroelectric crystals with ultrahigh piezoelectricity

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Transparent piezoelectrics are highly desirable for numerous hybrid ultrasonic-optical devices ranging from photoacoustic imaging transducers to transparent actuators for light applications^{1–4}. However, it is challenging to achieve high piezoelectricity and perfect transparency simultaneously because most high-performance piezoelectrics are ferroelectrics (containing high-density light-scattering domain walls). Here, through a combination of phase-field simulations and experiments, we demonstrate a relatively simple method of using an alternating current electric field to engineer the domain structures of originally opaque rhombohedral Pb(Mg_{1-x}Nb_x)O₃-PbTiO₃ (PMN-PT) crystals to simultaneously generate near-perfect transparency, an ultrahigh piezoelectric coefficient d_{33} (greater than 2,300 picocoulombs per newton), an excellent electromechanical coupling factor k_{33} (about 94 per cent) and a large electro-optical coefficient r_{33} (approximately 220 picometres per volt), which is far beyond the performance of the commonly used transparent ferroelectric crystal LiNbO₃. We find that increasing the domain size leads to a higher d_{33} value for the [001]-oriented rhombohedral PMN-PT crystals, challenging the conventional wisdom that decreasing the domain size always results in higher piezoelectricity^{5–10}. This work presents a paradigm for achieving high transparency and piezoelectricity by ferroelectric domain engineering, and we expect the transparent ferroelectric crystals reported here to provide a route to a wide range of hybrid device applications, such as medical imaging, self-energy harvesting, touch screens and insular robotic devices.

Achieving simultaneously high piezoelectricity and perfect transparency in a piezoelectric material has long been a challenge. For example, traditional high-performance piezoelectric transducers are typically made from piezoelectric ceramic transducers and/or polymers with chemical compositions that are low in their morphotropic phase boundaries (MPBs), such as Pb(Zr_{0.52}Ti_{0.48})O₃ (PZT) ceramics and domain-engineered PMN-PT crystals. These materials possess very high d_{33} and k_{33} values^{11–14}, but they are usually opaque in the visible light spectrum. On the other hand, the commonly used transparent piezoelectric LiNbO₃ crystal and polyvinylidene fluoride (PVDF) (polymer¹⁵) have good transparency but much lower d_{33} and k_{33} values (LiNbO₃: $d_{33} = 40$ pC N⁻¹, $k_{33} = 47\%$; PVDF: $d_{33} = 20$ pC N⁻¹, $k_{33} = 16\%$) that severely limit the acoustic wave level, bandwidth and sensitivity of the transducers.

In addition to the intrinsic effects, such as pinning and grain boundaries, which are ubiquitous in ceramics, the poor transparency in PZT ceramics and domain-engineered PMN-PT crystals is closely associated with light scattering and reflection from their ferroelectric domain walls. There are two possible approaches to reducing the

light-scattering domain walls. The first is to pole a ferroelectric crystal along the polar direction to achieve a single-domain state. However, the d_{33} value of such single-domain PZT ceramic is generally very low^{16–18}—much lower than that of [001]-poled multi-domain rhombohedral PMN-PT crystals (~2,300 pC N⁻¹). In principle, one could first pole a rhombohedral PMN-PT crystal along the [111] direction to achieve a single-domain state with good transparency, then rotate the crystal to the [001] direction to guarantee high longitudinal piezoelectricity. However, this approach is not feasible in practice (see Methods for a detailed explanation). The second approach is to dramatically reduce the domain sizes by breaking the domains into polar nanoregions with spatial sizes at the order of nanometres (much smaller than the wavelength of visible light), thus greatly improving their light transparency—as observed in Li-doped PZT¹⁹. However, improving the transparency using polar nanoregions is achieved at the expense of a markedly reduced remanent polarization and thus very low d_{33} values. In fact, despite more than 30 years of effort, optimal hybridization of high-performance piezoelectrics have not been realized.

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(11) Ultrafast control of vortex microlasers, Science

RESEARCH

OPTICS

Ultrafast control of vortex microlasers

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The development of efficient and quantum information-processing technology calls for on-chip, integrated sources of structured light. Although integrated vortex microlasers have been previously demonstrated, they remain static and possess relatively long lasing lifetimes, limiting their suitability for high-speed optical communication and computing. We introduce permittivity-based vortex microlasers and demonstrate their application in ultrafast optical switching at room temperature. By exploiting both time symmetry and helicity properties, we reveal that the vortex laser lasing can be initiated by linearly polarized beam lasing, *in vivo versa*, with switching times of 1 to 15 picoseconds and energy consumption that is orders of magnitude lower than is previously demonstrated in optical switching. Our results provide an approach that breaks the long-standing trade-off between low energy consumption and high-speed photonic switching, introducing vortex microlasers that are suitable at ambient temperatures.

Beamlines of their unusual, or topological, vortex beams with different topological charges have been proposed as an effective approach to revolutionize classical and quantum communications (1–8). Vortex beams with well-defined topological charges have been developed using external phase elements, such as optical glass plates, ring-shaped grating helix gratings, and metasurfaces (9–12), recently driven by the demand for compact displays and high-density integration. Structured vortex microlasers have attracted much attention (3–7, 13). Compact vortex microlasers are usually created by transferring conventional optical systems into optical waveguide (4) or micropillar chains (7) and correlating them with additional asymmetry elements (5). Although the structural performance in both directional output and generation efficiency of orbital angular momentum (OAM) beams is critical, the quality (Q) factors of such vortex microlasers are strongly reduced by scattering losses and, thus, their energy consumption is large (14–17). Additionally, on-chip integrated vortex microlasers are either static (3, 6) or incapable of sensitive chirality via the circularly polarized optical pump (7), making them unsuitable for ultrafast optical switches (8). Finally, because of the resolution limit of the optical pump, there

seems to be a trade-off between low energy consumption and high modulation speed in microlasers, which restricts their applications in modern optical computing and optical communication.

In this work, we solve these problems by employing bound states in the continuum (BICs), which represent special solutions of the wave equation where the wave function exhibits localization but radiates light (18, 19).

In optical systems, BICs appear through interference between localized resonances and radiation modes, and they have been observed in the form of quasi-BICs in many systems, ranging from isolated nanoparticles to periodic structures (20–22). In addition to ultrahigh Q factors and low-threshold lasing, it has been predicted that the BIC modes can possess vortex behavior with different topological charges, which are important for vortex beams (23–25). These features make BICs very attractive for application in active photonics. We employ the specific characteristics of the topologically protected BICs and demonstrate the ultrafast control of permittivity-based vortex microlasers at room temperature.

OAM microlasers are created using a 220-nm-thick permittivity-modulated (MAPM) film, patterned with a square array of circular holes (Fig. 1A). The whole structure is placed on a glass substrate ($n_{\text{glass}} = 1.5$, where n is the refractive index) and coated with polymer methacrylate (PMMA) ($n_{\text{PMMA}} = 1.49$). The radius of the air holes is $R = 105$ nm, and the lattice spacing is $a = 260$ nm. We calculate the resonance of the transverse magnetic (TM) polarized field within the permittivity-modulated film (Fig. 1B). Such a film supports Q factor within the gap region range of MAPM, permittivity (Fig. 1C). By choosing

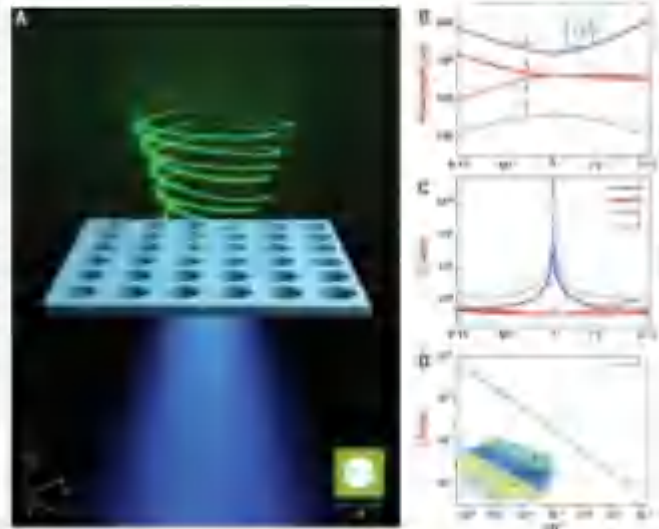


Fig. 1. Design and control of the quasi-BIC modes. (A) Schematic of the designed permittivity-modulated film. The microlaser is constructed by a thin film with a square lattice of holes on the surface of the substrate. (B) Dispersion relation of the TM modes. (C) Comparison of the dispersion relation of the TM modes. The red curve shows the real dispersion curve of the quasi-BIC modes, and the blue curve shows the imaginary dispersion curve of the quasi-BIC modes. (D) Field distribution of the TM modes for the quasi-BIC modes with a quality factor exceeding that of the conventional modes.

Author contributions: C.H., C.Z., S.X., and Q.S. designed the experiment. C.H., C.Z., S.X., Y.W., F.F., Y.L., N.Z., Q.Q., M.J., X.H., L.G., Y.-K.W., and Q.S. performed the experiment. C.H., C.Z., S.X., Y.W., F.F., Y.L., N.Z., Q.Q., M.J., X.H., L.G., Y.-K.W., and Q.S. analyzed the data. C.H., C.Z., S.X., Y.W., F.F., Y.L., N.Z., Q.Q., M.J., X.H., L.G., Y.-K.W., and Q.S. wrote the paper. C.H., C.Z., S.X., Y.W., F.F., Y.L., N.Z., Q.Q., M.J., X.H., L.G., Y.-K.W., and Q.S. contributed equally to this work. C.H., C.Z., S.X., Y.W., F.F., Y.L., N.Z., Q.Q., M.J., X.H., L.G., Y.-K.W., and Q.S. are co-first authors. C.H., C.Z., S.X., Y.W., F.F., Y.L., N.Z., Q.Q., M.J., X.H., L.G., Y.-K.W., and Q.S. are co-corresponding authors. C.H., C.Z., S.X., Y.W., F.F., Y.L., N.Z., Q.Q., M.J., X.H., L.G., Y.-K.W., and Q.S. are co-senior authors. C.H., C.Z., S.X., Y.W., F.F., Y.L., N.Z., Q.Q., M.J., X.H., L.G., Y.-K.W., and Q.S. are co-supervisors. C.H., C.Z., S.X., Y.W., F.F., Y.L., N.Z., Q.Q., M.J., X.H., L.G., Y.-K.W., and Q.S. are co-mentors. C.H., C.Z., S.X., Y.W., F.F., Y.L., N.Z., Q.Q., M.J., X.H., L.G., Y.-K.W., and Q.S. are co-advisors. C.H., C.Z., S.X., Y.W., F.F., Y.L., N.Z., Q.Q., M.J., X.H., L.G., Y.-K.W., and Q.S. are co-supervisors. C.H., C.Z., S.X., Y.W., F.F., Y.L., N.Z., Q.Q., M.J., X.H., L.G., Y.-K.W., and Q.S. are co-mentors. C.H., C.Z., S.X., Y.W., F.F., Y.L., N.Z., Q.Q., M.J., X.H., L.G., Y.-K.W., and Q.S. are co-advisors.

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Reprogrammable meta-hologram for optical encryption

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Meta-holographic encryption is a potentially important technique for information security. Despite rapid progresses in multi-tasked meta-holograms, the number of information channels available in metasurfaces is limited, making meta-holographic encryption vulnerable to some attacking algorithms. Herein, we demonstrate a re-programmable metasurface that can produce arbitrary holographic images for optical encryption. The encrypted information is divided into two matrices. These two matrices are imposed to the incident light and the metasurface, respectively. While the all-dielectric metasurface is static, the phase matrix of incident light provides additional degrees of freedom to precisely control the eventual functions at will. With a single Si metasurface, arbitrary holographic images and videos have been transported and decrypted. We hope that this work paves a more promising way to optical information encryption and authentication.

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Universal growth of ultra-thin III–V semiconductor single crystals

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Ultra-thin III–V semiconductors, which exhibit intriguing characteristics, such as two-dimensional (2D) electron gas, enhanced electron–hole interaction strength, and strongly polarized light emission, have always been anticipated in future electronics. However, their inherent strong covalent bonding in three dimensions hinders the layer-by-layer exfoliation, and even worse, impedes the 2D anisotropic growth. The synthesis of desirable ultra-thin III–V semiconductors is hence still in its infancy. Here we report the growth of a majority of ultra-thin III–V single crystals, ranging from ultra-narrow to wide bandgap semiconductors, through enhancing the interfacial interaction between the III–V crystals and the growth substrates to proceed the 2D layer-by-layer growth mode. The resultant ultra-thin single crystals exhibit fascinating properties of phonon frequency variation, bandgap shift, and giant second harmonic generation. Our strategy can provide an inspiration for synthesizing unexpected ultra-thin non-layered systems and also drive exploration of III–V semiconductor-based electronics.

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(15) Engineering active sites on hierarchical transition bimetal oxides/sulfides heterostructure array enabling robust overall water splitting, Nature Communications



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Engineering active sites on hierarchical transition bimetal oxides/sulfides heterostructure array enabling robust overall water splitting

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Rational design of the catalysts is impressive for sustainable energy conversion. However, there is a grand challenge to engineer active sites at the interface. Herein, hierarchical transition bimetal oxides/sulfides heterostructure arrays interacting two-dimensional MoD_x/MoS₂ nanosheets attached to one-dimensional NiO_x/Ni₃S₂ nanorods were fabricated by oxidation/hydrogenation-induced surface reconfiguration strategy. The NiMoO_x/NiMoS heterostructure array exhibits the overpotentials of 38 mV for hydrogen evolution and 186 mV for oxygen evolution at 10 mA cm⁻², even surviving at a large current density of 500 mA cm⁻² with long-term stability. Due to optimized adsorption energies and accelerated water splitting kinetics by theory calculations, the assembled two-electrode cell delivers the industrially relevant current densities of 500 and 1000 mA cm⁻² at record low cell voltages of 1.60 and 1.66 V with excellent durability. This research provides a promising avenue to enhance the electrocatalytic performance of the catalysts by engineering interfacial active sites toward large-scale water splitting.

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(16) Ultra-large electric field–induced strain in potassium sodium niobate crystals, Science Advances

SCIENCE ADVANCES | RESEARCH ARTICLE

MATERIALS SCIENCE

Ultra-large electric field–induced strain in potassium sodium niobate crystals

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Electromechanical coupling in piezoelectric materials allows direct conversion of electrical energy into mechanical energy and vice versa. Here, we demonstrate lead-free (K_{1-x}Na_x)/NiO₃ single crystals with an ultra-high large-signal piezoelectric coefficient d_{33}^* of 3000 pm V⁻¹, which is superior to the highest value reported in state-of-the-art lead-based single crystals (~2500 pm V⁻¹). The enhanced electromechanical properties in our crystals are realized by an engineered compositional gradient in the as-grown crystal, allowing notable reversible 100–1200-strain with stretch. Moreover, our crystals exhibit temperature-insensitive strain performance within the temperature range of 25 °C to 125 °C. The enhanced temperature stability of the response also allows the materials to be used in a wider range of applications that exceed the temperature limits of current lead-based piezoelectric crystals.

INTRODUCTION

Piezoelectricity, which arises from the polarization, electric field, and electric field, has many natural applications in the fields of information and communication, medical imaging, and diagnosis. Instead of their past of processing and low-cost piezoelectricity dominates the market. The piezoelectric properties of piezoelectric piezoceramics is an essential attribute of the piezoelectric response over all positive crystalline spin orientations of the corresponding single crystal. The crystal structure of the neighboring atoms associated with the piezoelectric response makes their property significant. This coupling in the piezoelectricity, which is electric field-induced strain of piezoceramics has generally failed to exceed a value of 1.0% (1, 2). Piezoelectric single crystals, whereas the piezoelectricity can be realized by stretching the crystal in an internal direction, are of great interest because of their superior piezoelectric coupling and piezoelectric properties (3–7). In particular, piezoceramics are the lead-based values, ferroelectric single crystals (1–xPbMg_{1-x}Nb_{1-x})/Bi_{1-x}Fe_{1-x}PO₄ (PBi_{1-x}Bi_{1-x}PO₄) and (1–xPb)(2–x)Nb_{1-x}Bi_{1-x}PO₄ (PZS-xPT) piezoceramics exhibit a large strain of 1.2% piezoelectric field of 12 kV cm⁻¹, giving the world's highest piezoelectric coefficient (d_{33}^*) of 2800 pm V⁻¹. This unusual achievement facilitates the development of electromechanical devices with growing importance in crystal growth technologies to manufacture the field and force sensors. While PZS-xPT and PZT-xPT single crystals are used in the high-end applications such as ultrasonic medical imaging and naval sonar, they have a limited future in consumer devices due to environmental regulations limiting the use of lead (8). This has driven an increasing research trend toward developing lead-free-based materials (9–12).

The ferroelectric single crystals (PMN-xPT and PZS-xPT) usually exhibit enhanced piezoelectricity when the composition approaches the morphotropic phase boundary (MPB), thereby a factorial increase in the piezoelectricity (13, 14). The piezoelectricity is significantly enhanced. Another important characteristic of these piezoelectric ferroelectric single crystals is the presence of polar nanoregions (PNRs), which contribute to 70 to 80% of the dipole and piezoelectric response (15). In addition to the inherent mechanism associated with the MPB and PNRs, alternative approaches include an engineered domain configuration (3, 16–20) and defect/strain-induced stretching (21) to enhance the piezoelectric response. However, there is a challenge for piezoelectric properties that can be achieved in PMN-xPT and PZS-xPT single crystals on the basis of the above-mentioned mechanisms and methods.

In this work, we propose to take advantage of an engineered compositional gradient in lattice because the combination of domain and microstructural defects-based ferroelectric (Pb_{1-x}Bi_{1-x}PO₄, PMN-xPT) single crystals. In this large large-signal piezoelectric coefficient d_{33}^* of 3000 pm V⁻¹ under a relatively small electric field of 7 kV cm⁻¹ is obtained. The response is at least three times larger than that of PMN-xPT single crystals under the same electrical stimulus. Moreover, temperature-insensitive ferroelectric field-induced strain performance is observed within the temperature range of 25 °C to 125 °C due to a high Curie temperature of 480 °C.

RESULTS AND DISCUSSION

The ultra-large piezoelectric properties

An ultra-large strain of 0.9% and an ultrahigh d_{33}^* of 3000 pm V⁻¹ under electric field of 7 kV cm⁻¹ are found in K_{0.97}Na_{0.03}(NiO₃) single crystals (KNSNi), as shown in Fig. 1B. Representative lead-based crystals were also compared (3, 8). The high large-signal piezoelectric coefficient d_{33}^* is at least three times the value of the related PE crystals. The hysteresis response was confirmed on KNSNi samples from repeated crystal growth and after by Slack Ziegler digital holographic interferometry, which shows consistent results with stress measured by a traditional force transducer (Fig. 1D). The temperature dependence of the d_{33}^* value is shown in Fig. 1E, which exhibits superior decrease pattern to that of PZS-xPT (3) and PMN-xPT crystals (Fig. S4, compare a high Curie temperature

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(17) High-performance position-sensitive detector based on the lateral photoelectrical effect of two-dimensional materials, Light Science & Applications

REVIEW ARTICLE

Open Access

High-performance position-sensitive detector based on the lateral photoelectrical effect of two-dimensional materials

Chang Hu¹, Nanjie Wang¹ and Bin Song²

Abstract

Two-dimensional (2D) materials such as graphene and transition-metal dichalcogenides have been extensively studied because of their superior electronic and optical properties. Recently, 2D materials have shown great practical application in position-sensitive detector (PSD), originating from the lateral photoelectrical effect of the materials of junctions. The high quantum sensitivity and elevated photoresponse of PSDs based on 2D materials, especially compatibility with Si technology, may enable various complementary applications. In this issue, recent studies of PSDs based on 2D materials are summarized, providing a promising route for high-performance PSDs.

Introduction

Since Schottky and Wallmark discovered and presented the lateral photovoltaic effect (LPE) in *p-n* junctions¹, intensive research has been conducted on the LPE in diverse applications, including space exploration, environmental monitoring, and optical engineering. In the LPE, when one side of a *p-n* junction is illuminated by a light spot, the optically excited electron-hole pairs will be separated up and down by the built-in electric field. Notably, the formation mechanism of the built-in field in position-sensitive detectors (PSDs) with different architectures is different. As shown in Fig. 1a, the PSD based on a *p-n* junction uses the built-in field formed at the interface due to different band structures to separate the photogenerated carriers. The built-in field in the PSD based on the insulator-coated substrate (such as a graphene transistor on a SiO₂/Si substrate) is formed by band energy banding induced by the localized interface states, and the photogenerated carriers accumulate at the SiO₂/Si interface and then diffuse laterally, as shown in Fig. 1b.

The mechanism of carrier lateral diffusion can be explained by the model shown in Fig. 1c. The photo-generated electrons enter the *n*-type semiconductor, while the photogenerated holes enter the *p*-type semiconductor. Then, these generated carriers diffuse laterally from the light spot to balance the electric field along the layer²⁻⁵. In general, the relationship between the total number of photogenerated electrons n_0 and that of photogenerated electrons transmitted into *n*-type semiconductor N_0 can be described as follows:

$$N_0 = n_0 [1 - P^{1-\alpha}] \tag{1}$$

where P is the probability of photogenerated electrons entering the *n*-type semiconductor, ρ is the laser power, and α is a time-dependent coefficient.

The electron diffusion equation at position r is

$$D_n \frac{d^2 N(r)}{dr^2} = \frac{N(r)}{\tau_n} \tag{2}$$

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Keywords

Position-sensitive detector, Lateral photoelectrical effect, Two-dimensional materials, Graphene, Transition-metal dichalcogenides

(18) Momentum-Topology-Induced Optical Pulling Force, Physical Review Letters

Momentum-Topology-Induced Optical Pulling Force

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We report an ingenious mechanism to obtain robust optical pulling force by k -space phase (over) engineering the topology of light momentum to the background. The underlying physics is found to be the carefully designed background, such as a photonic crystal structure. The ping-pong ball is pulled from the oblique direction towards optical pulling, and just the way for a new class of advanced optical manipulation technique, with potential applications of drug delivery and cell sorting.

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Optical manipulations using various forms of optical force have been widely used in diverse disciplines, including biology [1–6], atom physics [7–10], quantum physics [11–12], and nanotechnology [13–19]. Recently, the counterintuitive optical pulling force (OPF) has attracted broad interest due to its intriguing physics underlying and potential applications [20–26]. An OPF can attract objects towards the source over a long range, rather than pushing them away, which provides a new frontier of optical manipulation beyond trapping and pushing. According to the conservation of linear momentum, the realization of OPF requires the increase of forward momentum of the incident light when scattered by the object, which is the common challenge to be solved in all the mechanisms of OPF generation.

Except for exotic media, such as optically active medium [27–30], chiral medium [31–33], PT symmetry objects [34,35], resonant objects [36,37] and negative refractive-index objects [38–40], various structured optical beams are always needed in order to get an OPF, including non-paraxial Bessel beams [33,32,41–43], orbital beams [44–46], interference fields of multiple beams [47–49], the whispering mode in a waveguide [50,51] or near an interface [52–54]. The common feature of these engineered fields is their off-axial alignment (relative to the average Poynting vector) of the incident momentum as shown in Fig. 1(a). When the beam is scattered by an object, the off-axial momentum may be redirected partially or even totally in the incident direction. In this case, the light momentum

projection on z -direction increases obviously due to the convex shape of the iso-frequency contour as illustrated in Fig. 1(b), and the OPF is generated according to the linear momentum conservation. Actually, this has been the standard mechanism to get an optical pulling force, which was first proposed by Maxwell in acoustics more than ten years ago [55,56]. Certainly, we can get an effective OPF using the thermal effect of light on an absorptive object or background [57,58], which is, however, not due to the transfer of light momentum, but due to the Brownian motion.

From Figs. 1(b) and 1(b') and the above discussion, we find that the key factor that prevents the generation of an OPF by a plane wave is the convex-shaped iso-frequency contour which forbids the forward light momentum increase of the plane wave. This is also the reason of structured beams being indispensable in most previous OPF investigations. In order to surpass this restriction, we propose to engineer the normal feature of light momentum in an artificial background, where the iso-frequency contour is convex, rather than convex, as shown in Figs. 1(c) and 1(c').

The change of iso-frequency contours in k -space from one shape to another is called topological transition in the research of optical metamaterials [59–64]. Since k is proportional to the light momentum p by $p = \hbar k$ [65–67], the shape of iso-frequency contours can be traced to the topology of light momentum (TLM). In Fig. 1(c), when the incident plane wave is scattered by the object

(19) Breakup and Recovery of Topological Zero Modes in Finite Non-Hermitian Optical Lattices, Physical Review Letters

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Breakup and Recovery of Topological Zero Modes in Finite Non-Hermitian Optical Lattices

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The topological edge states (TESs) in 1D one-dimensional optical lattices are robust against field localization or scattering against the structural perturbations that would give rise to finite-order proximity eigenstates. However, the zero mode in a band of TES usually deviates from the exact zero-energy state in a finite Hermitian lattice due to the coupling between finite edge modes, which largely weakens the topological protection. Here, we first show that a breakdown of zero modes in finite Hermitian-Dirac optical lattices and then reveal that recovery by introducing non-Hermitian degeneracies with parity-time (PT) symmetry. We carry out experiments in 1D finite optical waveguide lattices, where PT symmetry was implemented with carefully controlled loss-gain transparency. The experimental results are fully compatible with the theoretical prediction. Our results show that the topological property of an open system can be tuned by non-Hermitian lattice engineering, which offers a route to enhance the topological protection in a finite system.

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Parity-time (PT) symmetric systems with non-Hermitian Hamiltonians have received increasing attention over the past decade [1–12]. On crossing the transition point [i.e., the exceptional point (EP)] [2,3], the spectrum is no longer purely real due to spontaneous PT symmetry breaking. Among various PT symmetric systems, PT photonics with intentionally modulated gain and loss has become an emerging topic [4,5], which has witnessed a series of novel optical phenomena and functionalities [6–12]. On the other hand, topological photonics has opened a great deal of interest with intriguing phenomena associated with the topological edge states (TESs) [13–15]. The dimensionless edge state with its eigenenergy pinned at the middle of a gapped band structure is referred to as the zero mode [16], which is expected to work in fault-tolerant waveguiding and quantum computation [17,18]. In fact, topological systems are usually of finite system size, which will weaken the topological protection. For example, zero modes may deviate from the exact zero energy in a finite Hermitian lattice due to the coupling effect [16], which inevitably leads to weakness of the intriguing topologically protected characteristics. It is necessary to strengthen these zero modes in a finite topological system. Recently, the topological properties of non-Hermitian systems have been studied [19–24] and TESs in non-Hermitian systems are demonstrated [25,26]. Non-Hermitian properties have been

proven to give rise to a lot of useful aspects in topological phenomena, such as selective enhancement of topological states [27], non-Hermitian-induced topological protection [28–29], topological insulator lasers [30–33], etc. They provide us with the foundation to investigate a variety of topological phenomena in open systems and a possible route for the recovery of zero modes by non-Hermiticity in finite systems.

In this Letter, we clearly demonstrate both the breakup and recovery of zero modes in a finite 1D topological optical system, where the recovery of zero modes was realized by introducing the loss to form a non-Hermitian topological configuration. We systematically investigated the interplay of PT phases and edge states in the non-Hermitian Su-Schrieffer-Heeger (SSH) model [34]. On crossing the EPs, the edge states are driven to undergo a transition from a trivial PT -symmetric edge state to a topological broken- PT one. The experimental results in silicon waveguide lattices with controlled losses are in good agreement with theoretical results. Our study reveals that the topological property of an open system can be tuned by non-Hermitian parameters and supplies some interesting physics and photonic applications in PT and topological systems.

We would like to start from a general 1D non-Hermitian SSH model defined by a finite coupled optical waveguide

(20) All-optical control of lead halide perovskite microlasers, Nature Communications



ARTICLE

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All-optical control of lead halide perovskite microlasers

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Lead halide perovskites based microlasers have recently shown their potential in nanophotonics. However, up to now, all of the perovskite microlasers are static and cannot be dynamically tuned in use. Herein, we demonstrate a robust mechanism to realize the all-optical control of perovskite microlasers. In lead halide perovskite microrods, deterministic mode switching takes place as the external excitation is increased: the onset of a new lasing mode switches off the initial one via a negative power slope, while the main laser characteristics are well kept. This mode switching is reversible with the excitation and has been explained via cross-gain saturation. The modal interaction induced mode switching does not rely on sophisticated cavity designs and is generic in a series of microlasers. The switching time is faster than 70 ps, extending perovskite microlasers to previously inaccessible areas, e.g., optical memory, flip-flop, and ultrafast switches etc.

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(21) Arbitrarily routed mode-division multiplexed photonic circuits for dense integration, Nature Communications



ARTICLE

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Arbitrarily routed mode-division multiplexed photonic circuits for dense integration

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On-chip integrated mode-division multiplexing (MDM) is an emerging technique for large-capacity data communications. In the past few years, while several configurations have been developed to realize on-chip MDM circuits, their practical applications are significantly hindered by the large footprint and inter-mode cross talk. Most importantly, the high-speed MDM signal transmission in an arbitrarily routed circuit is still absent. Herein, we demonstrate the MDM circuits based on digitized meta-structures which have extremely compact footprints. 112 Gbit/s signals encoded on each mode are arbitrarily routed through the circuits consisting of many sharp bends and compact crossings with a bit error rate under forward error correction limit. This will significantly improve the integration density and benefit various on-chip multimode optical systems.

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(22) Resonance-enhanced three-photon luminescence via lead halide perovskite metasurfaces for optical encoding, Nature Communications



ARTICLE

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Resonance-enhanced three-photon luminescence via lead halide perovskite metasurfaces for optical encoding

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Lead halide perovskites have emerged as promising materials for photovoltaic and optoelectronic devices. However, their exceptional nonlinear properties have not been fully exploited in nanophotonics yet. Herein we fabricate methyl ammonium lead tri-bromide perovskite metasurfaces and explore their internal nonlinear processes. While both of third-order harmonic generation and three-photon luminescence are generated, the latter one is less affected by the material loss and has been significantly enhanced by a factor of 60. The corresponding simulation reveals that the improvement is caused by the resonant enhancement of incident laser. Interestingly, such kind of resonance-enhanced three-photon luminescence holds true for metasurfaces with a small period number of 4, enabling promising applications of perovskite metasurface in high-resolution nonlinear color microprinting and optical encoding. The encoded information 'NANO' is visible only when the incident laser is on-resonance. The off-resonance pumping and the single-photon excitation just produce a uniform dark or photoluminescence background.

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(23) Room temperature electrofreezing of water yields a missing dense ice phase in the phase diagram, Nature Communications



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Room temperature electrofreezing of water yields a missing dense ice phase in the phase diagram

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Water can freeze into diverse ice polymorphs depending on the external conditions such as temperature (T) and pressure (P). Herein, molecular dynamics simulations show evidence of a high-density orthorhombic phase, termed ice χ , forming spontaneously from liquid water at room temperature under high-pressure and high external electric field. Using free-energy computations based on the Einstein molecule approach, we show that ice χ is an additional phase introduced to the state-of-the-art T - P phase diagram. The χ phase is the most stable structure in the high-pressure/low-temperature region, located between ice II and ice V, and next to ice V exhibiting two triple points at 6.06 kbar/131.23 K and 9.45 kbar/144.24 K, respectively. A possible explanation for the missing ice phase in the T - P phase diagram is that ice χ is a rare polarized ferroelectric phase, whose nucleation/growth occurs only under very high electric fields.

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(24) Quasiparticle interference and nonsymmorphic effect on a floating band surface state of ZrSiSe, Nature Communications



ARTICLE

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OPEN

Quasiparticle interference and nonsymmorphic effect on a floating band surface state of ZrSiSe

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Non-symmorphic crystals are generating great interest as they are commonly found in quantum materials, like iron-based superconductors, heavy-fermion compounds, and topological semimetals. A new type of surface state, a floating band, was recently discovered in the nodal-line semimetal ZrSiSe, but also exists in many non-symmorphic crystals. Little is known about its physical properties. Here, we employ scanning tunneling microscopy to measure the quasiparticle interference of the floating band state on ZrSiSe (001) surface and discover rotational symmetry breaking interference, healing effect and half-missing-type anomalous Umklapp scattering. Using simulation and theoretical analysis we establish that the phenomena are characteristic properties of a floating band surface state. Moreover, we uncover that the half-missing Umklapp process is derived from the glide mirror symmetry, thus identify a non-symmorphic effect on quasiparticle interferences. Our results may pave a way towards potential new applications of nanoelectronics.

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(25) Antiferromagnetic Order in Epitaxial FeSe Films on SrTiO₃,
Physical Review Letters

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Antiferromagnetic Order in Epitaxial FeSe Films on SrTiO₃

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Single monolayer FeSe film grown on $\sqrt{2}\times\sqrt{2}$ -RbTiO₃(STO) substrate shows the highest superconducting transition temperature ($T_c = 100$ K) and the broadest superconducting dome possible (width ΔT_c between 0 and T_c is only ~ 4 K). Although bulk FeSe crystals show similar magnetic order, the transition temperature (the so-called T_N) is ~ 100 K. Experimentally, presence of a lock of a d-orbital probe the magnetic state of FeSe/STO film reveals superconductivity is a superconducting state usually appears after the AFM order is suppressed [14]. The AFM spin fluctuations were proposed to play an important role in the pairing of iron pnictides [13, 5]. Among various iron pnictides, FeSe has the simplest crystal structure [8]. The T_N values of bulk FeSe is ~ 8 K and can increase to ~ 27 K under high pressure [7]. And the other iron pnictides bulk FeSe_{1-x}S_x also do not show AFM order [5] unless a certain pressure is applied [8–10].

The pairing mechanism of high-temperature superconductivity including cuprates and iron pnictides is one of the biggest challenges in modern physics. The antiferromagnetic (AFM) interaction has been long thought to be correlated with high-temperature superconductivity [3, 2] because the superconducting state usually appears after the AFM order is suppressed [14]. The AFM spin fluctuations were proposed to play an important role in the pairing of iron pnictides [13, 5]. Among various iron pnictides, FeSe has the simplest crystal structure [8]. The T_N values of bulk FeSe is ~ 8 K and can increase to ~ 27 K under high pressure [7]. And the other iron pnictides bulk FeSe_{1-x}S_x also do not show AFM order [5] unless a certain pressure is applied [8–10].

Surprisingly, a single monolayer (1-ML) FeSe film grown on a $\sqrt{2}\times\sqrt{2}$ -RbTiO₃(STO) [FeSe] will exhibit $\sqrt{2}\times\sqrt{2}$ -doped SrTiO₃(STO) [11] structure after electron doping. Although the magnetic pressure shows a large stepwise increase up to 20 meV [11] that increases up to ~ 100 K [12, 13]. Our magnetic studies below ~ 65 K have also been reported [14]. Recently, the π resonant magnetic measurements showed the T_N values of 1-ML FeSe/STO films can be as high as 100 K [15]. The mechanism of such a high T_N value is still an open question. Calculations have shown that the electron-doping coupling is significantly enhanced [16] due to the interfacial effect and therefore enhances the value of T_N in the $\sqrt{2}\times\sqrt{2}$ case [17], but the initial pairing mechanism is still unclear. First-principles calculations have shown that the FeSe/STO interface could enhance the AFM

interaction, which gives potential large spin fluctuations under nearly electron doping [18]–[24] and is enhanced by the localization of the electronic doping and the presence of another possible mechanism for the superconductivity [19]. Density functional theory (DFT) calculations have suggested that the magnetic ground state of 1-ML FeSe/STO film is AFM [18, 20, 21]. A group also claimed that 1-ML FeSe/STO could be an AFM order in terms topological superconductivity [22]. Therefore, it is very interesting to study the magnetic ground state of the 1-ML film before electron doping. Experimentally, the magnetic state of the FeSe/STO films is barely known. Previous angle-resolved photoemission spectroscopy (ARPES) measurements showed indirect signatures of the spin density waves [12], but they are indistinguishable from the band structure [25–28]. To determine the magnetic state, regular techniques such as neutron scattering, x-ray spin Hall effect and the Aharonov effect have failed so far to identify the magnetic film. In this Letter, we present direct evidence indicating that the magnetic ground state of the quasi-1-ML FeSe/STO film is AFM by using magnetic evanescent bias effect (MEBE) [29] measurements.

FeSe/STO films were grown following previous reports [11, 15]. Films before annealing are called “as-grown” films. The as-grown films were pre-annealed at ~ 300 °C for 3–6 h in air to make them superconducting. Before the films were transferred to another chamber to grow on Fe_{0.9}Se_{0.1} layer, a 30-nm-thick Sic protective layer [30] grows. The polycrystalline Fe_{0.9}Se_{0.1} film was grown by

(26) Self-induced backaction optical pulling force, Physical Review Letters

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Self-Induced Backaction Optical Pulling Force

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We achieve long-range and continuous optical pulling in a periodic photonic crystal background, which supports a unique Bloch mode with the subwavelength effect. Most remarkably, the pulling force reported here is mainly contributed by the intensity gradient force originating from the self-induced backaction of the object to the self-collimation mode. This force is sharply distinguished from the widely held conception of optical pulling forces based on the scattering force. Also, the pulling force is insensitive to the angle of incidence and can pull multiple objects simultaneously.

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Since the pioneering work of Ashkin (1-3), the concepts and techniques of matter trapping and manipulation by an optical force have experienced impressive progress in the past decades (4-6). Optical tweezers have become an indispensable tool in various disciplines, including biology (6-10), optics (11-13), and condensed matter (14) as well as quantum physics (15). In traditional optical trapping, however, it is extremely challenging to scale down below 100 nm due to both the diffraction-limited focusing of light and the r^{-4} decaying law of the optical force with the radius r of small particles. Under these circumstances, by setting the trapping operation inside a nanoscale, the self-induced backaction optical trapping was proposed in manuscripts (16-22), which can trap objects of size of nanometers (or even smaller, thanks to the active feedback of the trapping object to the trapping system).

Despite its power in nanometer trapping, the self-induced backaction method has not been explored in long-range optical pulling (or in optical tractor beams), because it influences the intensity gradient component of the coupling force in a nanoscale, while a tractor beam relies on the scattering component in a translation-invariant structure. A tractor beam, initially proposed in scenarios by Mariani *et al.* (23-25), is an intensity gradientless beam along the propagation direction, which can transport objects towards the source over a long distance rather than push them away. The counterintuitive optical pulling (26-29) occurs increasing

medium due to its sophisticated underlying physics and potential applications (26-48).

In the extensive studies up to now, except for the phonononic forces (49,50), all optical pulling forces have been regarded as optical scattering forces, like in the schemes of structured light beams (26-33,40), exotic media (34-42) (51), inhomogeneous background media (52,53), and guiding mode scattering (34,44,54).

In this Letter, we propose a novel approach based on the self-induced backaction-generated gradient force that achieves long-range optical pulling in a periodic photonic structure, namely, in a photonic crystal (PC) (55,56). Interestingly, the pulling force arising in the PC is not predominantly contributed by the scattering force as reported before but reduced by the intensity gradient force generated by the dynamic interaction of the object with the self-collimation (SC) mode of the PC (57-59). The SC mode is a unique kind of Bloch mode with a finite numerical size that can propagate infinitely long without diffraction. Hence, our results are sharply distinct to the widely held conception that the optical pulling force should be a scattering force, and shed new insightful insights concerning the optical force and momentum physics and technologies.

To illustrate the underlying physics, we address a two-dimensional PC structure with a square lattice as illustrated in Fig. 1. Figure 1(a) shows the outline of the entire configuration. The SC mode is excited by a Gaussian beam at an angle of incidence θ_0 and propagates along the x or

(27) Atomic-level insight into super-efficient electrocatalytic oxygen evolution on iron and vanadium co-doped nickel (oxy)hydroxide, Nature Communications



ARTICLE

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Atomic-level insight into super-efficient electrocatalytic oxygen evolution on iron and vanadium co-doped nickel (oxy)hydroxide

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It is of great importance to understand the origin of high oxygen-evolving activity of state-of-the-art multimetal oxides/(oxy)hydroxides at atomic level. Herein we report an evident improvement of oxygen evolution reaction activity via incorporating iron and vanadium into nickel hydroxide lattices. X-ray photoelectron/absorption spectroscopies reveal the synergistic interaction between iron/vanadium dopants and nickel in the host matrix, which subtly modulates local coordination environments and electronic structures of the iron/vanadium/nickel cations. Further, in-situ X-ray absorption spectroscopic analyses manifest contraction of metal-oxygen bond lengths in the activated catalyst, with a short vanadium-oxygen bond distance. Density functional theory calculations indicate that the vanadium site of the iron/vanadium co-doped nickel (oxy)hydroxide gives near-optimal binding energies of oxygen evolution reaction intermediates and has lower overpotential compared with nickel and iron sites. These findings suggest that the doped vanadium with distorted geometric and disturbed electronic structures makes crucial contribution to high activity of the trimetallic catalyst.

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(29) A new phase diagram of water under negative pressure: The rise of the lowest-density clathrate s-III, Science Advances

RESEARCH ARTICLE

WATER CHEMISTRY

A new phase diagram of water under negative pressure: The rise of the lowest-density clathrate s-III

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Ice and ice clathrate are not only omnipresent across polar regions of Earth as outer terrestrial moons but also ubiquitous in the solar system such as on comets, asteroids, and moons of the giant planets. Depending on the surrounding environment (temperature and pressure), ice/water exhibits an exceptionally rich and complicated phase diagram with 17 known crystalline polymorphs. Water molecules also form clathrate compounds with inclusion of guest molecules, such as cubic structure I (s-I), cubic structure II (s-II), hexagonal structure I (h-I), tetragonal structure I (t-I), and tetragonal structure II (t-II). Recently, guest-free clathrate structure II (s-III), also known as ice XVI located in the negative pressure region of the phase diagram of water, is synthesized in the laboratory and motivates scientists to reexamine other ice clathrates with low density. Using extensive Monte Carlo packing algorithms and dispersion corrected density functional theory optimizations, we predict a crystalline clathrate of cubic structure II (s-III) composed of two large pentadecahedral cages (10⁵Å³) and five small dodecahedral cages (3⁵Å³) per unit cell which is dynamically stable by itself and can be fully stabilized by encapsulating an appropriate guest/molecule in the large cavity. A new phase diagram of water ice with TIP4P/2005 four-point transferable intermolecular potential/2005 model potential is constructed by considering a variety of candidate phases. The guest-free s-III clathrate with ultralow density overrules s-I and s-II phases and emerges as the most stable ice polymorph in the pressure region below ~500 bar at 0 K and below ~1011 bar at 100 K.

INTRODUCTION

Water molecules are the third most abundant molecular species in the universe (1). The six face forms of water ice condensed into the most common molecular solid on Earth and can be also detected in the giant planetary interior. The existence of water and/or ice is imperative to the diversity of nature and possible presence of life. Because of the flexible hydrogen bonds, water ice exhibits an exceptionally rich and complicated phase diagram (2–10). Under different conditions of pressure (P) and temperature (T), there are 17 experimentally established crystalline phases of ice on Earth (3, 11). Among them, ice XI has the lowest mass density of 0.938 g/cm³ and forms a proton-ordered phase at very low temperature and ambient pressure (12). At ambient temperature (298 K), ice XI transforms into a proton-disordered phase, that is, the well-known ice II (13). In addition, many hypothetical ice phases are also predicted and well experimentally confirmed, for example, several ice s-I and s-II with low density (14), ice III as a precursor for ice nucleation (15), silica-like ice polymorphs (16), and partially amorphous (low-density) under extremely high pressure (15, 16, 17) (Fig. 1).

Apart from pure water (ice, water and vapor), gas hydrates (ices) are also stable. At H₂, CO₂, CH₄, C₂H₆, adamantane and n-butylcyclohexane can form a class of atomically ordered inclusion compounds, namely,

the gas hydrate or ice clathrate (18–19). In the ice clathrates, water molecules form loosely hydrogen-bonded frameworks constructed by pentahedral cages whose main cavities do not accommodate or partially occupied by guest molecules. So far, at least five types of ice clathrate structures have been experimentally synthesized (20, 21, 19) or theoretically proposed (22): cubic structure I (s-I) with 2 × 10² cages and 6 × 10² cages per unit cell, cubic structure II (s-II) with 16 × 10² cages and 8 × 10² cages, hexagonal structure I (h-I) with 45 × 10² cages, 2 × 10² cages, and 4 × 10² cages, tetragonal structure I (t-I) with 2 × 10² cages, and tetragonal structure II (t-II) with 6 × 10² cages, 4 × 10² cages, and 4 × 10² cages per unit cell, respectively. Usually, the clathrate hydrates are considered as hypothetical phases of water (23, 24) and figure to be stabilized by favorable van der Waals (vdW) bonds between guest molecules and host water cages (25–27, 28). However, Flory and coworkers (11) recently synthesized guest-free s-III clathrate (also named ice XVI) by leaching Xe atoms from the s-II clathrate. The guest-free s-III clathrate is mechanically stable at atmospheric pressure up to 100 K (29) and exhibits a very different expansion with good mechanical stability and large lattice constant than the filled clathrate. The experimental confirmation of structure s-III lattice still only confirms that water molecules themselves can form three crystalline phases on a density lower than ice XI but also motivates us to further explore other possible low-density polymorphs of water ice that are stable at low temperature and negative pressure.

To obtain the comprehensive knowledge of phase diagram of water under positive pressure, the negative pressure region of the diagram is much less explored (3), 22, 28, 29). Some efforts have been devoted to determining the driving mechanism for the fact that the dissolved liquid water can sustain before the event of nucleation occurs, namely the cavitation pressure (24–26). General prediction

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(30) Quasi-freestanding epitaxial silicene on Ag(111) by oxygen intercalation, Science Advances

RESEARCH ARTICLE

MATERIALS ENGINEERING

Quasi-freestanding epitaxial silicene on Ag(111) by oxygen intercalation

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Silicene is a monolayer allotrope of silicon atoms arranged in a honeycomb structure with massless Dirac fermion characteristics, similar to graphene. It merits development of silicon-based multifunctional nanoelectronic and spintronic devices operated at room temperature because of strong spin-orbit coupling. Nevertheless, silicene could only be epitaxially grown on conductive substrates. The strong silicene-substrate interaction may depress its superior electronic properties. We report a quasi-freestanding silicene layer that has been successfully obtained through oxidation of bilayer silicene on the Ag(111) surface. The oxygen atoms intercalate into the underlayer of silicene, resulting in isolation of the top layer of silicene from the substrate. In consequence, the top layer of silicene exhibits the signature of a 1 × 1 honeycomb lattice and hosts massless Dirac fermions because of negligible interaction with the substrate. Furthermore, the oxidized silicon buffer layer is expected to serve as an ideal silicene layer for electrocatalysis in electronic devices. These findings are relevant for the future design and application of silicon-based nanoelectronic and spintronic devices.

INTRODUCTION

Silicene can be grown in a variety of ways, ranging from theoretical calculations, including *ab initio* calculations, the quantum spin Hall effect, and possible superconductivity (1–20) to experimental low-temperature silicene growth using group-IV hydrides on Si(111) and Si(100) as appropriate substrate materials. There are also other approaches, such as the use of the monolayer silicene structure to be formed on several conductive substrates by epitaxial growth (21–23). However, the strong silicene-substrate interaction may markedly depress the superior electronic properties of the two-dimensional (2D) material (12, 24). For example, the hybridization between sp₂ and sp₃ orbitals results in a surface modification and depresses the Dirac fermion characteristics in an epitaxial silicene layer on an Ag(111) surface (12). Moreover, the conductive substrate makes it difficult to modulate the band level of silicene by electrocatalysis, and the bandy integration of silicene into microelectronic devices. Hence, how to overcome the substrate-induced band structure and silicene characteristics of epitaxial silicene has become a crucial issue for the development of silicene devices. Although some theoretical proposals have been put forward for heterostructure-based silicene on inert substrate (25–26), so far, there has hardly been any experimental success.

In bilayer silicene on Ag(111), the top sp²-sp² layer (with respect to (1 × 1) lattice) is considered to be hybridized to the sp²-sp³ layer (with respect to (1 × 1) Ag(111)) (27). Therefore, the lower sp²-sp³ layer could be regarded as a buffer layer. Previous experimental evidence suggests that the sp²-sp³ layer has a large hybridization state and a different chemical nature (12), whereas the sp²-sp² layer is highly sensitive to ambient gases, especially O₂ (19–21). Recently, oxidation

of the graphene-substrate interface has been achieved by introducing epitaxial graphene to metal (Pt, Pd, Au) contacts by using hydrogen or fluorine intercalation in the buffer layer (22, 28). Motivated by such a successful strategy, the interaction between the sp²-sp² silicene layer and the inert substrate is expected to be effectively reduced by an appropriate intercalation into the buffer sp²-sp³ silicene layer, which in turn may allow the top layer to recover the intrinsic properties of silicene.

Here, we report the oxidation of bilayer silicene on Ag(111) and found that the oxygen molecules intercalate into the sp²-sp³ buffer layer of silicene. As a result, the epitaxial silicene on the surface of the 1 × 1 lattice structure of “freestanding” silicene and recovers a clear Dirac fermion characteristic with low electron doping, as revealed by scanning tunneling microscopy (STM), x-ray photoelectron spectroscopy (XPS), and angle-resolved photoemission spectroscopy (ARPES) measurements. By comparing these results with theoretical calculations, we demonstrate that the top layer of silicene exhibits quasi-freestanding characteristics of isolated interaction with the substrate. Our study contributes a novel and simple way to obtain quasi-freestanding silicene on a substrate. In addition, the silicene oxide buffer layer may be used as the silicene layer for possible three-dimensional (3D) effect transistors (29,30) or the metal nanowire.

RESULTS

Figure 1A shows an STM image of periodic (1 × 1) silicene supported by the sp²-sp³ buffer layer on Ag(111). In the brown sample, the epitaxial silicene buffer layer shows a clear atomic structure that $a_1 = a_2 = a_3 = 0.354 \text{ nm}$ (31) and $\alpha_1 = \alpha_2 = \alpha_3 = 120^\circ$ (32). The top layer of silicene exhibits a (1 × 1) lattice with a lattice constant $a = 0.35 \text{ nm}$, which is approximately 1% from that of the 1 × 1 silicene epitaxial on a Si(111) (33). Figure 1B is an STM image of a single piece of silicene island with a top sp²-sp² layer collected after the sample was exposed to oxygen with a dose of 600 Langmuir (L). The height area of the island is higher than the substrate by about 0.3 Å, as shown in Fig. 1C. The high-resolution STM image (Fig. 1D)

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(31) Majorana Zero Mode Detected with Spin Selective Andreev Reflection in the Vortex of a Topological Superconductor, Physical Review Letters

Majorana Zero Mode Detected with Spin Selective Andreev Reflection in the Vortex of a Topological Superconductor

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Recently, many topological materials are proposed to possess parallel selective Andreev reflection (AR) in general quantum physics, which can be used to detect the Majorana zero mode (MZM) through spin-polarized scanning tunneling spectroscopy. In this study, we have applied the probe AR of MZMs in a topological superconductor of Bi₂Te₃/Sb₂Te₃ heterostructure. The spin bias probe of the tunneling differential conductance at the vortex center is observed to be spin-polarized when the tip polarization and the external magnetic field are parallel rather than antiparallel to each other. This spin-dependent tunneling effect provides the evidence of MZM and reveals its magnetic entropy in addition to the zero energy modes. The work will stimulate MZM research in various novel physical properties and, hence, is a very useful experimental guide of their synthesis and application in quantum computing.

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The Majorana fermion (MF) is a special type of fermion whose antiparticle is itself [1]. The MF was initially proposed in elementary particle physics, and the recent effort in searching such a genuine particle focuses on the neutrinos double beta decay experiment [2]. The Majorana zero mode (MZM), which fits to the Majorana equation that describes MFs and emerges as a novel solution in some condensed matter systems (MNs) when non-Abelian statistics may be used in robust building blocks by quantum computers [3–5]. Chiral p-wave superconductors [6] and the $\nu = 5/2$ fractional quantum Hall system [8] are possible candidates for MZMs [9]. Fu and Kane proposed the existence of MZMs at the interface of a topological insulator (TI) and an s-wave superconductor (SC) [7]. In recent years, a number of proposals have been proposed to detect the MZM, including the experiments in one-dimensional (1D) Kitaev semiconductor wire [9–11] and 1D Tsu–Kane chains in Pb [12]. Topological superconductivity may be induced on surfaces of 3D TIs such as Bi₂Se₃ [6], Te₂, and Bi₂Te₃ via the proximity effect, and the localized MZM at the vortex core has been studied by using scanning tunneling microscopy or spectroscopy (STM/STS) [13–15]. However, there are quasiparticle states inside the vortex core, which may mask signals cannot be distinguished from the zero modes within the present STM energy resolution. This introduces great difficulty in detection of the MZM by using normal STM/STS with a nonpolarized tip. In our recently proposed a novel

property (sMZM in a 3D system [16]), where the spin-selective AR is spin selective. They showed that MZM induces selective spin-polarized AR, in which incoming electrons with certain spin polarization in the lead are reflected as counter-propagating holes with the same spin [1]. The proposed spin-selective AR (SSAR) of the MZM in 3D systems can be generalized to a 2D topological superconductor (TSC). In the latter case, the spin of the MZM in the vortex core has a spatial distribution and is not conserved [17]; however, we show that the spin wave function at the center of the vortex core is $\pm 1/2$ of the Majorana mode is fully polarized along the diagonal magnetic field with \pm the lateral distance of the tip from the vortex center (see Fig. 3(a)). Therefore, the AR at $\nu = 1/2$ is expected to be spin selective, and can be probed in spin-polarized STS/STM/STS [18–20].

A simple advantage to probe SSAR in the TSC is that a small external field is sufficient to induce random and linear the spin-polarized MZM in the vortex core. On the other hand, the spin polarization of surface states and the vortex core states in the states is well conserved, which we can see in Fig. 3(b) below. In order to observe the SSAR, we apply a sufficiently large magnetic field to make the vortex topological in the *xy* plane, in which case the vortex core states are already spin-polarized. It will then be difficult to attribute spin-dependent spin-linear conductance to the MZM since both MZM (if present) and surface states and gap states would be spin-polarized.

(32) Epitaxial growth of two-dimensional stanene, Nature Materials

ARTICLES

INTEGRATED AI, MACHINE LEARNING AND DATA SCIENCE

nature
materials

Epitaxial growth of two-dimensional stanene

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Following the first experimental realization of graphene, other ultrathin materials with unprecedented electronic properties have been explored, with particular attention given to the heavy group-IV elements Sn, Ge and Sr. Two-dimensional buckled Sn-based stanene has been recently realized by molecular beam epitaxy growth, whereas Ge-based germanene was obtained by molecular beam epitaxy and mechanical exfoliation. However, the synthesis of Sn-based stanene has proved challenging so far. Here, we report the successful fabrication of 2D stanene by molecular beam epitaxy, confirmed by atomic and electronic characterization using scanning tunnelling microscopy and angle-resolved photoemission spectroscopy, in combination with first-principles calculations. The synthesis of stanene and its derivatives will stimulate further experimental investigation of their theoretically predicted properties, such as a 2D topological insulating behaviour with a very large bandgap, and the capability to support enhanced thermoelectric performance, topological superconductivity and the near-room-temperature quantum anomalous Hall effect.

Recent first-principles calculations^{1–5} have predicted that a layered structure of graphene-like stanene could be synthesized as a freestanding or heterostructure material. In fact, a bilayer structure of stanene has been experimentally synthesized, showing a bandgap of 0.2 eV, similar to that of germanene. However, the synthesis of stanene has proved challenging so far. Here, we report the successful fabrication of 2D stanene by molecular beam epitaxy, confirmed by atomic and electronic characterization using scanning tunnelling microscopy and angle-resolved photoemission spectroscopy, in combination with first-principles calculations. The synthesis of stanene and its derivatives will stimulate further experimental investigation of their theoretically predicted properties, such as a 2D topological insulating behaviour with a very large bandgap, and the capability to support enhanced thermoelectric performance, topological superconductivity and the near-room-temperature quantum anomalous Hall effect.

Stanene is a 2D material with a diamond-like structure¹ that is composed of a monolayer layer of a $\sqrt{3}\times\sqrt{3}$ triangular lattice of Sn atoms. It has a bandgap of 0.2 eV, similar to that of germanene². The buckling of stanene formed by the large difference between the top and bottom Sn atoms is typically around 0.25 eV depending on the thickness of the layer^{1,6,7}. The synthesis of stanene has proved challenging so far. Here, we report the successful fabrication of 2D stanene by molecular beam epitaxy, confirmed by atomic and electronic characterization using scanning tunnelling microscopy and angle-resolved photoemission spectroscopy, in combination with first-principles calculations. The synthesis of stanene and its derivatives will stimulate further experimental investigation of their theoretically predicted properties, such as a 2D topological insulating behaviour with a very large bandgap, and the capability to support enhanced thermoelectric performance, topological superconductivity and the near-room-temperature quantum anomalous Hall effect.

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(33) The structural origin of the hard-sphere glass transition in granular packing, Nature Communications



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The structural origin of the hard-sphere glass transition in granular packing

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Glass transition is accompanied by a rapid growth of the structural relaxation time and a concomitant decrease of configurational entropy. It remains unclear whether the transition has a thermodynamic origin, and whether the dynamic arrest is associated with the growth of a certain static order. Using granular packing as a model hard-sphere glass, we show the glass transition as a thermodynamic phase transition with a 'hidden' polytetrahedral order. This polytetrahedral order is spatially correlated with the slow dynamics. It is geometrically frustrated and has a peculiar fractal dimension. Additionally, as the packing fraction increases, its growth follows an entropy-driven nucleation process, similar to that of the random first-order transition theory. Our study essentially identifies a long-sought-after structural glass order in hard-sphere glasses.

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(34) Superconductivity above 100 K in single-layer FeSe films on doped SrTiO₃, Nature Materials

Superconductivity above 100 K in single-layer FeSe films on doped SrTiO₃

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Recent experiments on FeSe films grown on SrTiO₃ (STO) suggest that interface effects can be used as a means to reach superconducting critical temperatures (T_c) of up to 80 K (ref. 1). This is nearly ten times the T_c of bulk FeSe and higher than the record value of 56 K for known bulk Fe-based superconductors². Together with recent studies of superconductivity of oxide heterostructure interfaces^{3,4}, these results reinforce the long-standing idea that electron pairing at interfaces between two different materials can be tailored to achieve high-temperature superconductivity^{5,6}. Subsequent angle-resolved photoemission spectroscopy measurements of the FeSe/STO system revealed an electronic structure distinct from bulk FeSe (refs 13,14), with an energy gap vanishing at around 65 K. However, *in situ* electrical transport measurements¹⁵ have so far detected zero resistance—the key experimental signature of superconductivity—only below 20 K. Here, we report the observation of superconductivity with T_c above 100 K in the FeSe/STO system by means of *in situ* four-point probe electrical transport measurements. This finding confirms FeSe/STO as an ideal material for studying high- T_c superconductivity.

The search for superconductors without T_c above the liquid nitrogen temperature (77 K) led to the discovery of high- T_c cuprates with a T_c as high as 135 K, only three decades ago⁷. Since then, the value of T_c is now 30 K in the first Fe-based superconductor LaFeAsO (ref. 8), subsequent work on a series of Fe-based superconductors showed that the highest T_c under ambient pressure is 56 K (found in FeVAsO (ref. 9) or La_{1-x}FeAsO_{1-x} with a T_c of 57 K, here we limited to the cuprate-based materials) and has been considered to represent the superconductivity in Fe-based systems. However, the FeSe/STO system we reported in this paper (see Supplementary Information) exhibits an energy gap of FeSe/STO was found to vanish at approximately 65 K (ref. 15), which is 20 K lower than the superconducting critical temperature (T_c) of 100 K (see Supplementary Information). This seems together with the earlier work on superconducting oxide interfaces^{3,4}, demonstrating that the interface between two different materials provides not only a rich system to realize superconductivity^{5,6} but also provides a pathway to high- T_c superconductivity^{5,6}.

Direct-angle-angle-resolved photoemission spectroscopy experiments on the FeSe/STO system revealed a different electronic structure from that of bulk FeSe and a possible existence of

superconductivity around 100 K (ref. 15) (Fig. 1). The temperature dependence of photoemission spectroscopy (PES) intensity at 100 K and an insulating behaviour around a low-energy gap of 15 meV and an onset T_c of 100 K (ref. 15). Evidently, the addition of interface layer suppresses superconductivity in single-layer FeSe. In this work, we report the first successful preparation of single-layer films of FeSe grown on TiO_2 -doped STO, subsequent *in situ* electrical-transport measurements. We found that superconductivity could be obtained even at a temperature as high as 100 K.

Single-layer films of FeSe were grown on TiO_2 -doped STO substrate by the same method as reported previously¹⁶, involving a combination of molecular beam epitaxy (MBE) and pulsed laser deposition (PLD) capabilities. The growth process was monitored by *in situ* high-energy electron diffraction (HEED) (Fig. 2a), which allowed the *in situ* control of film growth (width of surface and thickness) (see Supplementary Information). The growth process of the films was confirmed by XPS imaging of both layers and atomic surface of film on the TiO_2 side (see Supplementary Information).

The PLD technique provides a flexibility of superconducting single-layer FeSe films grown on STO_{1-x}Ti_xO₃ substrate by changing x from 0.05 to 0.15. Fe-based superconducting heterostructures have been proposed with the structure $\text{FeSe}/\text{TiO}_2/\text{SrTiO}_3/\text{LaAlO}_3/\text{TiO}_2$ and $\text{LaFeO}_3/\text{TiO}_2/\text{SrTiO}_3/\text{LaAlO}_3/\text{TiO}_2/\text{FeSe}$ (ref. 17) and $\text{FeSe}/\text{TiO}_2/\text{SrTiO}_3/\text{LaAlO}_3/\text{TiO}_2/\text{FeSe}$ (ref. 18) and so on. The applicability of the PLD technique to superconducting films is well demonstrated¹⁹. Furthermore, due to depositing films grown on an insulating substrate, the interface region is covered by the carrier between the film and the substrate. Therefore, the PLD technique is a powerful tool to provide free interface superconductivity in these devices and to fabricate an artificial superconductor with a high T_c and a rich interface region.

Before IVP measurements, the surface heterostructure of the samples were examined, as shown in Supplementary Fig. 1 and all of them were found to have the capability of forming superconducting state at low pressure (see Supplementary Information). The *in situ* four-point probe electrical transport measurements (Fig. 3) showed that the T_c increased from 65 K (ref. 15) to 100 K (see Supplementary Information). The T_c was found to be 100 K, which indicates the observed results are not from disorder due to the real situation (see Supplementary Information). Figure 3b shows that the T_c values were obtained for two different orientations (Supplementary Fig. 4).

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(35) Demonstration of Nonlinear-Energy-Spread Compensation in Relativistic Electron Bunches with Corrugated Structures, Physical Review Letters

Demonstration of Nonlinear-Energy-Spread Compensation in Relativistic Electron Bunches with Corrugated Structures

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High quality electron beams with flat distributions in both energy and current are critical for many accelerator-based scientific facilities such as free-electron lasers and MeV ultrafast electron diffraction and microscopy. In this Letter, we report on using corrugated structures to compensate for the beam nonlinear energy chirp (induced by the existence of the radio-frequency field), leading to a significant reduction in beam energy spread. By using a pair of corrugated structures with orthogonal orientations, we show that the quadrupole wakefields, which otherwise, in some beam structures, can be effectively cancelled. This work also extends the applications of corrugated structures in the low beam charge (a few pC) and low beam energy (a few MeV) regime and may have a strong impact on many accelerator-based facilities.

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"Flat-flat" electron beams, i.e., flat in both beam energy and current [1], are critical for achieving high performance in many accelerator-based facilities such as seeded free-electron lasers (FELs) [2–6] and MeV ultrafast electron microscopes (UEMs) [7–8]. By either seeding the relativistic electron beams through undulators to produce coherent x-ray & rays or directly using the electrons as probes, these facilities are enabling new opportunities in an extraordinarily wide array of sciences. With the beam produced in a photocathode radio-frequency (rf) gun and accelerated by high energy rf structures (the electron beam longitudinal phase space, however, typically consists of nonlinear chirp (energy-time correlations) from the varying rf phase along the bunch. In seeded FELs, this nonlinear energy chirp deteriorates the performance of bunch compressors and broadens the FEL spectrum. In UEMs, the nonlinear energy chirp increases the beam global energy spread and, thus, leads to serious chromatic aberrations that reduce the microscope resolution. A similar detrimental effect also holds for ultrafast electron diffraction facilities where the nonlinear energy chirp sets the lower limit of the electron bunch length that can be achieved with velocity bunching [9].

The standard method for compensating for the nonlinear energy chirp is to use a harmonic cavity (see, for example, [10]). The idea is rather simple. Consider a beam passing through two rf structures with wave numbers at $k_{1,z}$, peak energy gain at $Z_{1,z}$ and phase at $\phi_{1,z}$, the energy of a

particle at longitudinal position z with respect to the reference particle can be written as

$$E(z) = E_1 \cos(k_1 z + \phi_1) + E_2 \cos(k_2 z + \phi_2) + \dots \quad (1)$$

Assuming the first rf structure is mainly used for beam acceleration and the second structure for cancellation of the nonlinear chirp, then the phase of the main acceleration structure should be correspondingly set at $\phi_1 = 0$ to provide maximal energy gain and the phase of the harmonic cavity should be set at the decelerating phase ($\phi_2 = \pi$). Under this condition, the energy chirp is cancelled up to the second order if $E_2 k_2^2 = E_1 k_1^2$, i.e., the required voltage for the harmonic cavity scales as $1/n^2$, where $n = k_2/k_1$ is the harmonic number.

Cancellation of the nonlinear chirp with active rf structures not only requires dedicated (expensive) rf modules, the relative phase of the rf structures needs to be accurately controlled as well. Another undesired outcome is that the beam energy will also be reduced by $1/n^2$ (e.g., if a C-band structure with frequency at 5712 MHz is used to cancel the nonlinear chirp from an S-band structure with frequency at 2856 MHz, the beam energy will drop by $1/4$). While reduction of beam energy by a fraction for FELs is typically not a big concern, it may result in stronger space-charge effect in UEMs since the beam energy is only a few MeV [7,8]. Furthermore, it may add considerable complexity and cost to such compact facilities.

(36) Multichromatic Narrow-Energy-Spread Electron Bunches from Laser-Wakefield Acceleration with Dual-Color Lasers, Physical Review Letters

Multichromatic Narrow-Energy-Spread Electron Bunches from Laser-Wakefield Acceleration with Dual-Color Lasers

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A method based on laser wakefield acceleration with controlled ionization injection (triggered by another frequency triple laser) is proposed, which can produce electron bunches with low energy spread. As the color pulses copropagate in the background plasma, the peak amplitude of the combined laser field is modulated in time and space during the laser propagation due to the plasma dispersion. Ionization injection occurs when the peak amplitude exceeds a certain threshold. The threshold is exceeded for limited duration periodically at different propagation distances, leading to multiple ionization injections and separated electron bunches. The method is demonstrated through multidimensional particle-in-cell simulations. Such electron bunches may be used to generate multichromatic x-ray sources for a variety of applications.

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Versatile x-ray sources with tunable brightness, spectrum range, pulse duration, and temporal-spatial coherence could be used for scientific instruments as well as medical and industrial applications. Continuous efforts are being made to push x-ray sources towards new limits such as coherent beams over 10 keV [1,2], attosecond pulses [3], two-color hard x-ray lasers [4], etc. However, currently most of these sources are based on conventional accelerators. More compact and low cost x-ray devices with comparable quality are highly desired. Recently, laser wakefield acceleration (LWFA) offers the possibility for a new generation of compact particle accelerators [5]. Single or multiple bunches can be generated for different applications [6–8]. LWFA-based compact and low cost α - and γ -ray sources also have attracted much interest [9,10].

In spite of the significant progress made in LWFA research in the past decade [11–16], it is widely recognized that the beam quality and stability still need to be improved considerably before its wide applications. The injection process is a key issue of high quality beam production. There are a few interesting schemes proposed. A cold optical injection scheme is proposed to produce ultralow energy spread beams as shown by two-dimensional (2D) simulations [17]. Besides, ionization-induced injection in mixed gases targets has been proposed and demonstrated as an attractive scheme [20–26] due to the relatively lower ionization threshold such that it is effectively pre-ionized and acts as background plasma, while at least another has inner shells with higher ionization thresholds (such as nitrogen and oxygen). The laser releases these inner shell

electrons at a location within the wake such that they can be easily trapped and accelerated. Generally, the final electron beam energy spread is related to the effective injection length [31–33], which often leads to energy spread much larger than 1% unless some technique, such as laser self-focusing, is used [34].

In this Letter, we propose a new electron injection scheme to produce ultralow energy spread and single-peaked multiple electron bunches with equally spaced energy peaks. We use a bichromatic laser to trigger sequential ionization injections. In our one-dimensional (1D) and multidimensional particle-in-cell (PIC) simulations using the code *osiris* [35], low energy spread beams are generated because the effective injection length is suppressed to a few hundred micrometers. A multipeak energy spectrum can also be observed if a specific condition is satisfied.

This scheme is illustrated in Fig. 1. A main pulse with the fundamental frequency is responsible for driving an accelerating wakefield in the blowout regime [36–38]. A second copropagating harmonic pulse with a smaller amplitude modulates the peak laser field strength and acts as a trigger of the high-Z gas K -shell ionization. Because of the laser dispersion in the plasma, the phase speed of the two frequency components are different. By tuning the amplitude of the two components, one can limit the K -shell ionization only occurring when the peaks of the two lasers overlap.

To understand this process, we first study the propagation of a bichromatic laser in plasmas. Consider two plane waves: $\mathbf{E} = \{1, 2\}$ with the normalized vector potential

$$a_i(z, t) = a_{0i} \sin(\omega_i t - k_i z + \phi_i), \quad (1)$$

(37) Experimental Detection of a Majorana Mode in the core of a Magnetic Vortex inside a Topological Insulator-Superconductor Bi₂Te₃/NbSe₂ Heterostructure, Physical Review Letters

Experimental Detection of a Majorana Mode in the core of a Magnetic Vortex inside a Topological Insulator-Superconductor Bi₂Te₃/NbSe₂ Heterostructure

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Majorana fermions have been intensively studied in recent years for their importance to both fundamental science and potential applications in topological quantum computing. They are predicted to exist in a certain case of superconducting topological insulators. However, it is extremely difficult to distinguish them experimentally from other quasiparticle states in the tiny energy difference between Majorana fermions and these states, which is beyond the energy resolution of most available techniques. Here, we circumvent this problem by systematically investigating the spatial profile of the Majorana mode and the bound quasiparticle states within a vortex in Bi₂Te₃ film grown on a superconductor NbSe₂. While the zero-bias peak in local conductance splits right off the vortex center in conventional superconductors, it splits after a finite distance ~20 nm away from the vortex center in Bi₂Te₃. This unusual splitting behavior has never been observed before and could be possibly due to the Majorana fermion zero mode. While the Majorana mode is destroyed by the interaction between vortices, the zero-bias peak splits at a conventional superconductor again. This work provides *in-situ* evidence of Majorana fermions and also suggests a possible route to manipulating them.

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Identical to their antiparticles, Majorana fermions (MF) were proposed in 1937 as an alternative to Dirac theory of ordinary fermions that carry opposite charge from their antiparticles [1]. Majoranas are the first candidate for MF in particle physics, but their Majorana status remains to be confirmed [2]. There are also proposals that quasiparticles in certain quantum condensed matter systems may be MFs. Examples include 5/2 fractional quantum Hall state, cold atoms, and chiral *p*-wave superconductors [3,4]. Experimental realization of MFs is of great significance in fundamental physics. MFs obey non-Abelian statistics, and thus can be used to develop topological quantum computation. The recent work by Fu and Kane predicted that MFs should be present as zero-energy bound states at vortex cores of an engineered heterostructure consisting of a normal *s*-wave superconductor (SC) and a topological insulator (TI) [5]. Cooper pairs are introduced via the proximity effect to the TI surface where spin and momentum are locked in the topological surface state (TSS) band [6,7]. This leads to an unusual *p*-wave-like paired state that is time-reversal invariant and robust against disorder [8]. Theoretical studies have showed that the MFs may also

reside at two ends of a semiconductor nanowire (NW) with strong spin-orbit coupling when it is contacted in *in-situ* with SC in a proper external magnetic field [9]. Several transport measurements revealed a signature of MFs, i.e., a sharp zero-bias peak in differential conductance spectrum, in various NW-SC junctions [10–14]. In InSb/Nb junction, an unconventional fractional ac Josephson effect was observed and attributed to the existence of MFs [15]. However, alternative explanations of these transport results based on disorder and/or band bending in the NWs have been proposed [16–19]. Very recently, S. Nadj-Perge *et al.* have reported their observation of MF on Fe atomic chains [20], as yet, no conclusive evidence has been established for the existence of MF [21].

In contrast, the disorder alone is unlikely to induce a zero-bias peak in a superconducting TI, which can be used to detect MFs without the complications mentioned above. Proximity effect-induced superconductivity in a TI surface has been demonstrated in several TI/SC heterostructures [22–25]. To obtain the evidence for the existence of MFs, a promising route is to detect the zero-bias bound state at vortex cores of a TI/SC heterostructure with spinning

(38) Quantum teleportation with independent sources and prior entanglement distribution over a network, Nature Photonics

Quantum teleportation with independent sources and prior entanglement distribution over a network

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Quantum teleportation¹ faithfully transfers a quantum state between distant nodes in a network, which enables revolutionary information-processing applications^{2–4}. This has motivated a tremendous amount of research activity^{5–24}. However, in the past not a single quantum-teleportation experiment has been realized with independent quantum sources, entanglement distribution prior to the Bell-state measurement (BSM) and feedforward operation simultaneously, even in the laboratory environment. We take the challenge and report the construction of a 30 km optical-fibre-based quantum network distributed over a 12.5 km area. This network is robust against noise in the real world with active stabilization strategies, which allows us to realize quantum teleportation with all the ingredients simultaneously. Both the quantum-state and process-tomography measurements and an independent statistical hypothesis test confirm the quantum nature of the quantum teleportation over this network. Our experiment marks a critical step towards the realization of a global ‘quantum internet’ in the real world.

Quantum teleportation¹ is at the heart of quantum technologies². Bennett et al. demonstrated quantum entanglement as the key to realize the transfer of information³. Quantum teleportation faithfully transfers the quantum state of a physical system (state) of the system being teleported from one node to another, which includes the preparation of distant quantum computers⁴ and quantum communication networks⁵.

Figure 1 depicts a remote quantum network. The central task for a quantum processor is state entanglement with many other nodes which constitute a state-sharing structure. The real use involves the central quantum processor by teleporting the quantum state to the remote nodes via the ready-to-use state. The state nodes perform an entanglement distribution and BSM and feedforward the transmission outcomes of the central processor. The structure can be regarded as a form of ‘ladder’ network. The quantum network, upon being made demand that each node has an independent quantum source. Therefore, of critical use in the quantum-network setting, a global quantum network is an explicit independent quantum source, even entangled distribution^{6,7} and state entanglement generation⁸. Simultaneously, in a single experimental setting of quantum teleportation in the real world, this remains an experimental challenge because of the maintenance of independent sources of being quite low

efficiency made in a quantum network with a high and dark quantum-number number. To do this, one can decrease the spectral and spatial distinguishability by applying spectral and spatial filtering. However, the real challenge is that the distinguishability of the spectral width of the quantum channel⁹ of the real world makes independent photons indistinguishable in a continuous, which was not a matter in previous experiments that were performed locally in a stable laboratory environment^{10–12}. To overcome this challenge we developed a number of experimental techniques to stabilize the quantum channel. In particular, we use the weak signal derived from the photon arrival time as the feedback loop to reduce the overall time delay between independent photons at distant nodes to 8 ps, which is equivalent compared with the 110 ps coherence time of photons. As such the quantum sources become indistinguishable photons from distant nodes in real-world in real degree of freedom and are experimentally compatible with high-speed detectors for weak-signal measurements.



Figure 1 | Structure of a quantum network with a central quantum processor. The central quantum processor is connected to many other nodes which constitute a state-sharing structure. The real use involves the central quantum processor by teleporting the quantum state to the remote nodes via the ready-to-use state. The state nodes perform an entanglement distribution and BSM and feedforward the transmission outcomes of the central processor.

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(39) Bimetallic Microswimmers Speed Up in Confining Channels, Physical Review Letters

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Bimetallic Microswimmers Speed Up in Confining Channels

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Synthetic microswimmers are envisioned to be useful in numerous applications, many of which occur in highly confined spaces. It is therefore important to understand how confinement influences swimmer dynamics. Here we study the motility of bimetallic microswimmers in linear and curved channels. Our experiments show swimmer velocities increase up to 5 times with the degree of confinement, and the relative velocity increase weakly on the fuel concentration and tank strength in solution. Experimental results are reproduced in a numerical model which attributes the velocity increase to electrostatic and electrohydrodynamic boundary effects. Our work not only helps to elucidate the confinement effect of phoretic swimmers, but also suggests that spatial confinement may be used as an efficient control method for them.

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Material transport at micro- or nanometer scales is conventionally driven by externally imposed fields, such as pressure or temperature [1,2]. In contrast, microorganisms move autonomously without external driving; they self-propel by converting local sources of energy into mechanical work [3–5]. Many synthetic microswimmers can mimic their biological counterparts in achieving autonomous motion [3–16]. For example, Janus particles generate local gradients of temperature [17], chemical concentration [12–14], or electric potential [15–19] and self-propel through respective phoretic mechanisms [20–22]. These phoretic microswimmers have demonstrated preliminary applications, such as cargo delivery [23,24], chemical sensing [25,26], and water purification [27,28]. They have also been recently used to explore the nonequilibrium physics of active matter [13,14,29].

Microswimmers are often required to move through narrow channel-like passages [30]. Hydrodynamic swimmers in such strong confinement have been investigated. Experiments showed that *Paramecium* [31] and *E. coli* bacteria [32] slow down in narrow channels; these results were later reproduced numerically [33,34]. We *et al.* showed numerically that amoeboid slow down in strong confinement [35]. Other numerical studies predicted that swimmer velocity may increase with confinement for a squimmer driven by acoustic deformations [37], a Taylor-like swimmer [36], and a rotating helix [17]. Lauder *et al.* [38] reported that elasticity of confining boundary may speed up a dipolar swimmer. These studies show hydrodynamic interaction with confinement can lead to rich swimmer dynamics.

In contrast with the numerical studies of hydrodynamic swimmers [31–38], the subject of synthetic swimmers in

confinement has remained largely unexplored until quite recently. Synthetic swimmers are found to strongly interact with boundaries, such as a flat plane [39–41], a corner [42,43,45], a colloid crystal [46], or a spherical confinement [47]. We are therefore motivated to investigate how strongly confining channels affect synthetic swimmers, particularly their motility. Our experiments show bimetallic swimmers significantly increase their velocity in channels. A numerical model reproduces our experimental results and reveals that the channels influence swimmers through electrostatic and electrohydrodynamic boundary effects. This knowledge not only helps us better understand the complex interplay between phoretic microswimmers and their environment, but also suggests control strategies for their practical applications in confined situations.

We fabricate two kinds of bimetallic microswimmers. One contains Au and Ru segments of equal length, the other Au and Rh segments. Both swimmers are rodlike and have a diameter $d_c = 0.3 \mu\text{m}$ and different length, denoted as l , (1.8 to 5.7 μm). A two-photon direct laser writing system fabricates microchannels with a submicron accuracy; the system's efficiency and flexibility allow us to systematically vary channel size and topology. Scanning electron images and detailed fabrication procedures can be found in the Supplemental Material [48]. After fabrication, we immerse swimmers and channels in a solution containing hydrogen peroxide (H_2O_2 , 5% to 30% by weight). Swimmer motion in the sample is recorded through a 60× objective at a rate of 30 frames per sec. with a camera (Basler acA2040-90um) and analyzed with standard particle tracking algorithms.

A bimetallic microswimmer propels through a self-electrophoretic mechanism. As sketched in Fig. S1a

(40) Controllable Terahertz Radiation from a Linear-Dipole Array Formed by a Two-Color Laser Filament in Air, Physical Review Letters

Controllable Terahertz Radiation from a Linear-Dipole Array Formed by a Two-Color Laser Filament in Air

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We demonstrate effective control on the carrier-envelope phase and angular distribution as well as the peak intensity of a single-cycle terahertz pulse emitted from a laser filament formed by two-color, the fundamental and the corresponding second harmonics, femtosecond laser pulses propagating in air. Experimentally, such control has been performed by tuning the filament length and the initial phase difference between the two-color laser components. A linear-dipole-array model, including the description of both the generation (via laser field ionization) and propagation of the emitted terahertz pulse, is proposed to present a quantitative interpretation of the observations. Our results contribute to the understanding of terahertz generation in a femtosecond laser filament and suggest a practical way to control the electric field of a terahertz pulse for potential applications.

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Terahertz (THz) radiation is of enormous interest for a variety of promising applications such as remote THz sensing and imaging, THz spectroscopy, THz nonlinear high field physics, and the diagnostic of a high-power laser-matter interaction [1–5]. High peak-power broadband THz radiances can be generated by two-color laser-induced air plasmas which can be precisely delivered to remote targets by laser self-guiding, avoiding thus the strong absorption of THz radiation by water vapor in the atmosphere [6–9]. THz radiation generation in air plasmas can be mainly described using four-wave rectification [10–12] and transient photocurrent [13–21] models. In particular, it is well known from previous studies that the electric field of a THz pulse emitted from a two-color laser plasma is usually with a single cycle waveform. Thus, a precise control of the carrier-envelope phase (CEP) is crucial for applications, e.g., THz nonlinear optics and spectroscopy on various systems [22–24].

In this Letter, we have performed an effective control of the spatiotemporal structure of a THz pulse emitted from a two-color laser-induced air-plasma channel (filament), including its electric field waveform (or CEP) and angular distribution as well as its peak intensity. Such control was realized through a change of the filament length as well as the initial phase difference between the two-color laser components. A theoretical model has been developed to describe both the generation and the propagation of the THz wave emitted from a filament, where the THz radiation is considered as a coherent superposition of THz waves emitted from a linear-dipole array (LDA) along the filament. A good agreement has been obtained between the experimental observations and the predictions from the theoretical model.

A laser pulse of 80 fs pulse duration and 3.5 mJ energy at 800 nm was focused by a plano-convex lens of 50 cm focal

length to form a filament in air. Before the laser pulse underwent filamentation [25,26], it was intercepted by a β -barium borate (BBO) crystal of 0.4 mm thickness, leading to second-harmonic (SH) generation as shown in Fig. 1(a). The initial phase difference ($\Delta\phi$) between the fundamental wave (FW) and its SH could be controlled by translating the BBO along the propagation axis of the laser.

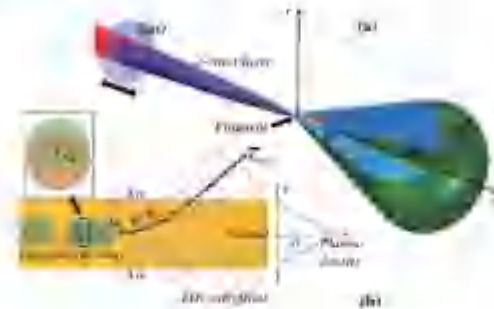


FIG. 1. (a) Schematic representation of the experimental setup at a polar coordinate (z, θ) , where the origin O corresponds to the geometric focus of the laser. The z axis corresponds to the propagation axis of the laser beam as well as the THz radiation. θ and r are used to describe the transverse propagation direction of the THz radiation. (b) Evolution of a THz wave front (red dashed curves) while a THz wave passes through different layers of a plasma channel. W_0, W_1, W_2, \dots are typical points on the front of a THz wave. The inset shows a dipole formed by charge separation inside a laser-plasma filament.

(41) Carrier-Envelope-Phase Characterization for an Isolated Attosecond Pulse by Angular Streaking, Physical Review Letters

Carrier-Envelope-Phase Characterization for an Isolated Attosecond Pulse by Angular Streaking

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The carrier-envelope phase (CEP) is a crucial parameter for a few-cycle laser pulse since it substantially determines the laser ionization. We propose for the first time to extend ultraviolet (UV) pulses, the unique strategy is directly characterizing the CEP of an isolated attosecond pulse (IAP) by angularly streaking the resulting ionization of a hydrogen atom in a combined IAP and phase-stabilized femtosecond-pulsed IR laser pulse. The fine modulations of the combined laser fields, due to the variation of the CEP of the IAP, are experimentally enlarged and the distinct time-dependent ionization ionization time, electron velocity is different ions with distinct tunneling ionization rates are angularly streaked in different directions. By measuring the resulting photoelectron momentum distribution, the CEP of the IAP can be retrieved. The characterization of the CEP of an IAP will open the possibility of generating attosecond-pulse trains.

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The CEP of a few-cycle laser pulse determines its electric field, which in turn, governs many strong-field processes observed in ultrashort physics. (1) The availability of few-cycle CEP-stabilized IR laser pulses [2,3] opens the possibility of exploring the subcycle dynamics in atoms and molecules. Over the last few decades, a series of interesting phenomena have been explored using few-cycle CEP-locked IR pulses. For example, Lindner *et al.* controlled the electron ejection from an atom at one of two possible instants, producing a Young's double-slit type interference in the photoelectron energy spectra [4]. Kling *et al.* studied the electron movement in a subcycle time scale and controlled the path of molecules breaking [5]. Eitanik *et al.* generated supercontinuum high harmonics [6], which could be synthesized into an IAP [7]. The IAP provides access to unprecedented high time resolution and has opened ultrashort science into attosecond time scales [8]. By combining an IAP and CEP-locked IR pulse, people have successfully observed many ultrashort processes with attosecond resolution, such as Auger decay [9], direct measurement of a few-cycle light wave [10], time delay in photoionization [10], tunneling processes [11], charge-directional streaking [12–14], control of absorption line shape [15], and many others.

The CEP of a few-cycle strong IR pulse has been characterized with several methods [16–20]. All these methods fundamentally depend on the nonlinearity ionization of atoms in strong laser fields. When the CEP of the IR laser pulse is varied, the symmetry of the laser waveform is modified. The slight sub-cycle modification of the laser waveform is magnified into the photoelectron since the

tunneling ionization is highly sensitive. By diagnosing the asymmetric photoelectron energy spectra, one may extend the CEP information [17] to a single shot. The ionized electron may rescatter with its parent ion and emit high harmonics, where the CEP information is also imprinted [11].

The successful characterization of the CEP in a few-cycle strong IR pulse motivates us to undertake the nontrivial task of characterizing the CEP of an IAP. Compared to a few-cycle CEP-stabilized IR field, a few-cycle CEP-stabilized IAP will make it possible to access dynamics in the sub-femtosecond period. Few-cycle IAPs have already been produced in many labs with different strategies [7,21–25]. Unfortunately, the concepts which work well for strong IR pulses cannot be directly applied to characterize the CEP of an IAP because the IAP is most likely to be extremely far too weak to sustain highly nonlinear processes. The characterization of attosecond pulses is usually performed using the reconstruction of temporal delay by interference of two-photon transitions (RABBIT) [26] or frequency-resolved optical gating (for complete reconstruction of attosecond burst (FROGCRAB)) [27] methods. Both methods are accurate but are limited to characterize the temporal envelope and the spectral phase of a pulse. The characterization of the CEP of an IAP is still a big challenge.

In this letter, we propose to use the angular streaking technique to characterize the CEP of an IAP. Note that angular streaking concepts have been used to study several interesting phenomena [28–31]. In our strategy, the prototypical hydrogen atom is tunneling limited by the

(42) Extreme Mechanics of Probing the Ultimate Strength of Nanotwinned Diamond, Physical Review Letters

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Extreme Mechanics of Probing the Ultimate Strength of Nanotwinned Diamond

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Recently synthesized nanotwinned diamond (NTD) exhibits unprecedented Vickers hardness exceeding 100 GPa [Q. Huang *et al.*, *Nature* (London) **510**, 250 (2014)]. This extraordinary finding challenges the existing understanding of material deformation and stress response under extreme loading conditions. Here we reveal by first-principles calculations (see) indenter-deformation generated stress-concentration mechanism that suppresses the propagation of bond collapse failure modes commonly known in strong covalent solids, leading to greatly enhanced peak stress and strain range in the indenter-diamond failure process. Moreover, the twin boundaries in NTD promote a strong stress concentration that drives preferential bond rupturing, producing a giant indentation strain softening. These results explain the exceptional indentation strength of NTD and offer insights into the extreme mechanics of the interface, interplay of the surface and strained crystal in probing ultrahard materials.

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Searching for materials harder and more stable than diamond has been an enduring quest driven by both practical needs and the fundamental desire to probe the ultimate limits of material strength. There have been extensive efforts exploring materials with different (but pure carbon) chemical compositions that exhibit high hardness and superior thermal stability. Notable examples include carbon nanotubes, boron nitride, and transition metal borides [1–4], but these materials fall short of reaching (or about exceeding) the strength and hardness of diamond. Another approach achieves strength enhancement in materials by controlling their microstructural size, known as the Hall-Petch effect [5,6], based on the idea that reducing grain (or twinning) size would suppress the nucleation and motion of dislocations by grain (or twin) boundaries. This approach has produced remarkable results in nanostructured metals [7] and strong covalent solids like boron nitride [8] and diamond [9], turning them harder than their bulk counterparts. In further reduced grain (or twin) sizes, however, an opposite phenomenon, known as the inverse Hall-Petch effect [10,11], takes place when the atomic-scale weaker interactions at the grain (or twin) boundaries dominate the plastic deformation process.

It was recently reported [12] that a nanotwinned diamond (NTD) structure exhibits unprecedented high Vickers indentation hardness (exceeding 200 GPa), which is twice the commonly quoted hardness—value of about 100 GPa for single-crystal cubic diamond. This result surpasses by a surprisingly large margin the hardness of all previously reported materials. More intriguing, this colossal hardness enhancement in NTD cannot be explained by the Hall-Petch effect since the typical twin boundary (TB) separation is $\lambda = 5\text{--}6\text{ nm}$ in the synthesized NTD and well inside the

reverse Hall-Petch region (below about 50 nm), where the TB intersections are expected to dominate. To elucidate the unexpected strengthening of NTD, it is essential to examine the structural and stress response to indentation loading.

In this work, we explore the extreme mechanics of probing the ultimate strength of NTD. We employ first-principles calculations to examine the deformation and stress response of both the indenter and indented material under a variety of loading conditions. From calculated compressive stress and we set up a compressible Vickers indenter model to account for the large but finite normal stress beneath the indenter, which has been assumed to be indented (i.e., incompressible indenter) in past studies. This new indenter-deformation generated stress-concentration presents the propagation of bond collapse failure modes commonly known in strong covalent solids, leading to greatly enhanced peak stress and strain range compared to those obtained from the incompressible indenter model. Moreover, we have identified a robust mechanism for bond configuration learning from a TB-induced stress concentration, which produces a giant indentation strain softening in NTD. These results explain the unprecedented high strength and hardness of NTD and the insights into the novel structural deformation and strengthening mechanisms advance the fundamental understanding of material behavior under extreme indentation loading conditions.

Stress-strain relations from first-principles calculations provide an accurate description of material deformation and strength [13–23], and the results can be directly compared to well-controlled nanoindentation measurements [24–26]. Using indentation, shear instability usually precedes the initiation of cracks and delamination [27], signaling the onset of incipient plasticity [28,29] and bond-collapse onset

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LETTER

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Granular materials flow like complex fluids

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Granular materials such as sand, powders and grains are ubiquitous in daily life and in industrial and geotechnical applications^{1–3}. These disordered systems form stable structures when unperturbed, but in the presence of external influences such as tapping of the container they relax, becoming fluid in nature. It is often assumed that the relaxation dynamics of granular systems is similar to that of thermal glass-forming systems^{4,5}. However, so far it has not been possible to determine experimentally the dynamic properties of three-dimensional granular systems at the particle level. This lack of experimental data, combined with the fact that the motion of granular particles involves friction (obscures the motion of particles in thermal glass-forming systems) does not, in general, allow an accurate description of the relaxation dynamics of granular materials to be lacking. Here we use X-ray tomography to determine the microscale relaxation dynamics of hard granular ellipsoids subject to an oscillatory shear. We find that the distribution of the displacements of the ellipsoids is well described by a gamma law⁶ (which is similar to a Gaussian distribution for small displacements but has a heavier tail for larger displacements), with a shape parameter that is independent of the amplitude of the shear strain and of the time. Despite this universality, the mean squared displacement of an individual ellipsoid follows a power law as a function of time, with an exponent that does depend on the strain amplitude and time. We argue that these results are related to microscale relaxation mechanisms that involve friction and memory effects to hereby the motion of an ellipsoid at a given point in time depends on its previous motion. Our observations demonstrate that, at the particle level, the dynamic behaviour of granular systems is qualitatively different from that of thermal glass-forming systems, and is instead more similar to that of complex fluids. We conclude that granular materials (as such) even when the driving strain is weak.

Static and driven granular systems exhibit an interplay like thermal glassy systems⁷. However, they also exhibit phenomena such as shear and force chains⁸, and show the kink and avalanche behaviour which demonstrates that their properties are also affected by many parameters, friction and adhesion. Although there are many studies on the microscale properties of granular materials^{9–11}, the investigation has focused on the structure and dynamics of granular systems at the particle level and in three dimensions, owing to the experimental difficulty of determining the position and velocity of individual particles^{12–14}. However, this information is crucial for obtaining a fundamental understanding of the properties of these systems at the microscopic. Previous information about the system can be obtained through simulation experiments and simulations that provide direct insight into propagation and ageing dynamics^{15–17}, memory effects¹⁸, dynamic heterogeneity¹⁹, cluster distribution²⁰, avalanches²¹ and the role of contact^{22,23}. Our X-ray tomography study shows that the relaxation dynamics (mean squared displacement of the particle displacements) is well described

by a gamma law⁶ and by a weak memory effect. Surprisingly, we find that the expected time-averaged behaviour of the dynamics of glass-forming systems, whereby particles are microscopically trapped by their nearest neighbours^{24–26}, is absent in the system that we study. Determining a fundamental mechanism of the microscale relaxation dynamics of thermal glass-formers and granular systems.

The particles that we study are hard elastic prolate ellipsoids (masses $m = 0.25$ g) made from an epoxy resin of 1.05 g cm⁻³. The particles are ground into a regular form with side walls that can be used to prepare a cubic glass in the system and a force chain in the top of the particles (see Methods). We used four strain amplitudes: $\gamma = 0.05, 0.1, 0.2$ and 0.26 , and a strain rate of angular $\dot{\gamma} = 10^{-3}$ s⁻¹. This is very slow compared to the inertial number i about 1.4×10^{-3} , meaning that the experiments can be considered as quasi-static²⁷. Before taking any measurements, we performed hundreds of oscillatory shear cycles to bring the system to a stationary state (see Methods). After each cycle we had 4 oscillatory scans and measured the positions and orientations of all particles²⁸ (Fig. 1a, Supplementary Tables 1 and 2) over the lifetime of the particles and recorded data (Fig. 2) shows some typical results. The following we measure the means of random cycles and their standard deviation.

The translational mean squared displacement (MSD) $\langle \Delta r^2(t) \rangle = \langle (x_i(t) - x_i(0))^2 + (y_i(t) - y_i(0))^2 \rangle$ of the position of particle i at time t of the particles is shown in Fig. 3b. We see that the MSD has a plateau at very small times even if the driving is weak, which demonstrates that the particles exhibit a thermal glassy behaviour²⁹ at short times. It is also controlled by the microscale memory function, which does not show any sign of a slow relaxation time (Supplementary Fig. S1). $\langle \Delta r^2(t) \rangle$ is proportional to time at small times, indicating normal diffusive behaviour. It is increased to 0.26, then the γ dependent case of $\langle \Delta r^2(t) \rangle$ can be well described by a power law with an exponent $\alpha = 0.61 \pm 0.02$ that is the different simulation. This result is surprising because it could be expected that stronger driving would induce order, such as force chains and force chains and heterogeneity. However, this is not the case and we observe the order in the microscopic simulation below. For $\gamma = 0.05$ and 0.1 , we find a power law with an exponent close to 1.0 (normal diffusion) instead of constant at $\alpha = 0.6$ and $\alpha = 1.0$ respectively. The γ dependence of α on this case can be well described by a power law with exponent close to 0.27. The data represent well for $\alpha = 0.59$. Some possible ways of explaining the presence of anomalous diffusion are to assume that the underlying configuration space has a fractal nature³⁰, but the constant velocity is about a 1.4×10^{-3} s⁻¹, so that the structure of the particles is stable^{31,32}, however, we show that from the microscopic simulation can be characterized by means of data in the dynamic.

The rotational dynamics is clearly more slow than the translational diffusion $\langle \Delta \theta^2(t) \rangle = 0.20 \log(t)$, where $\theta = 0$ is arbitrary. Despite a γ dependence, indicating that for this time scale the particles are the rotational diffusion will be the dynamics and of the order of $\log(t)$ for $\gamma = 0.1$ the rotational diffusion in the system

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(44) Experimental studies of vibrational modes in a two-dimensional amorphous solid, Nature Communications



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Experimental studies of vibrational modes in a two-dimensional amorphous solid

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The boson peak, which represents an excess of vibrational states compared to Debye's prediction at low frequencies, has been studied extensively, and yet, its nature remains controversial. In this study, we focus on understanding the nature of the boson peak based on the spatial heterogeneity of modulus fluctuations using a simple model system of a highly jammed two-dimensional granular material. Despite the simplicity of our system, we find that the boson peak in our two-dimensional system shows a shape very similar to that of three-dimensional molecular glasses when approaching their boson peak frequencies. Our finding indicates a strong connection between the boson peak and the spatial heterogeneity of shear modulus fluctuations.

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(45) Origin of Noncubic Scaling Law in Disordered Granular Packing, Physical Review Letters

Origin of Noncubic Scaling Law in Disordered Granular Packing

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Recent diffraction experiments on metallic glasses have revealed an unexpected noncubic scaling law between density and average interatomic distance, which led to the speculation of the presence of fractal glass order. Using x-ray tomography we identify how a similar noncubic scaling law in disordered granular packing of soft-sphere particles. We find that the scaling law is directly related to the contact neighborhood: the first nearest-neighbor shell and iterations, it is closely connected to the phenomenon of jamming. The seemingly unrelated scaling exponent around 2.5 arises from the contact neighborhood with a contact number around 8, and we argue that the exponent should not be universal.

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The origin of dynamic arrest and mechanical rigidity in amorphous materials remains one of the important unsolved questions in condensed matter physics [1–3]. Whether it has a structural origin or is just a dynamic phenomenon remains controversial [4,5]. For metallic glasses, it has long been speculated that dense local packing structures of short-range order serve as the building blocks in these systems [6]. However, how these local structures can be extended to medium or large scales remains as present a mystery due to the existence of geometric frustration or intrinsic chemical disorder [7–9]. Recently, it has been proposed that metallic glasses possess a medium-range fractal order, which could rationalize the commonly observed noncubic scaling law between the position of the first diffraction peak and the bulk density found in neutron and x-ray scattering experiments on these systems [10–11]. The first diffraction peak position is usually associated with the largest interplane distance in crystals or the typical nearest-neighbor distance in liquids [14–16], and it shows in these systems a power law of exponent β as a function of the bulk density since they are three-dimensional by nature. It is therefore surprising that the scaling exponent obtained for metallic glasses under density change induced by either pressure or compression tuning has instead values that are between 2.3 and 2.5 [10–11]. The origin of this anomalous scaling law has been attributed to the presence of a regular or statistical fractal network formed by glass order [6–10,12]. In this picture, the atoms move affinely relative to each other under deformation, and their coherent scattering intensity yields the noncubic law. However, real metallic glasses are in fact quite compact, while a large-scale fractal structure has zero mass density. Therefore, in order for this picture

to be valid, one requires that a substantial amount of atoms exist within the fractal structures, which do not contribute coherently to the sharp scattering peaks [11]. Another possibility is that the fractal structure only exists up to a finite length scale, above which the system is still homogeneous and three-dimensional [6,12]. These explanations are appealing since they naturally refer to a fractal medium-range glass order, such as percolatingicosahedral structures for metallic glasses, and therefore explain how glass order extends in space. However, the interpretation of the existence of the noncubic law based on the fractal picture is not without controversy [11], and sometimes one also finds deviations from the noncubic law [18,19].

In this work, we provide microscopic insight into this problem by studying the three-dimensional packing of spherical granular particles, which is a prototypical hard-sphere glass former and has long been considered as a structural model for metallic glasses [20–22]. We identify a noncubic scaling law in our system and provide evidence that its origin is local, i.e., without resorting to any fractal structure. Instead, it results from a complex structural evolution of the first-shell neighbors when the packing structure varies, controlled usually by the contact neighbors as required by mechanical stability, and the global behavior or a simple statistical coverage of the local ones. Therefore, such phenomenon is directly related to the jamming phenomenon and might be universal near the jamming criticality [23–24]. In the experiment, we used synchrotron x-ray computed tomography (CT) technique to obtain the packing structure of packing with a wide range of packing fractions Φ [25–27] (see Supplemental Material [35]). In the following, we use the average particle diameter as a unit of length.

(46) Dynamical Transition of Collective Motions in Dry Proteins, Physical Review Letters

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Dynamical Transition of Collective Motions in Dry Proteins

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Water is widely assumed to be essential for protein dynamics and function. In particular, the well-documented “dynamical” transition at ~ 200 K, at which the protein changes from a rigid, nonfunctional form to a flexible, functional state, as detected in hydrogenated protein by incoherent neutron scattering, requires hydration. Here, we report on coherent neutron scattering experiments on perdeuterated proteins and reveal that a transition occurs in dry proteins at the same temperature resulting primarily from the collective heavy-atom motions. The dynamical transition discovered is intrinsic to the energy landscape of dry proteins.

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Water is intimately involved in protein function [1–7]: in cozymes, for example, hydration enhances enzymatic rates, and noncatalytic water molecules can aid in transporting substrates and products into the catalytic sites [1,8,9]. Numerous experiments have suggested that a minimum hydration level of $h = 0.2$ (i.e., water/protein) is required for enzymes to function [1,10–12]. Hydration may also enhance the internal motions required for the functional chemical steps [1,7,8,10]. Indeed, $h = 0.2$ has also been reported to be the minimum hydration level required for the widely studied dynamical transition in proteins at ~ 200 K, also named as the “glass” transition in analogy to the glass-transition process in glass-forming systems, at which internal protein motions change from nonfunctional, glasslike, rigid, harmonic vibrations to incorporating liquidlike, flexible, anharmonic dynamics required for function [1,5,13–19]. A large body of experimental and simulation work has suggested that this transition is coupled to the activation of the translational motion of the hydration water on the protein surface [15,7,9,13,17–22]. However, the energy landscape of a globular protein is intrinsically complex and highly anharmonic, leading to the question as to whether the protein dynamical transition may be an intrinsic property of dry proteins.

Deuteration is key in the present observations. For simplicity, in what follows hydrogenated and perdeuterated samples are denoted using the prefixes H and D. The neutron data collected on H proteins are mostly incoherent

scattering signals ($\sim 90\%$, see Fig. S1(a) in the Supplemental Material [23]), resulting from self-correlations of the motions of hydrogen atoms. In contrast, the scattering from D proteins is primarily coherent ($\sim 90\%$, Fig. S1(b) in Ref. [23]), arising mostly from cross-correlations in motions between protein atoms [20–22], and is dominated by scattering signals from the heavy (non-H) atoms ($\sim 70\%$), especially those on the backbone (Fig. S1(c) in Ref. [23]).

Here, by performing neutron scattering and molecular dynamics (MD) simulation on lyophilized (hydration level, $h = 0.02$, Fig. S2 in Ref. [23]) perdeuterated cytochrome P450 (CYP) and green fluorescent protein (GFP) (Figs. 1(a) and 1(b)), we show that the motions of heavy atoms in an essentially dry protein themselves present a dynamical transition at ~ 200 K.

On the high-flux backscattering spectrometer HFBS at NIST, we measured $S(q, \Delta t)$, the intensity of the elastic peak of the dynamic structure factor, which is an estimate of the average amplitude of the atomic motions (up to the temporal-resolution time of the instrument, Δt , which is ~ 1 ns [25]). Figure 1 presents the temperature dependence of $S(q, \Delta t)$ for both dry ($h = 0.02$) H and D samples. The presence of a downward change in gradient in the temperature dependence of $S(q, \Delta t)$ indicates that an anharmonic dynamic process is activated in the corresponding temperature [24]. Figures 1(c) and 1(d) show that both H-CYP and H-GFP present a kink at ~ 125 K. This transition has been observed several times previously in various hydrogenated

(47) Strong Field Theories beyond Dipole Approximations in Nonrelativistic Regimes, Physical Review Letters

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Strong Field Theories beyond Dipole Approximations in Nonrelativistic Regimes

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The exact nondipole Volkov solutions to the Schrödinger equation and Pauli equation are found, based on which a strong field theory beyond the dipole approximation is built for describing the nondipole effects in nonrelativistic laser-driven electron dynamics. The theory is applied to investigate momentum partition laws for multiphoton and tunneling ionization and especially shows that the complex interplay of a laser field and Coulomb's action may reverse the expected photoelectron momenta along the laser propagation direction. The magnetic-spin coupling does not bring observable effects on the photoelectron momentum distribution and can be neglected. Compared to the strong field approximation versus the dipole approximation, this theory works in a much wider range of laser parameters and lays a solid foundation for describing nonrelativistic electron dynamics in both short wavelength and undiffracted regimes where nondipole effects are unavoidable.

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Two kinds of processes are very fundamental in strong laser-atom and laser-molecule interactions, i.e., the deposition of photon energies into atoms and molecules and the transfer of photon linear momentum into targets. Plenty of studies focused on the energy deposition and a series of intriguing ultrafast phenomena were explored, such as tunneling ionization [1], high harmonic generation and its synthesis of attosecond light pulses [2], nonsequential double ionization [3], coherent control of electron localization [4–7], and energy sharing between electrons and nuclei [8]. For the second one, the transfer of photon linear momentum attracts much less attention due to the fact that a photon is not an effective momentum carrier as compared to a nonrelativistic electron.

Thanks to the detection techniques with high resolutions developed recently, Saneek *et al.* [9] and Ludwig *et al.* [10] reported that photoelectrons acquire a momentum shift along the direction of laser propagation (termed as longitudinal momentum shift in the rest of the Letter) during tunneling ionization, and this photon momentum transfer is also reported in the time-dependent Dirac equation simulation [11], time-dependent Schrödinger equation (TDSE) simulations [12,13], and classical trajectory Monte Carlo calculations [14]. Some characters of the longitudinal momentum shift [15–20] were explored by calculations of strong field approximation (SFA). Though it is well recognized that the longitudinal momentum shift must be due to the nondipole laser-atom coupling, the exact nondipole Volkov solution was not built into these studies. To overcome this obstacle, some attempts have been made. Chelkowska *et al.* adopted the Volkov solution of the Kohn-Gordon equation in their SFA calculations [21]. Krajczyk *et al.* made an analogy with quantum electrodynamics in obtaining the transition amplitude [22], Crivello and others

reverted the space-time dependent laser vector potential as $\mathbf{A}(\mathbf{r}, t) = \mathbf{A}(t) + (\mathbf{r} - \mathbf{r}_0)/c \mathbf{E}(t)$ [15,22–24]. Here, $\mathbf{A}(t)$ and $\mathbf{E}(t)$ are the dipole approximated laser vector potential and laser electric field, and c is the light speed. To build a self-consistent strong field theory beyond the dipole approximation is still difficult because there is no exact nondipole Volkov-wave function to the Schrödinger equation until today, though the Volkov solution to the Schrödinger equation within the dipole approximation had been obtained around eighty years ago.

In this Letter, we analytically solved the TDSE for a free electron in a monochromatic laser field and achieved the exact nondipole Volkov solution. Based on this, we built the nondipole SFA theory, which is then applied to describing how the electron acquires the momentum along the laser propagation direction, and how rescattering alters the photon momentum partition between electrons and nuclei. For very wide ranges of laser parameters, this theory is self-consistent and is able to grasp the main dynamics and offer intuitive physical explanations with easy computations.

We used the light cone coordinate $\eta = t - (z/c)$ to describe the laser pulse by assuming it propagates along the z axis. The components of the laser vector potential are $\mathbf{A}(\mathbf{r}, t) = \mathbf{A}(\eta) = (A_x(\eta), A_y(\eta), 0)$. The behavior of a free electron embedded in such a laser field is described by the TDSE as [atomic units (a.u.) are used unless stated otherwise]

$$i \frac{\partial}{\partial t} \psi(\mathbf{r}, t) = \frac{1}{2} (\mathbf{p} + \mathbf{A}_\perp(\eta))^2 + p_z^2 \psi(\mathbf{r}, t), \quad (1)$$

with \mathbf{p} being the electron momentum operator. We recognized that the Volkov solution of Eq. (1) has the form

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(48) Structural and topological nature of plasticity in sheared granular materials, Nature Communications



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OPEN

Structural and topological nature of plasticity in sheared granular materials

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Upon mechanical loading, granular materials yield and undergo plastic deformation. The nature of plastic deformation is essential for the development of the macroscopic constitutive models and the understanding of shear band formation. However, we still do not fully understand the microscopic nature of plastic deformation in disordered granular materials. Here we used synchrotron X-ray tomography technique to track the structural evolutions of three-dimensional granular materials under shear. We establish that highly distorted coplanar tetrahedra are the structural defects responsible for microscopic plasticity in disordered granular packings. The elementary plastic events occur through flip events which correspond to a neighbor switching process among these coplanar tetrahedra (or equivalently as the rotation motion of 4-ring disclinations). These events are discrete in space and possess specific orientations with the principal stress direction.

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(49) Manipulation of polarizations for broadband terahertz waves emitted from laser plasma filaments, Nature Photonics



Manipulation of polarizations for broadband terahertz waves emitted from laser plasma filaments

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Polarization control of broadband terahertz waves is essential for applications in many areas, such as materials science, medical and biological diagnostics, near-field communications and public securities. Conventional methods for polarization control are limited to narrow bandwidth and often with low efficiency. Here, based on theoretical and experimental studies, we demonstrate that the two-colour laser scheme in gas plasma can provide effective control of elliptically polarized terahertz waves, including their ellipticity, azimuthal angle and chirality. This is achieved with a circularly polarized laser at the fundamental frequency and its linearly polarized second harmonic, a controlled phase difference between these two laser components, as well as a suitable length of the laser plasma filament. Flexible control of ellipticity and azimuthal angle is demonstrated with our theoretical model and systematic experiments. This offers a unique and flexible technique on the polarization control of broadband terahertz radiation suitable for a wide range of applications.

Plasmas driven by two-colour femtosecond-laser pulses are an avenue for tunable, on-demand, broadband terahertz waves with reasonable efficiency that are suitable for many applications. For many of these applications, effective control of the polarization state is a key prerequisite. In material science, for instance, the polarization direction of a circular wave plays a pivotal role in optical microscopy, spectroscopy, diagnosis^{1,2} and terahertz information technologies. Meanwhile, the ellipticity of the terahertz wave is an important parameter in biological research, immunity and its stability is critical for studying chiral (chirality of systems molecules and spin dynamics in solid materials).^{3,4} It remains a challenge to realize flexible control of the polarization states of the broadband terahertz pulses. Conventional methods, such as waveplate polarizers and Fresnel rhombs wave plates, have limited flexibility for polarization modulation and are mainly applicable to terahertz waves mainly a narrow spectral bandwidth. Although it was demonstrated that linear and circularly polarized terahertz pulses originate from laser plasma sources with application of an external magnetic field,^{5–7} however, terahertz pulses with linear and elliptical polarization have been observed by two-colour coherent control of the motion of surface electrons.⁸ Additionally, it has been shown that the manipulation of the polarization state always simultaneously involves the azimuthal angle and ellipticity in these schemes, which restricts the flexibility of polarization modulation for practical applications.

Here, we introduce a new route to realize flexible control of polarization states for terahertz pulses produced by the two-colour laser scheme based on the theoretical and experimental investigations. We demonstrate that polarized terahertz pulses with any desired azimuthal angle and ellipticity can be realized by controlling the relative phase between two-colour laser components and the length

of a filament, respectively. The polarization control can be flexibly realized by changing the polarization state of the pump laser.

Results and discussion

Figure 1a depicts the experimental setup geometry. A linearly polarized pulse with any polarization state in air plasma, excited by a two-colour laser pulse, composed of a fundamental wave (FW) and its second harmonic (SH), the two terahertz pulses and their four laser pulses, the polarization states are characterized by their ellipticity and azimuthal angle, which are the transverse or the xy of electric field vector projected on the xy -plane, denoted as (E_x^2/E_y^2) , $(\theta = 0^\circ)$ and $(\theta = 90^\circ)$ for the terahertz pulse and incident laser FW and SH, respectively. The axial relative phase between the FW and SH pulses (φ_r) was controlled by flexibly manipulating a fused silica waveplate with the phase compensator (PC) constituted an alpha-halometer from two birefringent crystals (cut at fixed angles) with a dual-wavelength plate (DWP). Polarization of the terahertz pulse was characterized by decomposing two orthogonal polarized electric field vectors of the terahertz pulse (Fig. 1b), measured by an electro-optical sampling method with a femtosecond DWP (DWP) crystal at two orthogonal directions.^{9–11} See Methods for measurement for terahertz polarization. Polarization of the FW and SH pulses (P^2) is controllably by simultaneously manipulating DWP and QWP located after the PC. In our scheme, the FW is linearly polarized with $\theta = 0^\circ$ and $\theta = 90^\circ$ while the polarization of the SH pulse remains linearly polarized with $\theta = 0^\circ$ and $\theta = 90^\circ$. As shown in the section “Manipulation of terahertz polarization”, these polarization states for the driving laser can provide the high flexibility of terahertz polarization. The ellipticity E_x^2/E_y^2 required for effective discovery of chiral spread terahertz pulses can be modified by inserting a crystal wafer along the filament. Because

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(50) Coulomb-Driven Relativistic Electron Beam Compression, Physical Review Letters

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Coulomb-Driven Relativistic Electron Beam Compression

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Coulomb repulsion between charged particles is a well-known phenomenon in many areas of research. In general, the Coulomb repulsion force broadens the pulse width of an electron bunch and limits the temporal resolution of many scientific facilities such as ultrafast electron diffraction and x-ray free-electron lasers. Here, we demonstrate a scheme that greatly makes use of the Coulomb force to compress a relativistic electron beam. Furthermore, we show that the Coulomb-driven bunch compression process does not introduce additional timing jitter, which is in sharp contrast to the conventional radio-frequency bunching technique. Our work not only leads to enhanced temporal resolution in electron-beam-based scientific instruments that may provide new opportunities in probing natural systems far from equilibrium, but also opens a promising direction for advanced beam manipulation through self-field interactions.

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Ultrafast electron diffraction (UED) in which a laser-pulsed electron probe is excited by a pump laser and then probed by a delayed electron pulse has been widely used for watching atoms in motion during structural changes of which the characteristic time scale extends from a few femtoseconds to a few picoseconds [1,2]. With the advent of ultrashort lasers, the temporal resolution in UED applications now depends primarily on the pulse width and timing jitter of the electron beam. In x-ray free-electron lasers (FELs [3–7]), because the x rays are produced by sending a high-brightness relativistic electron beam through a series of magnetic undulators [8–10], the temporal resolution in laser-pump x-ray-probe experiments is also largely dependent on the electron beam pulse width and timing jitter. Therefore, one of the long-standing goals in both the FEL and UED communities is to generate a high-brightness electron beam with a short pulse width and small timing jitter.

According to Coulomb's law, the magnitude of the electric force is directly proportional to the charge and inversely proportional to the square of the distance. So when millions of electrons are produced with a femtosecond laser pulse through photoemission, they tend to repel each other. The Coulomb force, in general, accelerates the electrons in the front part of the bunch while it decelerates those in the back, leading to pulse broadening [11]. Fortunately, the elongated electron beam can be compressed with an radio-frequency (rf) buncher [12–15]. In this method, the electron beam is first sent through an rf buncher cavity, where the bunch head is decelerated to a lower energy while the bunch tail is accelerated to higher energy after passing through a

dispersive element (typically, a drift for keV and MeV electrons) or a magnetic chicane for ultrarelativistic electrons (see, e.g., [16]), the bunch tail may catch up with the bunch head, leading to bunch compression. While this method has been widely used in the UED and FEL community, it is also realized that the phase jitter in the rf cavity leads to considerable timing jitter, which limits the temporal resolution achievable [15,16]. The physics behind this effect is rather simple. E.g., in practical cases, the bunch centroid does not always pass through the cavity at a constant phase, and the variation of beam energy from this phase jitter is translated to the variation of the time of flight after passing through the dispersive element.

It should be noted that the Coulomb force is not always detrimental. The Coulomb force has been used before to suppress beam shot noise in an optical wavelength [17,18], as well as to enhance the electron beam density modulation at a terahertz wavelength [19,20]. In this Letter, we demonstrate a scheme that makes use of the Coulomb force to compress a relativistic electron beam without introducing additional timing jitter. The setup is shown in Fig. 1. The laser is first split into two pulses with a beam splitter (BS1) so that the energy, pulse width, transverse size, and timing of each pulse can be independently adjusted. The main pulse (consisting of 95% of the energy) is used to produce the drive beam, and the remaining part is used to produce the target beam in the σ -band (2850 MHz) photo-cathode of gun. After BS1, the main pulse is sent through four π barium borate birefringent crystals with the first one (7.5 μm thick) used to produce a double pulse separated by 6.4 ps and the remaining three of them

(51) Experimental Machine Learning of Quantum States, Physical Review Letters

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Experimental Machine Learning of Quantum States

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Quantum information technologies provide promising applications in communication and computation, while machine learning has become a powerful technique for extracting meaningful structure in “big data.” A crossover between quantum information and machine learning represents a new interdisciplinary area attracting progress in both fields. Traditionally, a quantum state is characterized by quantum-state tomography, which is a resource-consuming process when scaled up. Here we experimentally demonstrate a machine-learning approach to construct a quantum-state classifier for identifying the separability of quantum states. We show that it is possible to experimentally train an artificial neural network to efficiently learn and classify quantum states, without the need of obtaining the full information of the states. We also show how adding a hidden layer of neurons to the neural network can significantly boost the performance of the state classifier. These results shed new light on how classification of quantum states can be achieved with limited resources, and represent a step towards machine-learning-based applications in quantum information processing.

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Over the past few years, there has been a significant advancement in an emerging field of quantum machine learning [1,2], where quantum information meets the modern information-processing technologies. On one hand, various modern quantum technologies, such as quantum communication [3–5], quantum computation [6,7], and quantum metrology, have been experimentally demonstrated by exploiting quantum entanglement as a resource [8,9]. On the other hand, machine learning, a modern technique for making predictions by mining information from “big data” [10], has been proven as one of the most successful achievements in artificial intelligence. Notable examples include self-driving cars and the famous AlphaGo, which surpasses the top human players at the game “Go” [11,12].

The key question in quantum machine learning is how to develop new ideas for applying technologies in machine learning to quantum information, or, vice versa, to gain advancements in both fields. In fact, several promising steps along both directions have already been taken in the community. In particular, machine learning can in principle exploit quantum superposition to enhance its performance. For example, quantum versions of the principal component

analysis [13] and support vector machines [14] have been proposed. Other quantum algorithms along with their experimental implementations [15–19] have also been demonstrated in recent years to broaden the versatility of machine learning.

Additionally, machine learning can also be applied to certain quantum tasks, from classifying separability of quantum states [20,21] to classifying phases in condensed matter physics [22,23], and even the development of new classical algorithms for solving many-body systems [24]. These results suggest that machine learning of quantum states represents a new platform for solving problems in quantum information science.

Here we report an experimental machine learning of quantum states where an artificial neural network (ANN) is trained for classifying the separability of some quantum states, given that only *partial information* about the quantum states is available. More specifically, based on the experimental data, we have constructed a quantum-state classifier [20], which generalizes the pattern recognition in learning theory for quantum data. In the classical setting, a set (x_i, y_i) of data x_i and label y_i are supplied as the training set for the machine-learning program, and the

(52) Frozen Condition of Quantum Coherence for Atoms on a Stationary Trajectory, Physical Review Letters

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Frozen Condition of Quantum Coherence for Atoms on a Stationary Trajectory

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The quantum coherence (QC) of two interacting atoms on a stationary trajectory is investigated. We develop a formalism to characterize the properties of atoms on a stationary trajectory. We give a criterion under which QC is frozen to a nonzero value. The frozen condition that vanishing super- or subradiant decay rate is not so sensitive to the initial condition of state. We show that enhanced QC and a subradiant state can be gained from the initial state.

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Introduction.—QC as a consequence of the quantum state superposition principle is one of the key features that results in nonclassical phenomena. It is a powerful resource in quantum information theory as well as entanglement and discord-type quantum correlation. Though it is important in fundamental physics, only recently have relevant steps been attempted to develop a rigorous framework to quantify QC for general states, such as the relative entropy of QC and the l_1 norm [1].

A long-standing and significant issue concerning QC is the decoherence induced by the inevitable interaction between the system and environment. In the past few years, several proposals have been suggested for fighting against the deterioration of QC, for instance, decoherence-free subspaces [2,3], quantum error correction codes [4], dynamical decoupling [5], or quantum Zeno dynamics [6,7]. Recently, the conditions of sustaining long-lived QC were investigated [8]. It was shown that QC can remain unchanged with time (freezing coherence) and the frozen condition (FC) for two qubits undergoing local identical bit-flip channels with Bell-diagonal initial states was only dependent on the initial condition of the states [9].

In a realistic physical system, atoms usually cannot be handled simply as noninteracting individual qubits when atomic spacing is small. Besides, for an ensemble of atoms, motion and temperature are also important factors which should be taken into account. Then, searching for a general FC for interacting atoms under normal conditions is necessary in practice.

In this Letter, we investigate the FC of QC for two identical two-level atoms on a stationary trajectory which has a characterization that the geodesic distance between two points on the trajectory depends only on the proper

time interval [10,11]. Thus, for a stationary trajectory, field correlation function is invariant under translations in time. Besides, the stationary trajectory guarantees that the atoms have stationary states. The inertial atom in a Minkowski vacuum or thermal bath and a uniformly or circularly accelerated atom viewed by an instantaneous inertial observer are all stationary. We find that the QC for interacting atoms on a stationary trajectory can also be long lived; however, the FC is not so sensitive to the initial condition of the single excitation state but to the super- or subradiant decay rate of atoms. We develop a formalism to describe atoms on a stationary trajectory and give the general relationship between the quantum characterizing properties of atoms. Besides, we show that enhanced QC and the subradiant state can be obtained from the initial state.

Formalism.—We consider two identical two-level atoms moving on stationary trajectories $\mathbf{x}_i(\tau) = (\mathbf{x}_i(\tau), \bar{x}_i(\tau))$ in a fluctuating vacuum electromagnetic field, where $(\mathbf{x}_i, \bar{x}_i)$ are the Minkowski coordinates of atom i referring to an inertial reference frame and τ denotes proper time of these two interacting atoms (see Fig. 1). The total Hamiltonian of the coupled system can be described by $H = H_0 + H_1 + H_c$. Here, H_0 is the free Hamiltonian of atoms, and its explicit expression in the Schrödinger picture is $H_0 = \sum_{j=1,2} \omega_j n_j \sigma_j^z$, where ω_j is the level spacing of the two-level atoms, $\sigma_j^+ = |\epsilon_j\rangle\langle g_j|$ and $\sigma_j^- = |g_j\rangle\langle \epsilon_j|$ are, respectively, the raising and lowering operators of the atom j . The free Hamiltonian H_c with respect to τ takes the form $H_c = \sum_{\mathbf{k}, \lambda} \omega_{\mathbf{k}} a_{\mathbf{k}, \lambda}^\dagger a_{\mathbf{k}, \lambda} (\partial/\partial t) |12\rangle$. Here $a_{\mathbf{k}, \lambda}^\dagger$, $a_{\mathbf{k}, \lambda}$ are the creating and annihilation operators for a photon with momentum \mathbf{k} , frequency $\omega_{\mathbf{k}}$, and polarization λ . The Hamiltonian H_1 that

(53) Gradual Crossover from Subdiffusion to Normal Diffusion: A Many-Body Effect in Protein Surface Water, Physical Review Letters

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Gradual Crossover from Subdiffusion to Normal Diffusion: A Many-Body Effect in Protein Surface Water

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Dynamics of hydration water is essential for the function of biomacromolecules. Previous studies have demonstrated that water molecules exhibit subdiffusion on the surface of biomacromolecules, yet the microscopic mechanism remains unclear. Here, by performing neutron scattering, molecular dynamics simulation, and analytic modeling on hydrated perdeuterated protein powder, we found water molecules jump randomly between trapping sites on protein surfaces, whose waiting times obey a broad distribution, resulting in subdiffusion. Moreover, the subdiffusive exponent gradually increases with observation time towards normal diffusion due to a many-body volume-exclusion effect.

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Introduction.—Water is the solvent of life, playing a crucial role in determining the native structure, dynamics, and function of biological macromolecules [1–3]. The diffusive motions of water molecules not only aid the transportation of the function-required essential ingredients, e.g., protons, ions, and substrates across membranes and into the catalytic site of enzymes [4,6,7], but also render internal flexibility to the biomacromolecules. This flexibility is likely to be crucial for the biofunction, as it is absent when the biomacromolecules are dehydrated [1–3,8]. On the other hand, as an active solvent, the biomacromolecule will significantly alter the structure and dynamics of the hydration water molecules surrounding it [9–13]. Both experiments and simulations showed that the diffusive motions of the hydration water on the surface of various biomacromolecules: DNA [10], RNA [14], proteins [11], and lipid membranes [12], are significantly retarded as compared to bulk water, and present an anomalous subdiffusion [11,12,15,16]. This subdiffusive motion is characterized by a fractional power-law dependence on time of the mean square atomic displacement (MSD), i.e., $\langle Y^2(\Delta t) \rangle \sim t^\beta$, with $\beta < 1$ [11,12,15,16]. Two plausible physical pictures have been proposed to explain subdiffusion of hydration water [17–19]: spatial disorder, i.e., the rough surface of the biomacromolecules forms a fractal, percolated network, to hinder the diffusion of water molecules; or temporal disorder, i.e., water molecules jump between traps on the surface of the biomacromolecules with a broad distribution of trapping times. Because of a lack of microscopic evidence, however, the precise physical mechanism remains largely unclear [17,19].

In this work, we studied the diffusive dynamics of surface water on hydrated perdeuterated green fluorescent protein (GFP) and cytochrome P450 (CYP) at physiological conditions by combining neutron scattering experiments with molecular dynamics (MD) simulations and analytic modeling. We showed that the dynamics of hydration water is subdiffusive and the subdiffusion exponent gradually increases with observation time. By analyzing the MD trajectories of individual water molecules, we directly observed the discrete trapping events of the water molecules on the protein surface and found that the associated trapping times obey a broad distribution, leading to the subdiffusion. Moreover, the deep trapping sites are mostly occupied, and thus water molecules prefer to jump among shallow traps, rendering a gradual increase of subdiffusive power law with the observation time. Finally, a lattice toy model, many-body continuous time-random walk, is developed to provide a complete, quantitative description of all the relevant features of diffusive dynamics of the hydration water.

Results and discussions.—Neutron scattering directly probes the fluctuation of the nuclear position and is highly sensitive to hydrogen atoms. Hence, it is a powerful tool for studies of diffusion of water in a variety of environments, from porous silica to carbon nanotube or the surface of biomacromolecules [11,16,20,21]. Here, the neutron scattering experiments were performed at 280 K on both perdeuterated GFP and CYP [see Figs. 1(a) and 1(b)] at a hydration level of 0.4 g water/g protein, using the backscattering spectrometer (BASIS) at Oak Ridge National

(54) Mapping Twisted Light into and out of a Photonic Chip, Physical Review Letters

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Mapping Twisted Light into and out of a Photonic Chip

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Twisted light carrying orbital angular momentum (OAM) provides an additional degree of freedom for wireless optics and an emerging resource for both classical and quantum information technologies. Its inherently infinite dimensions can potentially be exploited by using mode multiplexing to enhance data capacity for sustaining the unprecedented growth in big data and internet traffic and can be employed to build large-scale quantum computing machines in high-dimensional Hilbert space. While the emission of twisted light from the surface of integrated devices to free space has been widely investigated, the transmission and processing inside a photonic chip remains to be addressed. Here, we present the first laser-driven silicon waveguide being capable of supporting OAM modes and experimentally demonstrate faithful mapping of twisted light into and out of a photonic chip. The states OAM_0 , $OAM_{\pm 1}$, $OAM_{\pm 2}$ and their superpositions can transmit through the photonic chip with a total efficiency up to 60% with minimal insertion loss. In addition, we present the transmission of quantum twisted light state of single photons and measure the output states with single-photon imaging. Our results reveal OAM as a new degree of freedom to be harnessed and manipulated in a photonic chip for high-capacity communication and high-dimensional quantum information processing.

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The phase of an optical beam with a spatial degree of freedom of orbital angular momentum (OAM) is twisted like a helix around its axis of travel and the cancellation of light waves at the axis itself results in a “donut-like” intensity profile. The twisted light has a helical wave front with an azimuthal phase term $e^{i\ell\phi}$ [1], with which every photon can carry an OAM of $\ell\hbar$ (ℓ is topological charge, ϕ is azimuthal angle, and \hbar is the Planck constant h divided by 2π). Having the special feature of intensity structure (donut-like intensity), phase structure (optical phase front), and dynamic characteristic (carrying OAM), the twisted light has been widely applied into the field of optical manipulation [2], optical mapping [3,4], and optical tweezers [5].

In recent years, OAM has shown great potential in communication systems for increasing the channel capacity capacity [6,7]. The inherent topological charges and the inherent orthogonality may provide tremendous resources for mode multiplexing. The inherently infinite dimensions can also potentially be exploited to deliver high-dimensional quantum states with larger alphabets and to build quantum computing machines in high-dimensional Hilbert space [8–12].

Large-scale applications of OAM beyond proof-of-principle demonstrations require developing integrated devices to enable the generation, transmission, and processing of such a new degree of freedom. Previous works have demonstrated on-chip generation twisted light with integrated star couplers [13], microring resonators [14], and controlled phase arrays [15]. While the emission of twisted light from the surface of integrated devices to free space has been widely investigated, the transmission and processing inside a photonic chip remain to be solved.

In this Letter, we demonstrate a faithful mapping of twisted light into and out of a photonic chip by prototyping dielectric waveguides with femtosecond laser direct writing [16,17]. We couple the states OAM_0 , $OAM_{\pm 1}$, $OAM_{\pm 2}$ and their superpositions into and out of the photonic chip with a total efficiency up to 60% and verify the output states by measuring with the Gaussian reference beam and making projection measurements, which clearly demonstrate that the output state (locally) keeps the OAM of the input state. In addition, we present the transmission of single-photon quantum twisted light and measure the output states with single-photon imaging.

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(55) Multistage Coupling of Laser-Wakefield Accelerators with Curved Plasma Channels, Physical Review Letters

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Multistage Coupling of Laser-Wakefield Accelerators with Curved Plasma Channels

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Multistage coupling of laser-wakefield accelerators is essential to overcome laser energy depletion for high-energy applications such as 100-GeV electron-positron colliders. Current staging schemes feed subsequent laser pulses into stages using plasma mirrors while controlling electron beams focusing with plasma lenses. Here a more compact and efficient scheme is proposed to realize the simultaneous coupling of the electron beam and the laser pulse into a second stage. A partly curved channel, integrating a straight acceleration stage with a curved transition segment, is used to guide a fresh laser pulse into a subsequent straight channel while the electron continues straight. This scheme benefits from a shorter coupling distance and continuous guiding of the electrons in plasma while suppressing transverse beam dispersion. Particle-in-cell simulations demonstrate that the electron beam from a previous stage can be efficiently injected into a subsequent stage for further acceleration while maintaining high capture efficiency, stability, and beam quality.

PRL 120, 154801 (2018)

Laser-wakefield accelerators (LWFAs) have recently attracted considerable attention as a promising new accelerator technology [1–4]. They are capable of supporting enormous acceleration gradients, as high as hundreds of GeV/m. This makes it possible to build compact accelerators in university-scale laboratories for many applications, such as compact electron diffraction devices [5], high-energy particle accelerators [6–9], and tabletop radiation sources [10–13].

Among the many applications, perhaps the most intriguing and challenging application is a LWFA-based TeV-level electron-positron linear collider [14,15]. However, the energy gained by electrons in a single-stage LWFA is limited by several effects, including electron dephasing, laser diffraction, and laser energy depletion. Although a few schemes have been proposed to increase the single-stage energy gain by mitigating dephasing and diffraction, a single-stage LWFA is still limited by pump depletion. Currently, single-stage acceleration up to 10 GeV energy is thought to be a reasonable value, given present laser technology and plasma density scaling, and a worldwide effort is pursuing this goal [8,16–21]. For future TeV colliders, the multistage coupling of LWFAs is inevitable. Electron beams accelerated in a first stage should be injected into a second wakefield stage that is driven by a

fresh laser pulse. Because of the micrometer-size electron beams and wake structures, the synchronization precision at the femtosecond scale, and the limited coupling distance, multistage coupling is considered to be very challenging. Recently, staged acceleration has been demonstrated by Benke *et al.* [22]. By using a plasma mirror [23] to reflect a fresh laser pulse and a plasma lens [24] to refocus electrons, about 3.5% of the electron beam charge was coupled into a second stage, which produced 100 MeV energy gain. In this staging scheme, the dedicated plasma mirror and the lens had to be installed between the LWFA stages. Besides its complexity, the matching of the electron beam between stages is still very challenging, particularly in achieving efficient coupling. The coupling efficiency must be near 100% if one considers the requirements for up to a hundred stages. Thus, a simple and efficient multistage coupling scheme is highly desirable.

Besides plasma mirrors, bending plasma channels may be considered to guide lasers [25–28]. Kruusa [26] theoretically studied laser propagation in curved plasma channels and found an equilibrium laser centroid trajectory. Chen [27] and Polino [28] discussed applications of synchrotron radiation based on curved plasma channels, and one patent was published on compact undulators and a radiation source based on curved channels [29]. In this

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(56) Translational and Rotational Dynamical Heterogeneities in Granular Systems, Physical Review Letters

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Translational and Rotational Dynamical Heterogeneities in Granular Systems

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We use x-ray tomography to investigate the translational and rotational dynamical heterogeneities of a three-dimensional hard-spherical granular packing driven by oscillatory shear. We find that particles which translate quickly form clusters with a size distribution given by a power law with an exponent that is independent of the strain amplitude. Identical behavior is found for particles that are translating slowly, rotating quickly, or rotating slowly. The geometrical properties of these four different types of clusters are the same as those of random clusters. Different cluster types are considerably correlated or anticorrelated, indicating a significant coupling between translational and rotational degrees of freedom. Surprisingly, these clusters are formed already at time scales that are much shorter than the relaxation time, in stark contrast to the behavior found in glass-forming systems.

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The relaxation dynamics of most disordered materials, such as glass-forming liquids, polymers, foams, and granular materials, differs significantly from the Debye behavior found in simple liquids in that it shows a marked non-exponential time dependence [1,2]. Although several mechanisms can give rise to this type of time dependence, it is often the local disorder of the particle arrangement that is the origin for this behavior [3–5], i.e., each particle has a different local environment, and as a result, the relaxation dynamics depends strongly on the particle considered. Previous studies have shown that the slowly (or quickly) relaxing particles are not distributed uniformly in space but instead form clusters. This effect, named “dynamical heterogeneity” (DH) is believed to be a key ingredient in understanding the *α*-process of glass-forming systems and hence the phenomenon of the glass transition. As a consequence, many studies have been carried out to study the nature of the DHs [3–12].

Usually DH are associated with the translational degrees of freedom (TDOF) of the particles. The particles in molecular systems and granular materials have, however, also rotational degrees of freedom (RDOF). Since these are coupled with the TDOF it is evident that they will be important for the relaxation dynamics of the system as well [13–17]. However, in practice it is difficult to probe the RDOF in molecular systems since experiments do not allow us to track directly the orientation of individual particles. As a consequence, only indirect experimental probing of the RDOF has been possible so far [7,13] and

most of our current knowledge of them comes from computer simulations [18–20]. (Plastic crystals are an exception since they have no disorder in the TDOF [21].) The situation is not much better for the case of granular materials since these are usually opaque and hence it is very challenging to probe in three dimensions (3D) the displacement and reorientation of the particles [22–24]. Because of the non-spherical shape of the particles and the presence of friction, there is often a strong coupling between the TDOF and RDOF making the experimental study of the DH for both TDOF and RDOF indispensable for understanding of the relaxation dynamics [25]. In the present work we thus use x-ray tomography to investigate these DH in a three-dimensional granular packing driven by oscillatory shear, making it, to the best of our knowledge, the first experimental investigation to probe the DH of the TDOF as well as the ones for the RDOF.

Our system consists of 4100 hard-rod particles made of polystyryl chloride with an aspect ratio of 1.5 (polydispersity 0.9%), i.e., a shape that makes the crystallization of the system difficult [26] and allows for a rather strong *F-R* coupling. The dimension of the minor axis of the particles is $2b = 12.7$ nm, and in the following we will use b as the unit of length. The particles are in a rectangular box of dimension $8b/2b = 8.5b = 22.6b$ that can be divided in the y direction. More details of the setup are given in [13]. We drive the system to a steady state by cycling it many times at all strain amplitudes γ investigated ($\gamma = 0.26, 0.40, 0.60$ and 0.87 ; see Supplemental Material [12] for details).

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(57) Broadband THz to NIR up-converter for photon type THz imaging, Nature Communications



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Broadband THz to NIR up-converter for photon-type THz imaging

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High performance terahertz imaging devices have drawn wide attention due to their significant application in healthcare, security of food and medicine, and nondestructive inspection, as well as national security applications. Here we demonstrate a broadband terahertz photon-type up-conversion imaging device, operating around the liquid helium temperature, based on the gallium arsenide homojunction interfacial workfunction internal photoemission (HIWIP)-detector-LED up-converter and silicon CCD. Such an imaging device achieves broadband response in 4.2–20 THz and can absorb the normal incident light. The peak responsivity is 0.5 AW^{-1} . The light emitting diode leads to a 72.5% external quantum efficiency improvement compared with the one widely used in conventional up-conversion devices. A peak up-conversion efficiency of 1.14×10^{-2} is realized and the optimal noise equivalent power is $29.1 \text{ pWHz}^{-1/2}$. The up-conversion imaging for a 1000 K blackbody pinhole is demonstrated. This work provides a different imaging scheme in the terahertz band.

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(58) Direct Observation of Topology from Single-Photon Dynamics, Physical Review Letters

PHYSICAL REVIEW LETTERS 122, 193903 (2019)

Direct Observation of Topology from Single-Photon Dynamics

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Topology manifesting in many branches of physics deepens our understanding on state of matter. Topological photonics has recently become a rapidly growing field since artificial photonic structures can be well designed and constructed to support topological states, especially a promising large-scale implementation of these states using photonic chips. Meanwhile, due to the inapplicability of Hall conductance to photons, it is still an elusive problem to directly measure the integer topological invariants and topological phase transitions in photonic system. Here, we present a direct observation of topological winding numbers by using bulk-state photon dynamics on a chip. Furthermore, we for the first time experimentally observe the topological phase transition points via single-photon dynamics. The integrated topological structures' direct measurement in the single-photon regime and strong robustness against disorder add the key element into the toolbox of "quantum topological photonics" and may enable topologically protected quantum information processing in large scale.

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The introduction of topology into condensed-matter and material sciences originates from the connection of integer quantum Hall conductances with topological Chern invariants [1], which greatly expands our knowledge on state of matter. With the birth of topological insulators, searching topological state of matter in solid state materials [2,3] and photonic systems [4,5] has recently become a leading research field. In contrast to the challenging experimental requirements for realizing topological states in solid state materials, photonic systems provide a convenient and versatile platform to design various topological lattice models and study different topological states including topological insulator states [6–10] and topological Weyl points [11,12]. The found topological edge states and phase transition potentially can be utilized for developing inherently robust and efficient artificial photonic devices [13–20].

Although photonic topological edge states have been extensively studied, the photonic topological invariant defined on the bulk states, which is the essential observable for characterizing topological states, is still less explored. In fermion systems, the topological invariant can be revealed by conductance measurements, while the concept of Hall conductance is inapplicable in photonic systems. Several methods recently have been proposed for indirectly

detecting topological invariants, based on probing Berry curvature [21,22], non-Hermitian photon loss [23,24], bulk-edge correspondence [25,26], and phase spectroscopy [27]. Nevertheless, direct measuring photonic topological invariants remains elusive.

Waveguide lattice provides a natural platform by studying photon dynamics. The reason is that the propagation of photons in a photonic waveguide lattice is naturally equivalent to a time evolution. However, direct measuring topological invariants based on photon dynamics has never been reported before in topological waveguide lattices. Here, for the first time, we experimentally demonstrate that both the topological invariants and the topological phase transitions can be directly measured based on the propagation of photons in the bulk of a topological waveguide lattice. Our work generalizes the method proposed in linear optics-based quantum walks [28,29] to use continuous-time dynamics to directly measure topology. This dynamical method is very generic for observing topology, which recently also has been experimentally demonstrated in ultracold atoms [30].

To extend promised topological protection into the quantum regime, we have to find an appropriate system that is physically scalable and has inherently low loss when

(59) Parity-Induced Thermalization Gap in Disordered Ring Lattices, Physical Review Letters

PHYSICAL REVIEW LETTERS 122, 013903 (2019)

Parity-Induced Thermalization Gap in Disordered Ring Lattices

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The gaps separating two different states widely exist in various physical systems: from the electrons in periodic lattices to the analogs in plasmas, phonons, plasmons, systems, and even quasicrystals. Recently, a thermalization gap, an inaccessible range of photon statistics, was proposed for light in disordered structures [Nat. Phys. 11, 920 (2015)], which is intrinsically induced by the disorder-immune chiral symmetry and can be reflected by the photon statistics. The lattice topology was further identified as a decisive role in determining the photon statistics when the chiral symmetry is satisfied. Being very distinct from one-dimensional lattices, the photon statistics in ring lattices are dictated by its parity, i.e., odd or even sized. Here, we for the first time experimentally observe a parity-induced thermalization gap in strongly disordered ring photonic structures. In a limited scale, though the light tends to be localized, we are still able to find clear evidence of the parity-dependent disorder-immune chiral symmetry and the resulting thermalization gap by measuring photon statistics. While strong disorder-induced Anderson localization overwhelms such a phenomenon in larger-scale structures. Our results shed new light on the relation among symmetry, disorder, and localization, and may inspire new resources and artificial devices for information processing and quantum control in a photonic chip.

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Symmetry is crucial for establishing energy gaps for electrons in lattices and analogs [1–5]. Disorder often decreases such effects [1], but not for all cases. Certain disorder-immune chiral symmetries [6–9] would emerge in random matrix theory [10], such as chiral [6] and particle-hole symmetric ensembles [11], and play a key role in fields ranging from superconductivity [12] to quantum chromodynamics [13]. A distinctive characteristic of disorder-immune chiral symmetry is that the system Hamiltonian can be transformed into a block off-diagonal matrix form, a separate bipartite sublattice [10,14,15], and is disorder-immune behaving as holding for each realization from a disordered ensemble [16].

The light-binding lattices have been suggested to study disorder-immune chiral symmetry [16]. The resulting thermalization gap, set by parity, is circuitously confirmed in one-dimensional disordered lattices [17,18]. Inspired by the works on photonic thermalization gap [16] and the effect of lattice topology [19], we investigate the scenario of parity-dependent disorder-immune chiral symmetry in ring lattices, chiral symmetry breaking for even-sized ring lattices, and chiral symmetry breaking for odd-sized ones. The resulting thermalization gap can be

revealed by observing the photon statistics in cycle-parity disordered ring lattices.

In this Letter, we demonstrate experimental observation of the parity-induced thermalization (PIT) gap on a photonic chip. By using femtosecond laser direct writing, we are able to freely prototype waveguides as sites of lattices and introduce disorder by precise coupling control in three dimensions. We measure the evolved light distribution out of 120 lattices with randomly picked coupling parameters in a fixed disorder level to obtain photon statistics. We observe a thermalization gap in an even-parity disordered ring lattice where the cross-correlation cannot go below a certain limit due to disorder-immune symmetries. As a way to demonstrate such a parity-induced thermalization gap, we observe that the cross-correlation in even-sized lattices is significantly larger than the one in odd-sized lattices. This work may deepen the understanding of the relation among symmetry, disorder, and localization, and inspire applications for quantum integrated photonics.

The dynamics of light in photonic lattices can be described by a set of coupled discrete Schrödinger equations [20], which are derived from the Schrödinger-type paraxial

(60) Superhigh-Resolution Recognition of Optical Vortex Modes Assisted by a Deep-Learning Method, Physical Review Letters

PHYSICAL REVIEW LETTERS 123, 183902 (2019)

Superhigh-Resolution Recognition of Optical Vortex Modes Assisted by a Deep-Learning Method

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Orbital angular momentum (OAM) has demonstrated great success in the optical communication field, which theoretically allows an infinite increase of the transmitted capacity. The resolution of a receiver to precisely recognize OAM modes is crucial to expand the communication capacity. Here, we propose a deep-learning (DL) method to precisely recognize OAM modes with fractional topological charges. The minimum interval recognized between adjacent modes decreases to 0.01, which is far as we know is the first time this superhigh resolution has been realized. To exhibit its efficiency in the optical communication practice, we transfer an Ethernet protocol by a superhigh-resolution OAM multiplexing system. As the convolutional neuron networks can be trained by data up to an infinitely large volume in theory, this work exhibits a huge potential of generalization capability for next-generation DL-based ultrahigh OAM optical communication, which might even be applied to microwave, millimeter wave, and terahertz OAM communication systems.

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Vortex beams carrying orbital angular momentum (OAM) have been extensively investigated in optical manipulation [1], imaging [2], interaction between light and matter [3,4], and optical communication [5] since it was recognized in 1992 [6]. The helical wave front of such a vortex beam is described by a phase factor $\exp(i\ell\phi)$, where ϕ is the azimuthal angle and the topological charge ℓ is an unlimited integer or fractional value. Owing to the unbounded dimensional space, it provides high degrees of freedom for multiplexing information, which infinitely boosts the transmitted capacity in optical communication processes [7,8].

Because of its significant application, the generation of OAM states has aroused tremendous enthusiasm ranging from spiral phase plate [9], q plate [10,11], and metasurface [12] to integrated devices [13]. One of the most accomplished types of technology is to use spiral light modulator (SLM) uploading spiral holograms [14], which can simply create vortex beams and flexibly manipulate the phase information of light. Besides, the recognizable ability of the receiver is also crucial for both OAM shift keying (OAM-SK) [15] and OAM division multiplexing (OAM-DM) [16] in terms of OAM-based optical communication. Therefore, much work has been conducted focusing on efficiently developing an OAM sorter recently. Some techniques utilize traditional optics devices that transform the azimuthal position into transverse position [17] based on interferometers [18–20] and vortex diffractive gratings [21]. One currently efficient sorting method is to implement coordinate transformation that separates states with different topological charges by a distinct focal spot on the detect-

plane [22,23]. And a computational method to measure the transmission matrix (TM) provides an anti-interference approach to reverse the propagation of vortex beams [24]. In addition, the OAM-to-polarization coupling effect is also proposed to sort different OAM states of light [25]. All these methods pave the way for effective OAM-based optical communication.

However, owing to the limitation of the resolution, only eigenmodes whose topological charges are integer are considered in all methods mentioned above. With the increase of an integer topological charge value, the growing phase singularity and the diffraction effect enormously affects the intensity distributions of vortex beams, which extremely adds the difficulties of being focused in free-space and being coupled in fibers. This problem limits the development of the OAM-based optical communication. Consequently, it is of fundamental importance to expand communication capacity adequately with more OAM states but smaller phase singularity.

Our goal here is to develop a superhigh-resolution technique to precisely separate modes into subdivisible space between adjacent eigenmodes. That is, the minimum interval $\Delta\ell$ among recognized modes is a fractional value and can be as small as possible. The pioneering techniques of generating fractional OAM modes were proposed in 2004 [26,27]. However, the tiny variety of fractional OAM modes is further complicated to be discriminated compared to that of integer charges [28–30]. Recently, the deep learning (DL) method has been developed to possess the ability of extracting intrinsic features and dividing decision boundary according to data [31]. The original assignment

(61) Coupling of Edge States and Topological Bragg Solitons, Physical Review Letters

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Coupling of Edge States and Topological Bragg Solitons

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The existence of the edge states at the interface between two media with different topological properties is protected by symmetry, which makes such states robust against structural defects or disorder. We show that if a system supports more than one topological edge state at the interface, even a weak periodic perturbation may scatter one edge state into another without coupling to both modes. Due to the Bragg scattering of the edge modes, which in a topological system is highly selective, with closed bulk and backward scattering channels, even when conditions for resonant scattering are not satisfied. When such a system bears nonlinearity, Bragg scattering enables the formation of a new type of soliton—topological Bragg solitons. We report them in a spin-orbit coupled (SOC) Bose-Einstein condensate in a homogeneous two-particle Zeeman lattice. An interface supporting two edge states is created by two different SOC's, with the s (representing the synthetic magnetic field) having opposite directions on different sides of the interface. The reported Bragg solitons are found to be stable.

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Topological edge states are fundamental for understanding the physics of many phenomena including quantum Hall effects [1], anomalous [1], and quantum spin Hall effects; topological insulators [2], Majorana fermions [3], just to mention a few. Introduced in solid state physics, topological edge states were shown to be a universal wave phenomenon. In particular, they enable protected edge currents in photonic crystals as predicted in [4] and observed in [5–14] (see review [15]), in surface plasmon-polariton systems [16], in systems of cold atoms in optical lattices [14–16], in spin-orbit coupled Bose-Einstein condensates (SOC BECs) [17], and in optoelectronic systems, such as cyclotron-polariton condensates [18,19].

Topological edge states are robust against disorder [6], which distinguishes them from bulk states that can be manipulated by perturbations [20]. This makes such states promising for a variety of applications. Meanwhile, when it is necessary to selectively excite or remove edge states, or transform them in any other way, this robustness becomes a drawback. Topological edge states can be coupled by nonlinearity that gives rise to a rich set of phenomena, such as modulational instability [10] or envelope soliton propagating along the edge [17]. However, by enabling linear coupling one qualitatively enriches the physics of respective systems. Now nonlinear interactions, requiring simultaneous energy and momentum conservation laws, can be made resonant. That leads to a plethora of novel phenomena, which so far were not considered for edge states. In particular, linear mode coupling resulting from

periodic modulation of the system, i.e., from Bragg scattering, in the presence of nonlinearity, can lead to the formation of Bragg solitons [21,22], which are relevant for many applications [23], and whose properties qualitatively differ from those of the envelope solitons mentioned above.

In this Letter, we introduce an efficient mechanism of coupling and conversion of topological edge states based on Bragg scattering by periodic modulations of an interface between topologically different media. This mechanism works when a system supports more than one topological edge state per interface. We show that even a very weak periodic perturbation of such an interface may result in periodic transitions between two edge states moving with different group velocities. Such a coupling is highly selective, with closed bulk and backward scattering channels. Furthermore, we construct Bragg solitons propagating along the interface and representing a spatially localized envelope of two coupled edge states.

As a case example, we address an atomic SOC BEC [24,25], characterized by a spin-orbit parameter $\mathcal{M} = (\mathcal{M}^T, \mathcal{M}^T)^T$ (T stands for the transpose). The condensate is placed in a honeycomb lattice, which can be created experimentally by applying at least three laser beams [26,27]. Almost arbitrary 2D field distributions, and thus optical potentials, can be produced by the interference of quasi-parallel laser beams [28]. The lattice is characterized by inverted potential profiles for the spin components and can be created by periodically varying Zeeman splitting (see [29] for a possibility of experimental

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Research in Physics

Terahertz Oscilloscope for Recording Time Information of Ultrashort Electron Beams

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We propose and demonstrate a terahertz (THz) oscilloscope for recording time information of an ultrashort electron beam. By injecting a laser-driven THz pulse with circular polarization into a dielectric tube, the electron beam is swept radially such that the time information is uniformly encoded into the angular distribution that allows one to characterize both the temporal profile and timing jitter of an electron beam. The dynamic range of the measurement in such a configuration is significantly increased compared to detection with a linearly polarized THz pulse. With this THz oscilloscope, nearly 50-fold longitudinal compression of a relativistic electron beam is about 15 fs (rms) is directly visualized with its arrival time determined with 1 fs accuracy. This technique bridges the gap between streaking of photoelectrons with optical lasers and deflection of relativistic electron beams with radio-frequency deflectors, and should have wide applications in many ultrashort electron-beam-based facilities.

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The ability to characterize the time information of an ultrashort electron beam including both the temporal profile and arrival time is crucial for optimizing and enhancing the performance of many electron-beam-based scientific facilities such as free-electron lasers (FELs) [1–3], ultrafast electron diffraction (UED) [4,5], and microscopy (UEM) [6–9], laser-driven and beam-driven advanced accelerators [10–15], etc. In the accelerator community, rf deflecting cavities have been widely used to measure the temporal profile of relativistic electron beams with energy ranging from MeV to GeV (see, e.g., Refs. [16–18]). However, the information of beam arrival time with respect to an external laser as required in a pump-probe experiment cannot be directly measured with a rf deflector. In the attosecond-science community, streaking of photoelectrons with optical lasers has become a standard technique for characterizing the complete information of attosecond pulses [19]. Recently, this technique has been adapted to characterize femtosecond x-ray pulses in FELs with the streaking imparted by femtosecond and THz pulses [20–23]. However, this technique

does not apply to a relativistic electron beam as dictated by the Lawson-Woodward theorem [24]. Very recently, THz-streaking of keV and MeV electrons [25–27] in a subwavelength metallic structure has been used to measure both the temporal profile and the arrival time of electron beams. However, the small aperture used to enhance the THz field may significantly limit the number of useful electrons. Furthermore, with the streaking imparted by a linearly polarized THz pulse, the beam receives (anisotropic) angular streaking, and thus the measurement has a rather limited dynamic range (time window where the measurement is accurate) comparable to about one-quarter of the wavelength.

In this Letter, we demonstrate a laser-driven THz oscilloscope that allows one to record the complete time information of an ultrashort electron beam with both large dynamic range and high temporal resolution. This scheme exploits the assumed interaction between the THz pulse and the electron beam in a dielectric tube, allowing sufficient deflection to be generated without enhancing the THz field with a narrow slit. With accurate control of the THz polarization, the electron

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Connecting shear localization with the long-range correlated polarized stress fields in granular materials

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One long-lasting puzzle in amorphous solids is shear localization, where local plastic deformation involves cooperative particle rearrangements in small regions of a few inter-particle distances, self-organizing into shear bands and eventually leading to the material failure. Understanding the connection between the structure and dynamics of amorphous solids is essential in physics, material sciences, geotechnical and civil engineering, and geophysics. Here we show a deep connection between shear localization and the intrinsic structures of internal stresses in an isotropically jammed granular material subject to shear. Specifically, we find strong (anti)correlations between the micro shear bands and two polarized stress fields along two directions of maximal shear. By exploring the tensorial characteristics and the rotational symmetry of force network, we reveal that such profound connection is a result of symmetry breaking by shear. Finally, we provide the solid experimental evidence of long-range correlated inherent shear stress in an isotropically jammed granular system.

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Designer spin order in diradical nanographenes

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The magnetic properties of carbon materials are at present the focus of intense research effort in physics, chemistry and materials science due to their potential applications in spintronics and quantum computing. Although the presence of spins in open-shell nanographenes has recently been confirmed, the ability to control magnetic coupling sign has remained elusive but highly desirable. Here, we demonstrate an effective approach of engineering magnetic ground states in atomically precise open-shell bipartite/nonbipartite nanographenes using combined scanning probe techniques and mean-field Hubbard model calculations. The magnetic coupling sign between two spins was controlled via breaking bipartite lattice symmetry of nanographenes. In addition, the exchange-interaction strength between two spins has been widely tuned by finely tailoring their spin density overlap, realizing a large exchange-interaction strength of 42 meV. Our demonstrated method provides ample opportunities for designer above-room-temperature magnetic phases and functionalities in graphene nanomaterials.

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(65) Optical soliton formation controlled by angle twisting in photonic moiré lattices, Nature Photonics



Optical soliton formation controlled by angle twisting in photonic moiré lattices

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Exploration of the impact of synthetic material realizations featuring tunable geometrical properties on physical processes is a research direction that is currently of great interest because of the outstanding phenomena that are continually being uncovered. Twistronics and the properties of wave excitations in moiré lattices are salient examples. Moiré patterns bridge the gap between aperiodic structures and perfect crystals, thus opening the door to the exploration of effects accompanying the transition from commensurate to incommensurate phases. Moiré patterns have revealed profound effects in graphene-based systems¹, they are used to manipulate ultracold atoms² and to create gauge potentials³ and are observed in colloidal clusters⁴. Recently, it was shown that photonic moiré lattices enable observation of the two-dimensional localization-to-delocalization transition of light in purely linear systems⁵. Here, we employ moiré lattices optically induced in photorefractive nonlinear media^{6–8} to elucidate the formation of optical solitons under different geometrical conditions controlled by the twisting angle between the constitutive sublattices. We observe the formation of solitons in lattices that smoothly transition from fully periodic geometries to aperiodic ones, with threshold properties that are a pristine direct manifestation of flat-band physics⁹.

Broadly, large-scale linear transport and localization properties of electrons in a material are intimately determined by its inner symmetry and geometrical properties, including its periodic or aperiodic nature, its electronic¹⁰, atomic¹¹, optical¹² or two-dimensional (2D) material¹³ systems. The lattice is directly related to the nature of the transport of the system, which can be revealed or realized. When an underlying material exhibits a nonlinear response, the formation of localized excitations—such as solitons—becomes possible^{14–16}. Nevertheless, the properties of solitons are still strongly impacted by the basic symmetries of the system. Optical induction is a unique laboratory for the investigation of solitons in a diverse environment. Thus, 2D self-trapping and soliton formation have been investigated in fully periodic optical lattices^{17–20} or well-defined aperiodic photonic structures^{21–23} by means of numerical simulations^{24–26}. A common feature of all these results is the occurrence of a power threshold for 2D soliton excitation, the value of which can be tuned by varying the lattice configuration. However, the threshold excitation in the case of aperiodic lattices or the transition from aperiodic to periodic lattices has been barely reported previously^{27,28}.

In this Letter, we study a novel optical moiré lattice—after a potential platform realization with a study coupled with incommensurate superlattices, moiré patterns enable the localization of light over a wide range from disordered lattices to a crystal. The bands in their resulting geometry that strongly impacts the diffraction behavior such as bands. Thus, the same incommensurate energy due to a balance between diffraction and self-phase modulation induced by nonlinearity, more lattices allow investigation of the formation of solitons controlled by the geometry of the induced optical potential. Here, we provide experimental evidence of such a possibility by reporting qualitative differences in soliton excitation dynamics in commensurate and incommensurate moiré lattices. We present an observation consistent with the formation of threshold 2D optical solitons in an aperiodic system, which is enabled by the presence of extremely flat bands in the incommensurate geometry. In addition—and importantly—we demonstrate that in commensurate moiré lattices the threshold for soliton formation is directed by the angle of the moiré superlattice, which determines the width of the allowed bands. Our observations thus show that the properties of flat bands are not only determined by the characteristics of the respective gaps but also fundamentally depend on the properties of the allowed bands.

A photonic moiré lattice can be created in a photorefractive crystal by a shallow modulation of the refractive index in the xy plane induced by two mutually rotated or twisted periodic square lattices generated by light interference along tilted crystal lattices. Incommensurate lattices induced, where the linear dielectric effect determines the refractive index modulation experienced by light with different polarizations. The open space coefficients are 45 pm^{-2} in magnitude, that is, $\epsilon_{ij} = 2.87 \text{ pm}^{-2}$. Accordingly, we used a linear polarized light in the lattice induction so that the corresponding beam did not experience any noticeable self-action in the crystal and this process was conducted as a linear medium. In contrast, light with self-induced polarization experienced a nonlinear medium response. Preparation using the experimental layout leads to such polarizations is achieved by the nonlinear self-action operation for $\Delta n = \epsilon_{ij} \cos(\theta) \sin(\theta) \cos(2\theta) \cos(2\theta)$.

$$\frac{\partial \Psi}{\partial t} = \frac{1}{2} V_x \Psi + \frac{V_y}{1 + i(\Gamma - \Psi)} \Psi \quad (1)$$

Here, Ψ is optical field, V_x and V_y are coupling coefficients, Γ is the nonlinear refractive index coefficient induced by the observation of

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Article

Localization and delocalization of light in photonic moiré lattices

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Moiré lattices consist of two superimposed identical periodic structures with a relative rotation angle. Moiré lattices have several applications in everyday life, including artistic design, the textile industry, architecture, image processing, metrology and interferometry. For scientific studies, they have been produced using coupled graphene–hexagonal boron nitride monolayers^{1,2}, graphene–graphene layers^{3,4} and graphene quantum crystals on a silicon carbide surface⁵. The recent surge of interest in moiré lattices arises from the possibility of exploring many salient physical phenomena in such systems; examples include commensurate–incommensurate transitions and topological defects⁶, the emergence of insulating states owing to band flattening⁷, unconventional superconductivity⁸ controlled by the rotation angle⁹, the quantum Hall effect¹⁰, the realization of non-Abelian gauge potentials¹¹ and the appearance of quasicrystals at special rotation angles¹². A fundamental question that remains unexplored concerns the evolution of waves in the potentials defined by moiré lattices. Here we experimentally create two-dimensional photonic moiré lattices, which—unlike their material counterparts—have readily controllable parameters and symmetry, allowing us to explore transitions between structures with fundamentally different geometries (periodic, general aperiodic and quasicrystal). We observe localization of light in deterministic linear lattices (that is based on flat-band physics¹³), in contrast to previous schemes based on light diffusion in optical quasicrystals¹⁴, where disorder is required¹⁵ for the onset of Anderson localization¹⁶ (that is, wave localization in random media). Using commensurate and incommensurate moiré patterns, we experimentally demonstrate the two-dimensional localization–delocalization transition of light. Moiré lattices may feature an almost arbitrary geometry that is consistent with the crystallographic symmetry groups of the substances, and therefore afford a powerful tool for controlling the properties of light patterns and exploring the physics of periodic–aperiodic phase transitions and two-dimensional wavepacket phenomena relevant to several areas of science, including optics, acoustics, condensed matter and atomic physics.

One of the most salient properties of an engineered optical system is its capability to affect a light beam in a prescribed manner, such as to control its diffraction pattern or to localize it. The importance of wavepacket localization extends far beyond optics, and impacts all branches of science dealing with wave phenomena. Irregularities or strictly periodic linear systems cannot result in wave localization, and the latter require the presence of structure defects or nonlinearity. Anderson localization¹⁶ is a hallmark discovery in condensed matter physics. All electronic states in one- and two-dimensional potentials with topological disorder are localized. Three-dimensional systems with disordered potentials are known to have both localized and

delocalized eigenstates¹⁷, separated by an energy known as the mobility edge¹⁸. Coexistence of localized and delocalized eigenstates has been predicted also in regular quasiperiodic one-dimensional systems. First in the discrete Aubry–André¹⁹ model and later in continuous systems and matter wave systems^{20–22}. Quasiperiodic (or aperiodic) structures, even those that possess long-range order, fundamentally differ both from periodic systems, where all eigenstates are delocalized Bloch waves, and from disordered media, where all states are localized (in one or two dimensions). Upon variation of the parameters of a quasiperiodic system, it is possible to observe the transition between localized and delocalized states. Such a localization–delocalization transition (LDT)

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(67) Direct Visualizing the Spin Hall Effect of Light via Ultrahigh-Order Modes, Physical Review Letters

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Direct Visualizing the Spin Hall Effect of Light via Ultrahigh-Order Modes

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We report an experiment showing the ultrahigh-order Imbert-Fedorov shift from the ultrastrong spin-orbital angular momentum coupling, which is a photonic version of the spin Hall effect, by measuring the reflection of light from the surface of a birefringent symmetrical metal cladding planar waveguide. The light incident at a near-normal incident angle and excites resonant ultrahigh-order modes inside the waveguide. A 3D in-situ displacement of separated reflected light spots corresponding to two polarizations states is distinguishable by human eyes. In our experiment, we demonstrate the control of polarizations of light and the direct observation of the spin Hall effect of light, which opens an important avenue towards potential applications for optical sensing and quantum information processing, where the spin states of photons exhibit key features.

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The spin-orbit interaction of photons on the interface with the refractive index discontinuity corresponds to the interplay between the spin degree of freedom of light and the intrinsic orbital angular momentum, which is a photonic version of the spin Hall effect [1–3]. It leads to a perpendicular spin-dependent displacement of light, i.e., the so-called Imbert-Fedorov (IF) shift [4,5]. For many measurements of the spin Hall effect of light (SHEL), the displacement between two opposite spins is at the order of the wavelength of light, which cannot be distinguishable by human eyes. The quantum weak measurement technology [6–11] therefore has been utilized, which opens a promising experimental way and hence brings potential practical applications [9,11]. On the other hand, the general concept of SHEL associated with the polarization nature of light occurs universally in many typical systems [16]. It leads to the longitudinal Goos-Hänchen shift from the spatial dispersion of light and also the transverse IF shift [7,9], which originates from the geometric Berry phase and the conservation of normal components of the total angular momentum of light [10]. Because of the unique physical nature [5–9] as well as possible implications in the metrology and the spin-based nano-optics [10,11], the research of SHEL is of important interest [12–14]. In particular, a large number of experimental works with the quantum weak measurement have been performed to observe the IF shift at a variety of optical platforms including the air-dielectric interface [15–20], metamaterial with the photon tunneling process [12], metasurfaces [17,21–24], Weyl semimetals [25], heterostructure semiconductors [26], and thin films supporting leaky guided

modes [7,11], where the quantum state of light limits the efficiency of the measurement [11]. Recently, methods associated with the surface plasmon resonance (SPR) have been applied to enhance the SHEL [27–30]. For example, an enhancement with a factor of 50 is revealed for the light incident on the metal surface near the Brewster angle, which gives the IF shift $\sim 3.2 \mu\text{m}$ [30]. Moreover, a small incident angle for the enhancement of SHEL in hyperbolic metamaterials has been explored and the transverse beam shift is $\sim 9 \mu\text{m}$ [31]. A natural question arises, is there a system supporting an IF shift associated with the SHEL that is large enough to be distinguishable by human eyes?

In this Letter, we show an ultrastrong spin-orbital angular momentum coupling of light in the experiment by using the resonance of resonant ultrahigh-order modes (UOMs). A linear-polarized light is incident into the birefringent symmetrical metal cladding planar waveguide at the near-normal incident angle θ_i and excites a high density of UOMs inside the waveguide [32–35]. The experiment shows a transverse spatial deviation at a submillimeter scale in the reflected light corresponding to two opposite spins, which can be visibly seen on the CCD. Our work points out a unique platform for studying the fundamental physics of the ultrastrong spin-orbital angular momentum coupling of light without applying the quantum weak measurement, leading towards broad potential implications including quantum information, optical communications, and biosensing.

We start with a brief description on the physical picture of the spin-orbital angular momentum coupling of light by considering a plane wave incident on a surface. Under the

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Integrated Quantum-Walk Structure and NAND Tree on a Photonic Chip

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In the age of the post-Moore era, the next-generation computing model would be a hybrid architecture consisting of different physical components, such as photonic chips. In 2008, it was proposed that the solving of the NAND-tree problem can be sped up by quantum walk. This scheme is groundbreaking due to the universality of the NAND gate; however, experimental demonstration has not been achieved so far, mostly due to the challenge in preparing the propagating initial state. Here we propose an alternative solution by including a structure called a “quantum slide,” where a propagating Gaussian wave packet can be generated deterministically along a properly engineered chip. In our experimental demonstration, the optical NAND tree is capable of solving computational problems with a total of four input bits, based on the femtosecond laser 3D direct-writing technique on a photonic chip. These results overcome the main roadblock to photonic NAND-tree computation, and the construction of a quantum slide may find other interesting applications in quantum information and quantum optics.

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Quantum walk, the quantum generalization of the classical random walk, is a natural platform for discovering exotic quantum phenomena and developing quantum algorithms [1–8]. Quantum-walk-based schemes have provided speed-up for various problems of particular interest, including Boson sampling [9–12], black-box problem [13], element distinctness [14], binary addition [15], factoring integers [16], and machine learning [17]. It has also been shown that quantum walk is powerful enough to perform universal quantum computation [18,19].

A NAND tree is a binary tree of NAND gates containing a total of N inputs but only one output, enabling the computation of arbitrary Boolean function of the form $F: \{0, 1\}^N \rightarrow \{0, 1\}$. While there are many classical algorithms for the NAND-tree problem [20,21], this task can be trapped to the quantum walk to speed up the computation. In 2008, Farhi *et al.* proposed a continuous quantum-walk-based protocol for balanced NAND tree for the first time [22]. Inspired by it, several discrete quantum-walk schemes have been developed [23,24]. The schemes have also been generalized by the unbalanced NAND formula [25–27], which can, in principle, represent arbitrary Boolean functions [28].

The original scheme of the NAND tree in Ref. [22] is based on a continuous quantum walk on a graph. As shown in Fig. 1, the graph contains a chain of uniform coupled sites, called “runway,” the quantum NAND tree whose sites are connected in the structure as a binary tree, and a set of input nodes. Suppose the graph is represented by the adjacent matrix T , the time-independent Hamiltonian of the entire system is just $H = JG$, where J is the coupling constant. The input is encoded by the on-off coupling of the input nodes (leaves) and the nodes on the top layer of the tree. For an incident wave packet with zero energy in the runway [22], if the computation outcome is $F(x) = 1$, the wave packet will pass through the quantum NAND tree; otherwise, the wave packet will be reflected when $F(x) = 0$. Therefore, by measuring whether the wave packet has passed through or been reflected, the binary computation outcome can be obtained.

The original protocol [22] of the quantum NAND tree is elegant and simple, and there is also proposal of realizing it on the molecular platform [29] being developed recently. However, its experimental realization remains challenging and, to our best knowledge, has not been

(69) Subharmonic Entrainment Breather Solitons in Ultrafast Lasers, Physical Review Letters

PHYSICAL REVIEW LETTERS 125, 163901 (2020)

Subharmonic Entrainment Breather Solitons in Ultrafast Lasers

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We study the stability and experimentally the subharmonic entrainment (SHE) breather soliton in mode-locked lasers for the first time, in which the ratio of the breather period to the round-trip time is an integer. We build a non-Hermitian degeneracy map of breather solitons and illustrate that SHE arises between the two exceptional points (EPs). We obtain SHE at the cost of 2% increase the isolation of breather solitons when tuning the gain and/or cavity loss, and prove that this phenomenon can improve the stability of breather soliton. Our research brings insight into the EP physics of ultrafast lasers and makes the mode-locked laser a powerful test bed for non-Hermitian degeneracy, which may open a new avenue in ultrafast laser research.

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Solitons, self-localized coherent structures, are ubiquitous in nature and have been observed in most branches of nonlinear science [1–5]. Mode-locked lasers constitute an ideal test bed for investigating solitons [6], where solitons arise from the balance among dispersion, nonlinearity, and environmental energy exchange. Once given the medium parameters, solitons generally display a fixed profile over propagation [7–8], which enables the application in designing pulse sources [9]. In many systems, the amplitude and duration of solitons can oscillate or breathe periodically. Breather solitons were first demonstrated in fiber cavities [10] and then observed in optical microresonators [11, 12] and mode-locked lasers [13]. Because of strong connection with the Fermi-Pasta-Ulam recurrence [14], turbulence [15], modulation instability [16], and rogue wave formation [17], breather solitons have attracted considerable attention. Furthermore, the ability of breather soliton to increase the resolution of dual-comb source (composed of two frequency combs with slightly different repetition rates) implies that the breather soliton is attractive in practical applications [18].

Recently, the breather soliton generated in Kerr-resonant optical resonators was predicted to exhibit subharmonic entrainment (SHE): an integer ratio $T_b/T_r = n$ of the breather period T_b to the round-trip time T_r [19]. It is a physical frequency locking that occurs between the two exceptional points (EPs) [20]. The EP is a branch point singularity in the parameter space [21–23] that arises when two or more eigenvalues of the non-Hermitian Hamiltonian and their corresponding eigenvectors become degenerate. Such degeneracies are distinct features of non-Hermitian systems, which do not obey conservation laws because the systems exchange energy with environment. The breather soliton laser, which has two frequencies (the repetition frequency f_r and the

breather frequency f_b), can be deemed as a non-Hermitian system. The eigenvalues of the non-Hermitian Hamiltonian between the EPs are complex with the same real part [21], resulting in frequency locking in the breather soliton system, from which the SHE breather soliton arises. Although it was predicted in Kerr-nonlinear resonators [19], the SHE breather soliton has not been examined experimentally.

In this Letter, we present the first experimental observation of the SHE breather soliton. We find that the breather soliton laser exhibits non-Hermitian degeneracy in certain condition and that the SHE arises between the two EPs. Experimentally, we fabricate a mode-locked fiber laser, which is easily driven into the breather soliton regime by tuning the pump power and/or cavity loss. We record the change of breather soliton with the pump power and characterize SHE using dispersive Fourier transform (DFT) technique [24–26] and dispersive temporal interferometer (DTI) [27]. We prove that SHE can enhance the stability of breather soliton. Simulations also prove that SHE arises between the EPs.

In principle, a stable mode-locked laser only contains the modes at the free-running cavity frequency f_{00} (corresponding to the cavity round-trip time T_{00}) together with its harmonic (at the repetition frequency) of $n f_{00}$, n is an integer. Nevertheless, the breather soliton mode-locked laser (which is an unstable system, in which another mode, the free-running breather frequency f_{b0} , arises. Typically, $f_{b0} \approx f_{00}$) [19]. Figure 1(a) depicts the free propagation of f_{b0} , which shows slight difference between $n f_{00}$ and the free-running breather period T_{b0} by $\Delta = T_{b0} - n T_{00}$ with $|\Delta| \leq T_{00}/2$. In Fig. 1(b), the soliton is treated as repeating in a period $n T_{00}$ with uniformly distributed n solitons in each period. Each soliton appears at the free-running subharmonic frequency $f_{s0} = f_{b0}/n$

(70) Vector Vortex Beam Emitter Embedded in a Photonic Chip, Physical Review Letters

PHYSICAL REVIEW LETTERS 124, 153601 (2020)

Vector Vortex Beam Emitter Embedded in a Photonic Chip

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Vector vortex beams simultaneously carrying spin and orbital angular momentum of light promise additional degrees of freedom for modern optics and emerging scenarios for both classical and quantum information technologies. The inherently infinite dimensions can be exploited to enhance data capacity for sustaining the unprecedented growth in big data and internet traffic and can be employed to build quantum computing machines in high-dimensional Hilbert space. So far, much progress has been made in the emission of vector vortex beams from a chip surface into free space; however, the generation of vector vortex beams inside a photonic chip has not been realized yet. Here, we demonstrate the first vector vortex beam emitter embedded in a photonic chip by using femtosecond laser direct writing. We achieve a conversion of vector vortex beams with an efficiency up to 50% and scalar vortex beams with an efficiency up to 74% from Gaussian beams. We also present an expanded coupled-mode model for understanding the mode conversion and the influence of the imperfection in fabrication. The fusion of embedded generation makes vector vortex beams directly ready for further transmission, manipulation, and emission without any additional interconversion. Together with the ability to be integrated as an array, our results may enable vector vortex beams to become accessible inside a photonic chip for high-capacity communication and high-dimensional quantum information processing.

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Light can carry both spin and orbital angular momentum (OAM). Spin angular momentum is associated with optical polarization, which is one of the most prominent and well-known properties of light. OAM is an emerging degree of freedom, which has helical wavefronts described by an azimuthal phase term $e^{i\ell\phi}$ (ϕ is the azimuthal angle in a cylindrical coordinate system). The topological charge ℓ can take any integer value represented by the number of crossed spiral wavefronts when rotating around the axis once. Because of the admitted topological charges and the inherent orthogonality, OAM can provide larger alphabets for classical information [2–4] and quantum information processing (QIP) [5–11].

Recently, there has been an increasing interest in vector vortex light beams, with a polarization varying along the azimuthal coordinate and a central optical singularity. Such beams exhibit some unique characteristics such as field

structure, phase singularity and rotation invariance [12]. Vector vortex beams have been applied in wide areas including microscopy [13], optical trapping [14], precision measurement [15], quantum communication [12,16], and QIP [17–21].

Large-scale applications beyond proof-of-principle demonstrations require developing integrated techniques to enable the generation [22–26], transmission [27], and even processing of vector vortex beams on a photonic chip. Several pioneering works have demonstrated on-chip generation vector vortex beams with integrated microring resonators [22,23,25,26]. While the emission of vector vortex beams from the surface of integrated devices to free space has been widely investigated, generation and transmission inside a photonic chip remain to be realized and very challenging.

In this Letter, we demonstrate a direct generation of vector vortex beams inside a photonic chip based on mode

(71) Femtosecond Relativistic Electron Beam with Reduced Timing Jitter from THz Driven Beam Compression

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Femtosecond Relativistic Electron Beam with Reduced Timing Jitter from THz Driven Beam Compression

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We propose and demonstrate a method to reduce the pulse width and timing jitter of a relativistic electron beam through THz driven beam compression. In this method the longitudinal phase space of a relativistic electron beam is manipulated by a linearly polarized THz pulse propagating in a dielectric tube such that the bunch tail has a higher velocity than the bunch head, which allows simultaneous reduction of both pulse width and timing jitter after passing through a drift. In this experiment, the beam is compressed by more than a factor of 4 from 1.50 fs to 28 fs with the arrival time jitter also reduced from 97 fs to 4 fs, opening up new opportunities in using pulsed electron beams for studies of ultrafast dynamics. This technique provides an effective way to manipulate beam longitudinal phase space with a THz pulse and may have a strong impact on accelerator and medical science facilities that require femtosecond electron beams with tight synchronization to external lasers.

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Ultrafast electron beams with small timing jitter with respect to external lasers are of fundamental interest in accelerator and ultrafast science communities. For instance, such beams are essential for laser and THz driven accelerators with an externally injected electron beam [1–4] where the beam energy spread and beam energy stability largely depend on the electron bunch length and injection timing jitter, respectively. For MeV ultrafast electron diffraction (UED) [5–12] where electrons with a few MeV energy are used to probe the atomic structure changes following the excitation of a pump laser, the temporal resolution is primarily limited by the electron bunch length and timing jitter. Similar limitations exist for x-ray free-electron lasers (FELs) [13–15] too, since the properties of the x-ray pulses depend primarily on that of the electron beams. Therefore, one of the long-standing goals in accelerator and ultrafast science communities is to generate ultrafast electron beams with small timing jitter.

A photocathode rf gun is the leading option for producing a high brightness ultrafast electron beam for FELs and MeV UED (see, e.g., Refs. [16,17]). Because of the space charge effect, the beam pulse width is broadened and therefore bunch compression is needed to reduce the pulse width. Bunch compression requires first a mechanism to imprint an energy chirp (correlation between an electron's

energy and its longitudinal position) and then sending the beam through a dispersive element such that the longitudinal displacement of the electrons is changed in a controlled way for reducing the pulse width. For a MeV beam, this is typically achieved by first sending it through a rf buncher cavity in a zero-crossing phase where the bunch head ($t < 0$) is decelerated while the bunch tail ($t > 0$) is accelerated. This imprints a positive chirp $\delta = \delta E/c\beta > 0$ in the beam longitudinal phase space, where δ is the relative energy difference of an electron with respect to the reference electron and c is the speed of light. Then the electron beam is sent through a drift after which the electrons at the bunch tail catch up with those at the bunch head, leading to compression in the pulse width. Full compression is achieved when $\delta R_{50} = 1$, where $R_{50} = L/\beta^2$ is the momentum compaction of the drift with length L and β is the Lorentz factor of the electron.

Recently, sub-10 (nan) fs beams have been produced with this technique [18,19]. However, the rf phase jitter results in increased timing jitter after compression. While THz pulse-based time-stamping techniques have been developed to measure and correct the electron beam arrival time jitter [19,20] for UED, the detector response time limited this shot-to-shot correction technique to low repetition rate. It is highly desired (in particular for those

(72) Young's Double-Slit Interference in a Hydrogen Atom

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Young's Double-Slit Interference in a Hydrogen Atom

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We demonstrate the possibility of realizing Young's double-slit interference in a hydrogen atom via *in situ* ionizations. By exposing the hydrogen atom to a high-frequency intense laser pulse, the bound state divides into a dichotomic Kramers-Henneberger state whose photoelectron momentum distribution imprints a double-slit interference structure. The dichotomic hydrogen atom presents molecular peculiarities, such as charge-resonance enhanced ionization, electron spin flipping due to the non-Abelian Berry phase. In return, the photoelectron momentum distribution carrying the double-slit interference structure provides unambiguous evidence on the existence of Kramers-Henneberger states, and thus the *in situ* stabilization.

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Double-slit interference, with which Young explained light as waves in 1801, is one of the most vital experiments in history. Though more than two hundred years had passed, double-slit experiments continue to provide insights into both fundamental research and practical applications nowadays. For example, which-way information experiments provide testimony to the basic principle of quantum mechanics [1–3], a matterless double-slit experiment explored the fluctuation of vacuum [4], and the interference in the electron pair demonstrates the universal nonlocality of quantum entanglement [5]. As practical applications, double-slit interference can measure light coherence [6], detect molecular structures [7], and retrieve the motion of nuclei [8].

The double-slit experiment can be performed in diatomic molecules, in which the two Coulomb centers work as the double slits, and the valence electron plays the role of light in the original Young's experiment. The theoretical prediction made by Cohen and Fano [9] had been observed in molecules such as H_2 [10,11], N_2 [12], neon dimer [13], etc. It sounds counterintuitive to have such kind of double-slit interference in an atom since one atom is impossible to offer two Coulomb centers. However, there existed non-trivial analogies in the time domain. Lindner *et al.* [14] designed such kind of double slits, in which the ionization bursts produced by two successive wave crests of a few-cycle pump-stabilized pulse interfere with each other. These time-domain double slits were further studied by adding another orthogonally polarized streaking pulse [15,16]. Recently, Puschmann *et al.* [17] considered a two-path ionization of a rubidium atom via intermediate $5P$ and $6P$ states, in which the excitation path is undetermined and the interference in wave-vector space is analogous to double-slit interference.

The existence of metastable Kramers-Henneberger states (KHSs) [18–20], which are the eigenstates of the time-averaged Hamiltonian in a high-frequency laser field [21], provides a possibility of realizing the double-slit interference in an atom. The KHS is known to be relevant for ionization stabilization [22], low-energy photoelectron generated by the nonadiabatic coupling [23,24] and dynamic interference [25–27]. The KHS plays an important role as long as the systems have exact or approximate Floquet symmetry [28]. However, while the acceleration of neutral atoms [29,30] and the amplifying of air being [31,32] provide tantalizing hints [26], there is still no “smoking gun” evidence for a direct experimental confirmation of the KHS up to now. The ionization stabilization of Rydberg atoms [33] does not provide convincing evidence since the nonadiabatic coupling [23,24] in intensive fields always populates a superposition of Rydberg states, as their energy gaps are small, interfering stabilization [34] might be responsible for the observed phenomena. The unavoidable focal intensity volume average effect in practice implies that the ionization signals are always from a mixture of KHSs and laser-free bound states. Thus, shifts of the photoelectron kinetic energy [35] and fine structures of the photoelectron momentum distribution (PMD) from the KH atom [36] are blurred after the focal intensity volume average.

In this Letter, we demonstrate the possibility of realizing molecular Young's double-slit interference using a single hydrogen atom. By exposing the atom to a high-frequency-intense laser field, the distorted bound state forms a dichotomic KHS [17,38]. When the atom gets ionized by a probe pulse, the PMD carries its structural information. The distinct angularly distributed PMD is robust against the focal intensity volume average, which provides

(73) Engineering of Magnetic Coupling in Nanographene

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Engineering of Magnetic Coupling in Nanographene

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Nanographenes with arbitrary imbalance have a net spin according to their structure or topological charges. Here, we report deterministic synthesis of atomically precise nanographenes and their atomic-scale characterization on a gold substrate by using low-temperature scanning tunneling microscopy and scanning tunneling spectroscopy. Our results clearly confirm individual nanographenes have a single spin of $S = 1/2$ via the Kondo effect. In covalently linked nanographene dimers, two spins are antiferromagnetically coupled with each other as revealed by inelastic spin-flip excitation spectroscopy. The magnetic excitation fabrication in dimer can be well engineered by tuning the local spin density distribution near the connection region, consistent with mean-field Hubbard model calculations. Our work clearly reveals the engineering of magnetism in nanographenes and provides an efficient way to further explore the carbon-based magnetism.

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Magnetism is mostly associated with the $3d$ or the $4f$ elements, and very rare materials (e.g., Fe, Cr, and Ni) are ferromagnets at room temperature. In the past decade, many works explored the emergence of magnetism in carbon-based materials [1–27], and some systems even show Curie temperatures above the room temperature [17,20,24]. Magnetic carbon-based nanostructures are particularly promising for applications in spintronics [1,28], due to the versatility of organic chemical synthesis offering flexible ways to control the structure and tune the magnetic properties, which is hard to achieve from inorganic materials. In addition, carbon-based nanomaterials are expected to have high magnitudes of spin-wave stiffness [29], weak spin-orbit coupling [30], and hyperfine couplings [31] which promise for high temperature magnets, transport of spin-polarized currents [32], and spin-based quantum information processing [1,17,28].

Intuitively, carbon atoms are unlikely to form magnetic order since sp electrons prefer to pair up to form covalent bonds. Theoretical calculations addressing carbon-based magnetism have mainly focused on impurities, point defects, or reduced dimensionality, such as the ferromagnetic order of graphene zigzag edges [32–34], the large net spin in graphene nanostructures with arbitrary imbalance [35,34–38], and magnetic order in topological frustrated graphene nanostructures [39–40]. Magnetism of

graphene-based materials depends crucially on their atomic structure and surrounding environments. Because of the difficulties encountered in fabricating and characterizing atomically precise graphene nanostructures, the direct comparison between experimental results and theory calculations is impeded, leaving many questions unclear. Thus, it is of great demanding to study the magnetic properties of graphene nanostructures at the atomic level.

In this Letter, we report the engineering of magnetism in atomically precise nanographene monomers and covalently linked dimers. Scanning tunneling microscopy or spectroscopy (STM/STS), non-contact atomic force microscopy (nc-AFM), together with single particle mean-field Hubbard model and spin-polarized density functional theory (DFT) calculations are combined to determine the magnetic properties of the achieved nanographenes with atomic precision. Our studied nanographenes are fabricated by using the on-surface synthesis approach, which has been widely adopted to prepare atomically precise graphene nanostructures [5,6,11,25,41–44] [see detailed synthesis steps in Supplemental Material (SM), Figs. (S1) and (S2)]. Molecular precursors of 9,10-bis(4-bromo-2,6-dimethylphenyl)anthracene were thermally deposited on clean Au(111), and the sample was rapidly annealed to ~ 1000 °C to induce C-C addition reactions and cyclodehydrogenation reactions, forming atomically precise magnetic

(74) Multiple In-Gap States Induced by Topological Surface States in the Superconducting Topological Crystalline Insulator Heterostructure $\text{Sn}_{1-x}\text{Pb}_x\text{Te-Pb}$

PHYSICAL REVIEW LETTERS 125, 136802 (2020)

Multiple In-Gap States Induced by Topological Surface States in the Superconducting Topological Crystalline Insulator Heterostructure $\text{Sn}_{1-x}\text{Pb}_x\text{Te-Pb}$

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Superconducting topological crystalline insulators (TCIs) have been proposed to be a new type of topological superconductor where multiple Majorana zero modes may coexist inside the insulation of insulator systems. The bulk superconductivity of TCI has been realized, but it is quite challenging to detect the superconductivity of topological surface states inside their bulk superconducting gaps. Here, we report high-resolution scanning tunneling spectroscopy measurements on lateral $\text{Sn}_{1-x}\text{Pb}_x\text{Te-Pb}$ heterostructures using superconducting tips. Both the bulk superconducting gaps and the multiple in-gap states with energy differences of ~ 0.3 meV can be clearly resolved on TCI $\text{Sn}_{1-x}\text{Pb}_x\text{Te}$ at 0.38 K. Quasiparticle interference measurements further confirm the in-gap states are gapless. Our work demonstrates that the unique topological superconductivity of a TCI can be directly distinguished in the density of states, which helps to further investigate the Dirac and Majorana fermions inside the superconducting gap.

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Research on topological superconductors (TSCs) is rapidly developing due to their significance in the quantum science and technology involving Majorana zero modes (MZMs) and topological qubits [1–5]. Evidence for topological superconductivity including MZMs has been found in superconducting topological insulators (TIs) such as Bi_2Te_3 -type $\text{Tl}_2/\text{NbSe}_2$ heterostructures [6–9], corner-based superconductors (SCs) [10–13], transition metal dichalcogenide 2H-WS_2 [14,15]. In contrast to the TI protected by time-reversal symmetry, a topological crystalline insulator (TCI) possesses the topological surface states protected by crystalline symmetry [16–17]. The multiple surface Dirac cones are observed in the SnTe -type IV–VI semiconductors with a rock-salt crystal structure [18–20]. Recently, superconducting topological crystalline insulators (TCIs) have been predicted to be a new type of TSC containing MZMs [21,22]. The hybridization between MZMs is forbidden by lattice symmetry, so multiple MZMs are expected to coexist in a single vortex. The bulk superconductivity of TCIs can be induced through chemical doping [23–24] and superconducting proximity effect [25–27]. Tight-binding models predict that a superconducting TCI with topological superconductivity has multiple surface Andreev bound states inside its superconducting gap [28], which is topologically distinct from a conventional SC and a superconducting TI [29]. However, these unique properties have not been observed due to the challenges in the preparation of good samples with proper

doping levels and the characterization of topological surface states inside such a small superconducting gap of only ~ 1 meV.

In this Letter, we report high-resolution scanning tunneling microscopy (STM) and scanning tunneling spectroscopy (STS) measurements on $\text{Sn}_{1-x}\text{Pb}_x\text{Te-Pb}$ heterostructures. Atomically flat lateral $\text{Sn}_{1-x}\text{Pb}_x\text{Te-Pb}$ heterostructures are prepared by molecular beam epitaxy. The intrinsic superconductivity of both TCI $\text{Sn}_{1-x}\text{Pb}_x\text{Te}$ and conventional SC Pb can be directly detected. In order to increase the resolution, superconducting tips and low-temperature down to 0.38 K are used. The STS spectra are simulated with the standard superconductor-insulator-superconductor (S-I-S) tunneling model. The in-gap resonances have energy differences of ~ 0.3 meV. Quasiparticle interference measurements taken inside the superconducting gap of $\text{Sn}_{1-x}\text{Pb}_x\text{Te}$ further confirm the in-gap states are gapless. The *in situ* low-temperature STS study provides the first direct experimental evidence of multiple in-gap states induced by topological surface states in a superconducting TCI.

Our experiments were performed in a commercial STM system (UltraScan) with a ^3He cryostat. Details about the growth of $\text{Sn}_{1-x}\text{Pb}_x\text{Te-Pb}$ heterostructures ($x = 30\%$) are described elsewhere [30]. Figures 1(a) and 1(b) are STM topography and schematic diagram of a lateral $\text{Sn}_{1-x}\text{Pb}_x\text{Te-Pb}$ heterostructure respectively. The temperature dependence of the superconducting gaps taken on Pb is

(75) Strain Tunable Semimetal–Topological-Insulator Transition in Monolayer 1T0 -WTe2

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Strain Tunable Semimetal–Topological-Insulator Transition in Monolayer 1T-WTe₂

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A quantum spin hall insulator is manifested by its conducting edge channels that originate from the nontrivial topology of the insulating bulk states. Monolayer 1T-WTe₂ exhibits the quantized edge conductance in Dirac semiconductors, but because of its semimetallic nature, the coherence length is restricted to around 100 nm. To overcome this restriction, we propose a strain engineering technique to tune the electronic structure, where either a compressive strain along the *a* axis or a tensile strain along the *b* axis can drive 1T-WTe₂ into an full gap insulating phase. A combined study of molecular beam epitaxy and *in situ* scanning tunneling microscopy or spectroscopy then confirmed such a phase transition. Moreover, the topological edge states were found to be very robust in the presence of strain.

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Monolayer transition metal dichalcogenides (TMDs) in 1T' structure are predicted to be quantum spin hall insulators (QSIs) [1]. Within this family, WTe₂ and WSe₂ have been verified experimentally using angle-resolved photoemission spectroscopy (ARPES) [2–4] and scanning tunneling microscopy or spectroscopy (STM/STS) [2–5]. Moreover, monolayer 1T'-WTe₂ has already been confirmed to be topological nontrivial by various nonlocal transport measurements [6] and exhibits a fully quantized conductance up to 400 K [7]. However, the coherent edge channel length is found to be around 100 nm because of the semimetallic nature of the bulk states [7]. Pursuing both fundamental research and potential applications, to overcome this hindrance, it is necessary to manipulate the electronic structure of monolayer WTe₂ to realize a full gap insulator. In principle, the electronic structure of a material can be engineered using various external fields such as electric or mechanical/strain fields. Recently, electrical tuning has been applied extensively to TMD materials for many structures [8], superconductivity [9,10], and Berry curvature dipole [11]. For low-dimensional systems, strain engineering is a very general and effective approach and is technologically feasible. Examples include strain-induced topological phase transition [12], strain engineering Dirac states [13,14], strain-induced flat band [15,16], strain tunable magnetism in domain walls [17], strain-induced surface-dominated conduction [18], and strain-enhanced superconductivity [19]. However, related studies in the TMD family are relatively limited and unproved [1,20,21]. So far, no method has been found to tune the band structures of monolayer 1T'-WTe₂ effectively.

In this work, based on orbital analysis and the first-principles calculations, we proposed by *ab initio* to optimize the band structure of monolayer 1T'-WTe₂, and we found that a compressive strain along the *a* axis or a tensile strain along the *b* axis can convert monolayer WTe₂ from a semimetal to an insulator. This was further confirmed by a systematic study using molecular beam epitaxy (MBE) and STM, where strains were introduced by the distorted substrate. The strain status of WTe₂ films was measured by STM, and the corresponding band structures were characterized by differential conductance dI/dV spectra. Based on experimental data, we find that a moderate in-plane strain can open a full insulating gap in WTe₂, and the gap size is positively related with lattice constant *b* and negatively related with *a* (*a* and *b* are marked in Fig. 1(a)), consistently with the theoretical results. Moreover, the topological edge states are obtained on both semimetallic and insulating WTe₂ islands, indicating that the nontrivial topology is robust in the presence of external strains. Our results unambiguously demonstrate the feasibility of strain engineering to realize the fully gapped QSI in monolayer 1T'-WTe₂, making 1T'-WTe₂ a suitable material platform for developing tunable topological electronic devices with a long coherence length.

Monolayer 1T'-WTe₂ has a rectangular unit cell as shown in Fig. 1(a), where one W atom layer is sandwiched between two Te atom layers. Moreover, there is a Γ -point band inversion along the *a* axis, inducing a large band inversion at the Γ point and driving monolayer 1T'-WTe₂ into a Z₂ topological phase [1]. However, a finite overlap between the valence

OPTICS

A hybrid quantum memory-enabled network at room temperature

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Quantum memory capable of storage and retrieval of flying photons on demand is crucial for developing quantum information technologies. However, the devices needed for long-distance links are different from those envisioned for local processing. We present the first hybrid quantum memory-enabled network by demonstrating the interconnection and simultaneous operation of two types of quantum memory, an atomic ensemble-based memory and an all-optical loop memory. Interfacing the quantum memories at room temperature, we observe a well-preserved quantum correlation and a violation of Cauchy-Schwartz inequality. Furthermore, we demonstrate the creation and storage of a fully operable heralded photon chain state that can achieve memory built-in combining, swapping, splitting, tuning, and chopping single photons in a chain temporally. Such a quantum network allows atomic excitations to be generated, stored, and converted to broadband photons, which are then transferred to the next node, stored, and faithfully retrieved, all at high speed and in a programmable fashion.

INTRODUCTION

A quantum network consisting of a large number of nodes of quantum nodes and interconnecting channels systems an ever-growing need for quantum information science. Such a network could be used for quantum-enhanced distributed systems that promise to supersede classical systems in the fields of communication, computing, and metrology (1). Unfortunately, quantum channels suffer from experimental loss and high latency (2, 3). In addition, the prohibitive generation of quantum states limits the scale of quantum systems (4, 5). Quantum memories capable of storing quantum states for periods quantum communication need long distance as well quantum repeaters (6, 7) and avoid the considerable resource of quantum state generation by synchronizing stochastic events (8, 9).

Various works have been proposed on quantum memory design and relevant physical implementations, such as optical delay lines and cavities (10–12), electromagnetically induced transparency (EIT) (13–16), the Rydberg atom EIT (17) (DLCZ) protocol (18, 19), the photon-echo quantum memory (20, 21), the resonant Far-Off-Resonance (FOT) atomic memory (22–27) and lattice-Twisted spin (LTS) memory (28). To make a quantum memory device practical for scalable and high-speed quantum information processing, the key requirements that have to be satisfied include broad acceptance bandwidths, high efficiencies, long lifetimes, low noise level, and the capability of operating at room temperature.

However, for the past 20 years, it has been proven very challenging to meet all these requirements simultaneously. Considerable efforts have been made in pushing the bandwidth from the kilohertz and

megahertz regime to the gigahertz regime and compressing from millikelvin to room temperature (29–32, 33–37). As some major progress, EIT and most of resonant Far-Off-Resonance (FOT) (38) have been made that is impossible to be obtained independently (39). As for the off-resonant thermal memory, while it maintains high efficiency, a two-photon noise related to quantum-to-classical conversion strategy (40).

Recently, we have realized a broadband and room temperature far-off-resonant DLCZ (41) with quantum memory, capable of operating with a high fidelity in the quantum regime (42). Here, we are presenting a quantum memory that is broadband, room temperature, and more importantly, compatible with the FOT quantum memory for further storing heralded entangled photons (i.e., swapping in) and retrieving them (i.e., swapping out) for controlled operations without any additional noise. Progress toward such a memory has been made by using several kinds of atoms (43, 44) labeled with a limited lifetime.

Atomically loop an ensemble has shown the capability of capturing single photons by using addressable optical detuning, also it can avoid introducing any additional noise (38, 45, 46, 47). The delay of single-photon can be traced back to 1993 (48), and although being extremely straightforward, it is actually very challenging to realize in practice. Continuous efforts have been made in various implementations and feasibility (10, 12, 23, 24–27, 49, 50) (mainly), recently low loss has been achieved, enabling high efficiency single-photon (46, 45) and multiphoton generation (27). A loop quantum memory, consisting of an all-optical storage loop and controllable polarization switches, can capture broadband flying single photons by and/or at room temperature, which thus complements the FOT quantum memory.

Here, we propose and experimentally demonstrate a broadband room temperature and hybrid quantum memory-enabled network. The hybrid network combines a broadband FOT quantum memory with a compatible all-optical loop memory. The network generates a wideband of optical nodes in a built-in and time-tunable fashion in a warm atomic ensemble and traps these nodes in an all-optical loop (see Fig. 1A). The hybrid quantum memory system is capable of combining, swapping, splitting, and chopping single-

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(77) A scalable photonic computer solving the subset sum problem

SCIENCE ADVANCES | RESEARCH ARTICLE

OPTICS

A scalable photonic computer solving the subset sum problem

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The subset sum problem (SSP) is a typical nondeterministic polynomial-time (NP)-complete problem that is hard to solve efficiently in time with conventional computers. Photons have the unique features of high propagation speed, strong robustness, and low detectable energy level and therefore can be promising candidates to meet the challenges. Here, we present a scalable chip built on photonic computer to efficiently solve the SSP. We map the problem into a three-dimensional waveguide network through a femtosecond laser direct writing technique. We show that the photons sufficiently disperse into the networks and search for solutions in parallel. In the case of six resolve photons, our approach exhibits a dominant superiority in time consumption even compared with supercomputers. Our results confirm the ability of light to realize computations intractable for conventional computers and suggest the SSP as a good benchmarking platform for the race between photonic and conventional computers on the way toward “photonic supremacy.”

INTRODUCTION

NP-complete problems (1) are typically defined as the problem solvable in polynomial time on a Turing machine (Turing machine (TM), which defines the class problems) or computationally hard on conventional electronic computers. A general type of nondeterministic Turing machines. The subset-sum problem (SSP) with practical application in resource allocation (2) is a benchmark NP-complete problem (3), and its intractability has been demonstrated to be resistant to quantum attack (4, 5). Given a list of N integers, the SSP asks whether there is a subset of k whose sum is equal to the target T . Apparently, the number of subset grows exponentially with the problem size N , which leads to an exponential time taking and thus severely limiting the size of the problem that can be tackled in reality.

To cope the immense difficulty, some researchers attempt to solve NP-complete problems in polynomial time with polynomial resources. A noncomputing machine (6, 7) is presented to an NTM has been demonstrated, while the ambitious claim is not valid in a realistic environment with the variable time (8). Things or an NTM, where the magical words *P*, *N*, *NP* are realized by simultaneously exploring all computation paths, are proposed (11, 12). In the case of space or material parallel exploration provides an alternative to the great time consumption. As time is irreversible, non-reversible and completely out of use charge, it is reasonable to take physical resources for it. Besides the above NTM programs, similar mechanisms have been taken. For instance, the dramatically powerful electronic supercomputers with an integration of an increasing number

of processors (13) and molecule-based computation using large quantum or DNAs or amino molecules (14–18). Furthermore, optimized algorithms are applied to specific instances (19–21).

Although improvements have been made, conventional electronic computers are ultimately limited by their dissipation problem (22), which is also a possible candidate for non-computing machines consisting of continuous structures (23). The major obstacle in computation is limited by the slow movement (24–26) of the long-distance process (14, 18). Quantum computation is still hindered by its dimension and scalability (22, 23). Other proposals are still in the stage of theory (11, 12, 24–28). However, we notice that photons have been extensively applied to prove the principle demonstrations of supercomputing (27) even without quantum speed-up, including an NP problem such as prime factorization (29) and NP-complete problems such as traveling salesman problem (30), Hamiltonian path problem (31–33) and Hamiltonian cycle problem (34), the P-complete problem, binary counting (34–39), other computational functions (40), and algorithms (41, 42) are also demonstrated in a photonic regime. The successful applications imply that photons are potential excellent candidates to solve the NP.

Here, we present a photonic computer constructed with chip-coupling as processing units to solve the SSP of a physically scalable system. Like the conventional electronic computers or the molecule in molecular computers, photons (coupled in the optical source) are treated as individual computation carriers. They travel in chips along layered waveguide networks to perform parallel computations. The specific instances of the problem are successfully encoded into the networks according to particular rules. The existence of large sums is judged by the arrival of photons to the corresponding output ports of the network. We further investigate its scalability and performance in time consumption, showing the photon embedded advantage.

RESULTS

Configuration of the photonic computer for the SSP

The proposed photonic computer solving the SSP can be divided as a non- γ -on-Neuron Architecture (see the Supplementary Material

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(78) Extremely brilliant GeV γ -rays from a two-stage laser-plasma accelerator

SCIENCE ADVANCES | RESEARCH ARTICLE

PHYSICS

Extremely brilliant GeV γ -rays from a two-stage laser-plasma accelerator

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Recent developments in laser-wakefield accelerators have led to compact ultrashort X γ -ray sources that can deliver peak brilliance comparable with conventional synchrotron sources. Such sources normally have low efficiencies and are limited to 10^{22} photons/sr at the keV to MeV range. We present a novel scheme to efficiently produce collimated ultrashort γ -ray beams with photon energies tunable up to GeV by focusing a multi-petawatt laser pulse into a two-stage wakefield accelerator. The high-intensity laser enables efficient generation of a multi-GeV electron beam with a high density and tens-nC charge in the first stage. Subsequently, both the laser and electron beams enter into a higher-density plasma region in the second stage. Numerical simulations demonstrate that more than 10^{22} γ -ray photons/sr are produced with energy conversion efficiency above 10% for photons above 1 MeV, and the peak brilliance is above 10^{26} photons $s^{-1} mm^{-2} mrad^{-2}$ per 0.1% bandwidth at 1 MeV. This offers new opportunities for both fundamental and applied research.

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INTRODUCTION

High sources of high-energy γ -rays are essential tools (1–3) that are applied in basic science ranging from fundamental research (4–7) in astrophysics, particle and nuclear physics, to high-resolution imaging (8–9) in industry, biology, medicine, material science, and industry. These applications can benefit greatly from the availability of further compact γ -ray sources with low divergence, short pulse duration, high energy, and high peak brilliance. At present, widely used synchrotron (10) and X-ray free electron laser (XFEL) (11) can deliver X-ray pulses with peak brilliance in the range of 10^{21} – 10^{22} photons $s^{-1} mm^{-2} mrad^{-2}$ per 0.1% bandwidth (BW) (respectively). However, they are normally limited to photon energies ranging from a few keV to hundreds of keV, and pulse powers in the multi-fW level. In addition, the cost and size of these large-scale facilities are much more expensive than those of compact laser-wakefield accelerators.

On the other hand, compact laser-wakefield acceleration (LWFA) (12–16) has been developed rapidly over the past two decades (17) and offers a radically different approach—the acceleration length is plasma is often three orders of magnitude smaller as compared to conventional accelerators, providing the ability to drive the acceleration and radiation of high-energy particles on a much smaller scale. Against nV electron beams have been produced using LWFA (18), and femtosecond-scale X γ -ray pulses in the range of keV to MeV can be produced via LWFA-based betatron radiation (19–22) and Compton backscattering (23–25). The resulting radiation sources have typical peak brilliance of 10^{21} – 10^{22} photons $s^{-1} mm^{-2} mrad^{-2}$ per 0.1% BW, while the photon energies per shot is limited

to 10^3 photons with the laser to photon energy conversion efficiency at a very low level on the order of 10^{-4} . Although enhanced methods have been developed to enhance laser to radiation yield by using energetic particle beam-driven plasma-wakefields (26,27) and increasing transverse oscillating amplitudes (28), it remains a great challenge to significantly increase the energy conversion efficiency and to generate collimated γ -rays with high peak brilliance on the order of the MHz level and with high energies in the MeV to GeV regime. Currently, many cutting-edge applications and scientific research (29–31) require γ -rays with ultrahigh brilliance and photon energies far exceeding 1 MeV. These applications include exploring elementary particles (3), probing nuclear structures and photo-nuclear physics (32, 33), and studying quantum processes (34), which may finally give γ -ray sources in the MeV to GeV range.

Continuous development in ultrashort laser technology (35) provides possibilities for producing brilliant high-energy γ -ray sources. So far, considerable theory and simulation efforts have been made to develop such photon sources based on emission of energetic electrons accelerated in intense laser fields using laser-magnetically well, near-critical-density plasmas (36–38), laser-driven radiation pressure (39–40), laser-induced solid transients (39–41), laser scattering off electrons (42, 43), and the emission of micro-magnets cascades (44–45). However, there are unavoidable physical limitations on the γ -ray peak brilliance with these methods such as a very large divergence in direct laser excitation with electrons (46), emission of micro-magnets (47) with high laser intensity or 10^{20} – 10^{21} W/cm² (even at our center of magnitude higher than the highest quantum yieldable to date) is required to produce GeV photons. This requires tens-of-petawatt (PW) laser pulses with focused counter-diffraction-limited spots, which is very challenging. As soon as the laser intensity is reduced to the order of $\sim 10^{18}$ W/cm² (which is the highest intensity level of reliable operation of current high-power laser systems), the method mentioned above become increasingly inefficient for γ -ray emission. In addition, the ability to tune the photon energy, power, and brilliance is limited. It has been recently proposed that collimated γ -rays may be produced when ultradense relativistic electron bunches interact with emission via beam-plasma instabilities (48).

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(79) Multifunctional antiferromagnetic materials with giant piezomagnetism and noncollinear spin current



ARTICLE

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OPEN

Multifunctional antiferromagnetic materials with giant piezomagnetism and noncollinear spin current

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We propose a new type of spin-valley locking (SVL), named C-paired SVL, in antiferromagnetic systems, which directly connects the spin/valley space with the real space, and hence enables both static and dynamical controls of spin and valley to realize a multifunctional antiferromagnetic material. The new emergent quantum degree of freedom in the C-paired SVL is comprised of spin-polarized valleys related by a crystal symmetry instead of the time-reversal symmetry. Thus, both spin and valley can be accessed by simply breaking the corresponding crystal symmetry. Typically, one can use a strain field to induce a large net valley polarization/magnetization and use a charge current to generate a large noncollinear spin current. We predict the realization of the C-paired SVL in monolayer V_2Se_2O , which indeed exhibits giant piezomagnetism and can generate a large transverse spin current. Our findings provide unprecedented opportunities to integrate various controls of spin and valley with nonvolatile information storage in a single material, which is highly desirable for versatile fundamental research and device applications.

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(80) Sierpiński Structure and Electronic Topology in Bi Thin Films on InSb(111)B Surfaces

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Sierpiński Structure and Electronic Topology in Bi Thin Films on InSb(111)B Surfaces

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Deposition of Bi on InSb(111)B reveals a striking Sierpiński triangle (ST) like structure in Bi thin films. Such a fractal geometric topology is further shown to mirror the intricate electronic topology in a thin film. Relaxation of a large misfit strain of about 50% to 40% between Bi adlayer and substrate is revealed to drive the ST-like island formation. A Fermi-Korotky model is developed to illustrate the enhanced strain relief in the ST islands (offering the additional step energy cost). Besides a sufficiently large tensile strain-relieving ST-like structure, also require large adlayer-substrate and inter-adlayer elastic stiffnesses, and weaker inter-adlayer interatomic interactions.

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The Sierpiński triangle (ST) [1] is a self-similar fractal with a fractional Hausdorff dimension $d_H = \log 3 / \log 2 = 1.58496\dots$ and with fascinating mathematical properties [2]. Historically, such beautiful patterns were often utilized for aesthetic reasons, especially in architecture, e.g., on the medieval doors in Rome [3]. Many fundamental and intriguing phenomena in ancient and modern mathematics and physics, as well as computer science, are related to the STs, such as Pascal triangles [4], towers of Hanoi [5], chaos games [6], cellular automata [7], DNA computing [8,9], etc.

Theory and experiments have shown that fractal materials can exhibit peculiar and unusual thermal, mechanical, electric, electronic, electro-magnetic, and optical properties which are often desirable for prospective nanoscale applications [10,16]. Fractal quasicrystals have even been predicted to possess nontrivial electronic topology usually pertained to crystals [17,18]. Recently, there has been also considerable interest in fabricating ST-like fractal materials, although the synthesis of such systems is notoriously difficult [19]. In 2014, a terpyridine-based architecture mimicking ST was synthesized in solution [20]. Subsequently, several research groups have successfully fabricated the fractal supramolecular materials featuring the ST pattern on the various metal surfaces: Ag(111) [21], Ag(100) [22], Au(111) [23–28], Au(100) [29,30], and Cu(111) [30], as well as on a graphite surface [31]. These fabricated fractal supramolecules can be organic, metal organic, CO compounds,

etc. [11]. Theoretically, Monte Carlo simulations show that the formation of molecular STs in these systems is essentially related to directional intermolecular bonds [32,33].

As indicated above, previous fabrication of STs in experiments is by assembling small carbon-based molecules into supramolecules. There has not yet been any report of the synthesis for STs composed of single-element atoms. Obviously, experimental fabrication of such pre-shaped islands providing directional intermolecular interactions [32,33] in this work, by depositing Bi atoms on an InSb(111)B surface around 400 K, we have observed ST-like Bi films with a thickness of one to three monolayer (MLs) from our scanning-tunneling-microscope (STM) experiments. Our proposed theoretical model reveals that the formation of a Bi ST island requires an unusually large lattice mismatch between a monatomic Bi adlayer and substrate, relatively large adlayer-substrate and inter-adlayer elastic stiffnesses, and relatively weak inter-adlayer interatomic interactions. This formation mechanism essentially differs from that of the supramolecular STs synthesized previously.

In addition, it has been predicted theoretically that an ultrathin Bi film exhibits a subtle nontrivial topological property [34–36], which makes the investigations of heteroepitaxial growth of ultrathin Bi films fundamentally interesting. Experiments have demonstrated the existence

(81) Controlling Cell Motion and Microscale Flow with Polarized Light Fields

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Controlling Cell Motion and Microscale Flow with Polarized Light Fields

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We investigate how light polarization affects the motion of photosensitive algae, *Euglena gracilis*, in a dimensionally polarized field, cells swim approximately perpendicular to the polarization direction and from a normally isotropic velocity profile. When light polarization varies spatially, cell motion is modulated by local polarization. In such light fields, cells present complex spatial distribution and transport patterns which are controlled by topological properties of the underlying field. We further show that ordered cell swimming can generate directed microscale flow. Our experimental results are quantitatively reproduced by an active Brownian particle model in which preferential direction is continuously coupled to local light polarization.

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Natural microswimmers, such as bacteria and algae, can achieve autonomous motion by converting locally stored energy into mechanical work [1–15]. Such cellular motility is not only an essential aspect of life but also an important resource to develop artificial microswimmers, which propel themselves through self-generated fields of temperature, chemical concentration, or electric potential [1, 3–7, 12, 13]. Both natural and artificial microswimmers have been used in a wide variety of applications [16–19].

In properly limited in a fluctuating heterogeneous environment, microswimmers tend to adjust their motility in response to external stimuli [20–23]. For example, intensity and direction of ambient light can induce a variety of motility responses in photosynthetic microorganisms [24–48] and artificial microswimmers [49–55]; these responses have been frequently used to control microswimmer motion [27, 31, 54–56, 55–58]. Besides intensity and direction, light polarization can also affect microswimmer motility and lead to phototaxis. *Euglena gracilis* cells align their motion direction perpendicular to the light polarization, possibly to maximize the light absorption [54, 55]; artificial microswimmers consisting of two delicate transparent legs in the polarization direction [44]. These previous experiments have focused on uniform light fields [44, 54, 55]. The possibility to use complex polarization patterns to control phototactic microswimmers has not been explored.

In this Letter, we investigate *Euglena gracilis* cell motion in various polarized light fields in a quantitative and systematic fashion. Our experiments show that, while spatially uniform polarization aligns cells into a global nematic state with no net motion, spatially varying fields can induce both local nematic order and mean cell motion. Further, we show that ordered cell swimming motion generates fluid flow that can transport passive particles.

Using the experimental data of individual cells, we construct a model to describe the influence of local light polarization on cell orientation dynamics and quantitatively reproduce all experimental observations.

Experiments.—*Euglena gracilis* are unicellular flagellated microorganisms with a rod-shaped body of a length $\sim 50 \mu\text{m}$ and a width $\sim 5 \mu\text{m}$. As shown in Fig. 1(a) and Movie S1 in the Supplemental Material [59], cells swim at a mean speed $\sim 90 \mu\text{m/s}$ (with a standard deviation of $10 \mu\text{m/s}$) while rolling around their long axis at a frequency of 1–2 Hz [57]. A photoreceptor on *Euglena* cell surface, marked as a red dot in Fig. 1(b), senses surrounding light and generates signals to modulate flagellar beating pattern [57, 58].

In our experiments, *Euglena* culture is sealed in a disk-shaped chamber ($\sim 150 \mu\text{m}$ in thickness and 24 mm in diameter) which is placed in an illuminating light path, as shown in Fig. S1 [59]. A collimated blue light beam is used to excite cell photoreceptors; the default light intensity is $100 \mu\text{W}/\text{cm}^2$. Various polarized optical fields can be generated by using different birefringent liquid crystal plates and by changing relative angles between optical elements [59]. Cell motion is recorded by a camera mounted on a microscope. Default system cell density ($n_0 = 5 \text{ cells}/\text{mm}^2$) is sufficiently low that we can use a standard particle tracking algorithm [60] to measure position, orientation, and velocity of cells. The current work mainly focuses on steady-state dynamics that is invariant over time.

Uniformly polarized light field.—*Euglena* photoreceptor contains dichroically oriented chromoproteins which lead to polarization-dependent photo responses [57, 58, 59, 61]. As shown in Fig. 1(c), cells in a horizontally polarized field tend to orient and swim perpendicular to the polarization [54]; we denote such a targeted direction for cells as D_{\perp} .

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(82) X-Ray Tomography Investigation of Cyclically Sheared Granular Materials

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X-Ray Tomography Investigation of Cyclically Sheared Granular Materials

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We perform combined x-ray tomography and shear force measurements on a cyclically sheared granular system with highly transient behavior, and obtain the co-evolution of microscopic structures and macroscopic shear force during the shear cycle. We explain the macroscopic behavior of the system based on microscopic processes, including particle-level structural rearrangement and frictional contact evolution. Specifically, we show how contact friction can induce large structural fluctuations and cause significant shear dilation effect for granular materials, and we also construct an empirical constitutive relationship for the macroscopic shear force.

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Dense granular materials exhibit complex behaviors upon shear. Notable examples include significant dilation [1], formation of shear bands [2], emergence of anisotropic force networks [3], and critical state [4]. However, the theoretical understanding of sheared granular materials remains very challenging since conventional continuum mechanics are no longer applicable due to the disordered nature of granular materials, as well as the complex role friction plays [5,6]. In practice, empirical constitutive theories are normally employed, including the famous Mohr-Coulomb criterion [7] and more sophisticated viscoplastic models [8]. However, these models are mainly macroscopic ones with little microscopic basis despite the fact that it is well known experimentally that microscopic structure and dynamics can strongly influence the macroscopic response [9]. Granular materials belong in general to the family of disordered materials like metallic and colloidal glasses, foams, emulsions, etc., and physicists have tried to treat the flow of dense granular materials within the general framework of disordered materials [9]. For example, theories based on free volume [10], shear transformation zones [11], and soft glass rheology (SGR) model [12], which are built upon microscopic and mesoscopic information to predict the plastic behaviors of certain disordered materials, have been directly applied to the granular case [13,14]. However, granular materials possess some peculiar properties, e.g., force chains [15,16], significant volume and stress fluctuations [6,17,18], and complex microscopic dynamics [19], which are not shared among all disordered materials. These properties are mainly induced by the frictional interparticle contacts. Therefore, to construct the correct constitutive theory for the flow of dense granular materials, physics on particle and contact levels have to be considered [20].

In this Letter, we apply cyclic shear to a three-dimensional (3D) disordered granular system. The system displays

highly transient dilatational behavior and mechanical response since the system always has to reorganize upon sudden shear reversal [14,21]. By combining x-ray tomography and shear force measurements, we obtain both microscopic structure evolution and macroscopic mechanical response of the system simultaneously. Based on these results, we can explain the macroscopic behaviors of sheared granular matter through the microscopic structural information, including both particle-level structural rearrangement and frictional contact evolution. We can then understand why granular materials demonstrate significant shear dilation as compared to other disordered materials, and also construct an empirical constitutive relationship to account for the macroscopic shear force behavior.

Our granular system consists of bidisperse smooth plastic beads ($\mu = 0.8$, diameters 5 and 6 μm , 7000 particles of each size). The bidisperse nature ensures no crystallization occurs in the system. The particles are placed in a shear cell mounted on a linear stage as shown in Fig. 1(a). The shear cell is in simple shear geometry and has an inner dimension of $24d = 24d_x = 24d_y$, where



FIG. 1. (a) Schematic of the experimental setup. granular packing is sheared by the stage on the linear stage, and the macroscopic shear force is measured by the force sensor. (b) Evolutions of the global volume fraction ϕ and shear force F versus one cycle (ensemble averaged over 12 realizations).

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(83) Connecting Packing Efficiency of Binary Hard Sphere Systems to Their Intermediate Range Structure

PHYSICAL REVIEW LETTERS 127, 278001 (2021)

Connecting Packing Efficiency of Binary Hard Sphere Systems to Their Intermediate Range Structure

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Using computerized x-ray tomography we determine the three-dimensional (3D) structure of binary hard sphere mixture as a function of composition and size ratio of the particles q . Using a recently introduced k-point correlation function we reveal that this 3D structure has an intermediate and large length scales of surprisingly regular order, the symmetry of which depends on q . The initial structural correlation length has a maximum at the composition in which the packing fraction is highest. At this composition also the number of different local particle arrangements has a maximum, indicating that efficient packing of particles is associated with a structure that is locally maximally disordered.

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Hard sphere (HS) systems play a paramount role in statistical mechanics and material science since they are important models to study the behavior of many-particle systems like liquids, colloids, and metallic glasses [1–4]. Also in the domain of granular materials many experimental and theoretical studies have used HS-like systems since granular particles usually have a high modulus and the simplicity of the particle shape allows us to probe the influence of friction and roughness on the properties of granular assemblies [5,6]. Since for HS the interaction energy between the particles is trivial, the properties of such systems is encoded in their structure and hence many previous studies have aimed to characterize it on various length scales [1,4,6,7,8]. Many of these investigations focused on one-component systems since this choice facilitates the description of the packing structure [1,7]. However, in the context of the glass transition one often uses weakly polydisperse samples or a slightly asymmetric binary mixture since this suppresses the crystallization of the liquid while keeping its structure close to the well understood one-component case [9–10]. These studies have shown that, depending on the composition and the packing fraction, the local structure, i.e., the first nearest neighbor shell, shows a surprisingly rich variety [10–13]. In contrast to these almost one-component systems, the case of mixtures in which the particles have a size ratio that differs significantly from unity has been studied much less [16–22]. Understanding this structure is, however, important since such highly asymmetric systems are relevant for describing the behavior of real granular materials which are usually highly polydisperse as well for the glass-forming ability of multicomponent systems like metallic glasses in which the atoms can have very

different radii [23–29]. In addition it has recently been found that such strikingly asymmetric mixtures can have local structures with unexpected symmetries which can, e.g., be used to create novel materials via self-assembly [30]. The main reason for our lack of understanding of these systems is that their theoretical description is significantly more complex and experiments on colloidal systems are hampered by the precise control of size ratio, while granular systems are prone to the phenomenon of phase separation [31,32]. A further problem that one faces with these systems is the difficulty to characterize their structure on intermediate length scales since the presence of two particle sizes gives rise to a highly complex distance-dependency of the partial radial distribution functions, i.e., the standard quantities that are used to characterize the structure of many-body systems [19,33–35]. Because of these difficulties there is at present little insight on how the size ratio or the composition affects the packing density in the structure of asymmetric HS systems.

In this work we use the computerized tomography (CT) technique to probe how composition and size ratio affect the packing fraction of a binary granular system. Using a recently developed method to characterize the structure of disordered systems at intermediate and large length scales we are able to show that a high packing fraction is intimately related to the presence of a short structural correlation length.

Our system is a binary mixture of particles (arbitrary butadiene styrene plastic), denoted in the following as “big” (b) and “small” (s) particles, with size ratio $q = d_b/d_s$, where d_b and $d_s = 1$ (in) are the respective diameters. The number fraction of small particles will be denoted by ϕ . We consider two size ratios: $q = 1.33$, which corresponds

(84) Dynamical Transitions and Critical Behavior between Discrete Time Crystal Phases

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Dynamical Transitions and Critical Behavior between Discrete Time Crystal Phases

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In equilibrium physics, spontaneous symmetry breaking and elementary excitation are (or) intrinsically closely related with each other: the symmetry and its spontaneous breaking not only control the dynamics and spectrum of elementary excitation, but also determine their underlying structures. In this Letter, based on an exactly solvable model, we propose a phase-ramping protocol to study an excitation-like behavior of a non-equilibrium quantum matter, a discrete time crystal phase with spontaneous temporal translational symmetry breaking. It is shown that slow ramping could induce a dynamical transition between two Z_2 symmetry breaking time crystal phases in pure domain, which can be considered as a temporal analog of the soliton excitation spatially sandwiched by two degenerate charge density wave states in polyacetylene. By tuning the ramping rate, we observe a critical value at which point the transition duration diverges, resembling the critical slowing down phenomenon in nonequilibrium statistical physics. We also discuss the effect of stochastic sequences of such phase-ramping processes and its implication to the stability of the discrete time crystal phase against noisy perturbations.

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Usually, the ground state energy of an equilibrium system does not have much to do with its observable behavior. What is physically important are the properties of low-lying excited states, which are likely to be excited owing to weak external fields or relatively low temperatures. For example, the thermal and elastic properties of a solid are determined by a few number of lattice wave excitations known as phonons [1]. Nonequilibrium quantum matter fundamentally differs from its equilibrium counterparts, and has received considerable interest in various fields ranging from ultracold atoms [2] to solid state physics [3] over the past decade. However, compared to equilibrium cases, much less is known about the “excitation” of nonequilibrium quantum matter, even its definition may be questionable. In above its relationship with fundamental properties of nonequilibrium quantum matter, e.g. the symmetries and their spontaneous breaking.

As a prototypical example of nonequilibrium quantum matter, the time crystal (TC) phase has offered new possibilities for the spontaneous symmetry breaking (SSB) paradigm [4] and have attracted considerable interest in its different forms [5–27]. Such a intriguing state, despite being proven to be forbidden in equilibrium [23,24], has been experimentally realized in nonequilibrium settings with periodic driving [25,26]. Its physical observables develop persistent oscillations whose periods are an integer multiple of the Hamiltonian period, thus spontaneously break the discrete temporal translational symmetry (DTTS). In equilibrium systems, the SSB is

closely related to the elementary excitation: it does not only affect the dynamics and spectrum of an elementary excitation, but it also determines its structure. For example, for a one-dimensional 1D system, a spontaneous discrete spatial translational (e.g., Z_2) symmetry breaking allows certain solitonlike excitation: a topological defect sandwiched by two Z_2 symmetry breaking phases [27]. A profound question is how to generalize such an idea into the nonequilibrium quantum matters with intriguing SSB absent in equilibrium physics (e.g., DTTS breaking).

In this Letter, we address this issue by studying a driven 1D interacting bosonic model that can manifest a subharmonic response for physical quantities, a signature of a discrete time crystal (DTC) phase. We impose a time-dependent perturbation on top of the periodical driving, which transiently breaks the original time translational symmetry, and we then monitor the response of the physical observable. It is shown that a slow perturbation may induce a dynamical transition between two Z_2 symmetry breaking DTC phases, as shown in Fig. 1. By tuning the ramping velocity, one can observe a critical point, at which the transition duration diverges, a reminiscence of the critical slowing down phenomenon [28,29]. Finally, we discuss the effect of multiple random phase-ramping processes and its implication to the stability of the discrete time crystal phase against noisy perturbations.

Model and method.—We consider a 1D hardcore bosonic model with an infinite long-range interaction. The Hamiltonian reads as follows:

(85) Experimentally Detecting Quantized Zak Phases without Chiral Symmetry in Photonic Lattices

PHYSICAL REVIEW LETTERS 127, 147401 (2021)

Experimentally Detecting Quantized Zak Phases without Chiral Symmetry in Photonic Lattices

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Symmetry plays a major role in identifying topological phases of matter and in establishing a direct connection between protected edge states and topological bulk invariants via the bulk-boundary correspondence. One-dimensional lattices are deemed to be protected by chiral symmetry, exhibiting quantized Zak phases and protected edge states, but not for all cases. Here, we experimentally realize an extended Su-Schrieffer-Heeger model with broken chiral symmetry by engineering one-dimensional zigzag photonic lattices, where the long-range hopping breaks chiral symmetry but ensures the existence of inversion symmetry. By the averaged mean displacement method, we detect topological invariants directly in the bulk through the continuous-time quantum walk of photons. Our results demonstrate that inversion symmetry protects the quantized Zak phase (no odd states can disappear in the topological nontrivial phase, thus breaking the conventional bulk-boundary correspondence). Our photonic lattice provides a useful platform to study the interplay among topological phases, symmetries, and the bulk-boundary correspondence.

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Topological phases of matter are fascinating states that escape from the standard description of Ginzburg-Landau theory, exhibiting protected edge states and quantized topological properties in the bulk [1–5]. A paradigm of topological phases in condensed-matter physics is the integer quantum Hall state in a two-dimensional electron gas [1]. Topological order is of major relevance in a variety of nonelectronic systems as well, including mechanical platforms, cold atoms, acoustics, trapped ions, and photonics [6–15]. A central result in the theory of topological matter is that bulk topological invariants are connected to the number of protected edge states (bulk-boundary correspondence) [4,5]. Hence, topological phases can be probed either in the bulk or at the edges, as demonstrated in several recent experiments for Hermitian and even some non-Hermitian systems [15–34]. However, the common belief that a nontrivial topological phase implies the presence of protected edge states is not a general result, and even in Hermitian systems it is known that there could exist nontrivial topological phases that do not exhibit edge states at all [35–37]. In such systems, edge dynamics alone cannot thus provide a safe diagnostic of a topological phase. A notable example is provided by one-dimensional

(1D) and 2D lattices with broken chiral symmetry but with preserved inversion symmetry [3,38–41]. In the Altland-Zirnbauer (AZ) classification, nontrivial topological phases are identified by three main symmetries: time-reversal, chiral, and particle-hole symmetries [3,42]. In this scheme, a topologically nontrivial 1D system should possess chiral symmetry, and the bulk topological invariant is the winding number corresponding to a quantized Zak phase [43]. A paradigmatic example is provided by the famous Su-Schrieffer-Heeger (SSH) model of polyacetylene [7,4], where the bulk-boundary correspondence holds [3,44]. Experimental demonstrations of quantized Zak phases in the SSH model with chiral symmetry have been reported in recent works [29,45] using different platforms and methods to extract the Zak phase from bulk dynamics. Long-range hopping in the SSH model then breaks chiral symmetry makes the system topologically trivial. However, when inversion symmetry is preserved, the Zak phase remains quantized and one can introduce a \mathbb{Z}_2 topological number that identifies two distinct topological phases, which are not caught by the AZ classification [36–38]. Interestingly, unlike the chiral SSH model, in the extended SSH model with inversion symmetry only, the quantized Zak phase can

Symmetry-Induced Error Filtering in a Photonic Lieb Lattice

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Quantum computation promises unassistedly parallel information processing capacity by harnessing the superposition and entanglement of quantum states. However, it is still challenging to build universal quantum computation due that the reliability and scalability are limited by unavoidable noise on qubits. Meritiful topological properties like quantum Hall phases are found capable of offering protection, but require stringent conditions of topological band gaps and broken time-reversal symmetry. Here, we propose and experimentally demonstrate a symmetry-induced error filtering scheme, showing a more general role of geometry in protection mechanism and application. We encode qubits in a superposition of two spatial modes on a photonic Lieb lattice. The geometric symmetry induces the system with topological properties featuring a flat band touching, leading to distinctive transmission behaviors of π -valued qubits and 0-phase qubits. The geometry exhibits a significant effect on filtering phase errors, which also enables it to monitor phase deviations in real time. The symmetry-induced error filtering can be a key element for encoding and protecting quantum states, suggesting an emerging field of symmetry-protected universal quantum computation and noisy intermediate-scale quantum technologies.

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Quantum systems have been considered as powerful platforms for processing massive information and solving classically intractable problems by encoding qubits [1,2]. Analog quantum computation has demonstrated the supremacy over classical computation [3,4], whereas the step to universal quantum computation is still blocked by noises [5]. Fault tolerance is crucial for realizing reliable quantum computation, which has been explored on various systems, such as trapped ions [6], spins in semiconductors [7], superconducting circuits [8], solid state [9] and quantum optics [10]. However, implementation of error correction usually needs ancillary qubits with extra hardware overhead [11], which limits the scalability of quantum computation. For example, a qubit encoded in a superposition of two photonic modes may undergo phase errors in logical operations, which can be corrected by cluster states [12,13] but put daunting constraints on quantum systems including sources and measurements.

Geometries play a key role widely in natural and artificial materials, and induce extraordinary particle behaviors featured by localization and regular distribution [14–21]. Compared with dispersionless localization [22–24], the self-stabilization of distribution modes protected by geometries is more attractive for quantum computers, such as quantum Hall and anomalous quantum Hall phases with topological band gaps and broken

time-reversal symmetry [25,26]. Regardless of the band topology, many symmetric geometries can equally create self-stabilizing states. Two-dimensional materials with rotational and inversion symmetries can bring about Dirac points possessing extraordinary properties of topological charges [27–29]. A Lieb lattice builds a square-lattice geometry with fourfold rotational symmetry [30,32] and builds a topological system with a flat band touching at Dirac points, which enables stable states with special phase features and makes it possible to implement a filter capable of exciting and stabilizing specific quantum states [33–35]. Such a filter will be able to eliminate phase errors and improve the purity of quantum states.

Our work is based on the photonic Lieb lattice with three sites forming a right triangle in each unit cell. Waveguides consisting of periodic unit cells are fabricated by femto-second laser direct writing and are arranged in a ribbon, as shown in Fig. 1(a). The evolution of quantum states through the photonic lattice can be described by an equivalent Schrödinger equation $i\partial_t\Psi = H\Psi$, with evolution length z considered as time t . The Hamiltonian is expressed by the tight-binding limit as [21],

$$H = \sum_{\langle i,j \rangle} J_{ij} c_i^\dagger c_j + \text{H.c.} \quad (1)$$

(87) Microscopic structure and dynamics study of granular segregation mechanism by cyclic shear

SCIENCE ADVANCES | RESEARCH ARTICLE

CONDENSED MATTER PHYSICS

Microscopic structure and dynamics study of granular segregation mechanism by cyclic shear

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Granular mixtures with size difference can segregate upon shaking or stress. However, the quantitative study of this process remains difficult because it can be influenced by many mechanisms. Conflicting results in similar experimental systems are frequently obtained when the experimental conditions are not well controlled, which is mainly due to the fact that many mechanisms can be at work simultaneously. Moreover, it is often that microscopic or empirical measures, which lack microscopic physical bases, are used to explain the experimental findings and therefore cannot provide an accurate and complete depiction of the overall process. Here, we carry out a detailed and systematic microscopic structure and dynamics study of a cyclically sheared granular system with rigorously controlled experimental conditions. We find that both convection and arching effect play important roles in the segregation process in our system, and we can quantitatively identify their respective contributions.

Supporting data for this article are available at <https://doi.org/10.1126/sciadv.abe0000>.
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INTRODUCTION

Granular mixtures consisting of different-size particles tend to segregate under shaking or stress, which has important implications for many industrial processes such as the mixing or segregation of a stream (1). It is also relevant to geophysical processes like landslides, avalanches, and block size distribution on asteroids (2,3). Granular segregation can be induced by size, density, friction, and inertia (4), while size remains the dominant factor (4–6). Owing to its ubiquitous existence and importance, it is crucial to gain a fundamental understanding of segregation mechanisms based on systematic experimental investigations. Continuous segregation models (7) subsequently developed based on them (7b–7d). However, the experimental investigations of granular segregation are highly nontrivial. This is due to the fact: (i) Complex dynamics associated with shaking or stress cannot be fully resolved in three dimensions (8D). Early investigations used particle tracking or the particle surface (18). These approaches were done at a stress without a deep understanding in the underlying mechanisms (18). X-ray tomography, magnetic resonance imaging, and interferometric scattering method have provided some gain insights in 3D segregation process (19–21). However, more work is needed to combine microscopic structure and dynamic information to comprehensively investigate the underlying mechanisms. Unfortunately, most experimental investigations are limited to 2D systems (16) using binary and monodisperse (or unimodal) mixtures only for the segregation process. The most well-known factor is particle size but density, the presence of material fluid or friction (22) and influence the segregation process (1, 23–25). Even when only size disparity (size mechanism including arching effect) (19–21), convection (21) and thermal gradient (26) have all been demonstrated to play important roles in the segregation process. By slowness gradient in the system can lead to segregation, and all these mechanisms can induce granules in the system to a certain

way (21). It is therefore vital to have controlled experiments that can separate their individual contributions.

In the present study, we use X-ray tomography to investigate granular segregation in a quasi-steadily sheared granular medium. When the system is under quasi-steady shear and therefore particle velocity or kinetic energy is constant, mechanisms like thermal diffusion (24, 25) or kinetic mechanisms (19, 21) can be excluded. This is in contrast to agitation by shaking in which granules in both steady-state and velocity axis, the heterogeneously varying the shear rate, we obtain particle-level structure and dynamics. We find that convection and granules arching mechanisms are the two main mechanisms that contribute to segregation in the dense packing and dilute shear limit.

RESULTS

In our system, small particles segregate from background particles upon cyclic shear and each cycle will not disturb them when seismic effect system is applied. This segregation process is a combination effect of several segregation mechanisms working in different time scales ranging from single to thousands of cycles. It is experimentally very costly to do tomography scans covering the full time scale. To improve experimental efficiency, we follow some specific experimental protocols to address specific time-scale properties of the system, and corresponding segregation mechanisms. Further details can be found in Materials and Methods, which can substantially reduce the scan number needed. In the following, we focus in detail how different mechanisms contribute together in the segregation process of our system.

Convection-induced global segregation

One of the main mechanisms proposed for granular segregation by convection (27), by tracking the motion of each of the background particles using processed 2D cut can obtain the global convection pattern as shown in Fig. 1C. This convection mainly consists of particles flowing in the center and moving down at the sides at two shear walls and to form a like dipole at low stresses. The most distinct of the convection convection starts continuously as depth increases. Figure 1B shows the typical length l of convection for $H = 12$ mm

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(88) Anomalous Contribution to the Nematic Electronic States from the Structural Transition in FeSe Revealed by Time- and Angle-Resolved Photoemission Spectroscopy

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Anomalous Contribution to the Nematic Electronic States from the Structural Transition in FeSe Revealed by Time- and Angle-Resolved Photoemission Spectroscopy

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High-resolution time- and angle-resolved photoemission spectroscopy were made on FeSe superconductors. With smallest photoemission, two critical transition features that correspond to two distinct electronic phase transitions were found only in the d_{xy} -orbital-derived band near the Brillouin zone center within ~ 40 meV and energy resolution. Upon compression to the double temperature dependent magnetism, we conclude that there are two equilibrium-like phase transitions (at approximately 90 and 120 K) above the superconducting transition temperature, and an anomalous contribution on the scale of 10 meV to the nematic states from the structural transition is experimentally determined. Our observations strongly suggest that the electronic phase transition at 120 K involves d_{xy} orbitals, and in the energy band development of FeSe, and furthermore, the contribution of the structural transition plays an important role in the nematic phase of iron-based high-temperature superconductors.

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Nematic order in iron-based superconductors refers to an electronic phase with broken rotational symmetry but preserved translational symmetry, and usually occurs at decreasing temperature combined with a lattice structure transition from tetragonal to orthorhombic phase. Such a nematic phase may play a pivotal role in high-temperature superconducting pairing, and its underlying physics is a predominant topic in the study of iron-based systems [1–6]. The electronic nematic phase can be a result of the tetragonal-to-orthorhombic structural phase transition, orbital correlation, or the set-in of magnetic order, and early theoretical and experimental work suggested that the transition to orthorhombic lattice distortion may be too small to drive the nematic electronic properties, such as in-plane resistivity [7–12]. However, there is still a lack of evidence of a purely electronic (nematic) phase transition, and how much the structural transition contributes to the band renormalization is still not clear experimentally.

FeSe with a superconducting transition at approximately 9 K, undergoes a tetragonal-to-orthorhombic lattice deformation and a breaking of rotational symmetry of electronic order (nematic phase) at the same temperature of approximately 90 K without the long-range magnetic order as discovered in many other iron-based materials, allowing us

to study the pure nematic phase over a wide temperature range [1,13–18]. However, the extent of the individual contributions of the structural transition and the purely electrical order to the nematic electronic states is still unknown. As most characterizations of nematic order are measured at thermal equilibrium conditions [19–24], it is difficult to differentiate the nematic electronic order from that of the structural transition and determine the contribution to the nematic states from the structural transition. Taking advantage of ultrafast experiments by using strong femtosecond photon pulses to excite only the electronic states, retaining the preserved orthorhombic lattice structure, it is possible to drive only the electronic phase transitions in an ultrafast manner without deforming the crystal lattice [25,26]. Thus, it is possible to isolate the electronic states from the lattice, and to clarify the origin of both the nematic and structural phase transitions and estimate the contributions of the structural transition to the nematicity.

In this Letter, we report two distinct purely electronic phase transitions and two equilibrium phase transitions above the superconducting transition temperature in FeSe revealed by time- and angle-resolved photoemission spectroscopy (TRAPES). In particular, with improved energy- and pump fluence resolution [27], we demonstrate that the

Nonlinear Thouless Pumping: Solitons and Transport Breakdown

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One-dimensional topological pumping of matter solitons in two coupled optical lattices moving with respect to each other is considered in the presence of intrinsic nonlinearity. It is shown that there exists a threshold nonlinearity level above which the matter transfer is completely arrested. Below this threshold, the transfer of both dispersive wave packets and solitons occurs in accordance with the predictions of the linear Chern numbers ($\pm\pi, 0$) quantized and distributed by the local dynamical Chern numbers of the lowest bands. The breakdown of the transport is also explained by nontrivial topology of the bands. In that case, the nonlinearity induces Rabi oscillations of atoms between two (or more) lowest bands. If the sum of the dynamical Chern numbers of the populated bands is zero, the oscillatory dynamics of a matter soliton in space occurs, which corresponds to the transport breakdown. Otherwise, the sign of the Chern numbers of the nonlinearly excited bands determines the direction and magnitude of the average velocity of matter solitons that remain quantized (and admit fractional values). Thus, even in the strongly nonlinear regime the topology of the linear bands is responsible for the evolution of solitons. The transition between different dynamical regimes is accurately described by the perturbation theory for solitons.

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Controlled unidirectional transfer of matter, heat, or physical quantities like charge or spin in a medium is a fundamental problem. It was discovered by Thouless [1] that one-dimensional (1D) pumping of electrons by a slowly moving periodic potential over one pumping cycle is quantized with quanta being determined by the Chern numbers considered in the extended coordinate-time space. To date, Thouless pumping was experimentally observed in systems of cold bosonic [2–4] and fermionic [5] atoms, spin systems [6], optics [7–9], acoustics [10], and plasmonics [11].

Although topological pumping was originally introduced for essentially linear systems [1, 12], physical realizations mentioned above provide opportunities for investigation of the phenomenon in nonlinear settings, for example, in the presence of optical Kerr nonlinearity or two-body interactions in Bose-Einstein condensates (BECs). Interatomic interactions were intrinsically present in experiments with BECs in bipartite magnetic lattices where (un)quantized transport was observed [13] and in the formation of the Mott insulator phase involving light binding of atoms upon pumping [3]. Nonlinear effects were explicitly included in the single-band light-binding approximation for a gas of several spinor states of fermions in optical lattices [14] and in model of a nonlinear interferometer [15] where the obtained pump was fractional. Topological transport of interacting photons was described in [16]. Qualitatively new effects appear in the

presence of sufficiently strong interactions. Using a Rice-Mele model, accounting for spin and the Hubbard interaction term in [17] the authors reported the breakdown of the Thouless pumping, explained by closure of the gap computed using amplitude-dependent many-body ground states under twisted boundary conditions. Apparently, the effect may be dependent on the specific model, since in another realization of a spinful Rice-Mele model subject to both open and twisted boundary conditions, breakdown of pumping was not found [18].

All above results dealt with repulsive interatomic interactions and thus with stable nonlinear states of constant amplitudes. In this Letter, we address nonlinear Floquet pumping of matter waves in a BEC with a negative scattering length, loaded in an optical superlattice created by two lattices, counterpropagating with respect to the other. In this regime a well-localized matter soliton is formed and excitations from the upper bands of the linear lattice spectrum acquire special importance [19,20]. Our two main findings are as follows: First, there exists a threshold amplitude below which the pumping closely follows the prediction of the linear theory [1], even when solitons are formed. Above the threshold amplitude, the transport subturns to a fractional, acquires direction opposite to the moving sublattice velocity, or is arrested, resulting in oscillatory motion of a soliton. Second, oscillations of a soliton, relatively well described by the perturbation theory [21], have topological nature. They occur due to the

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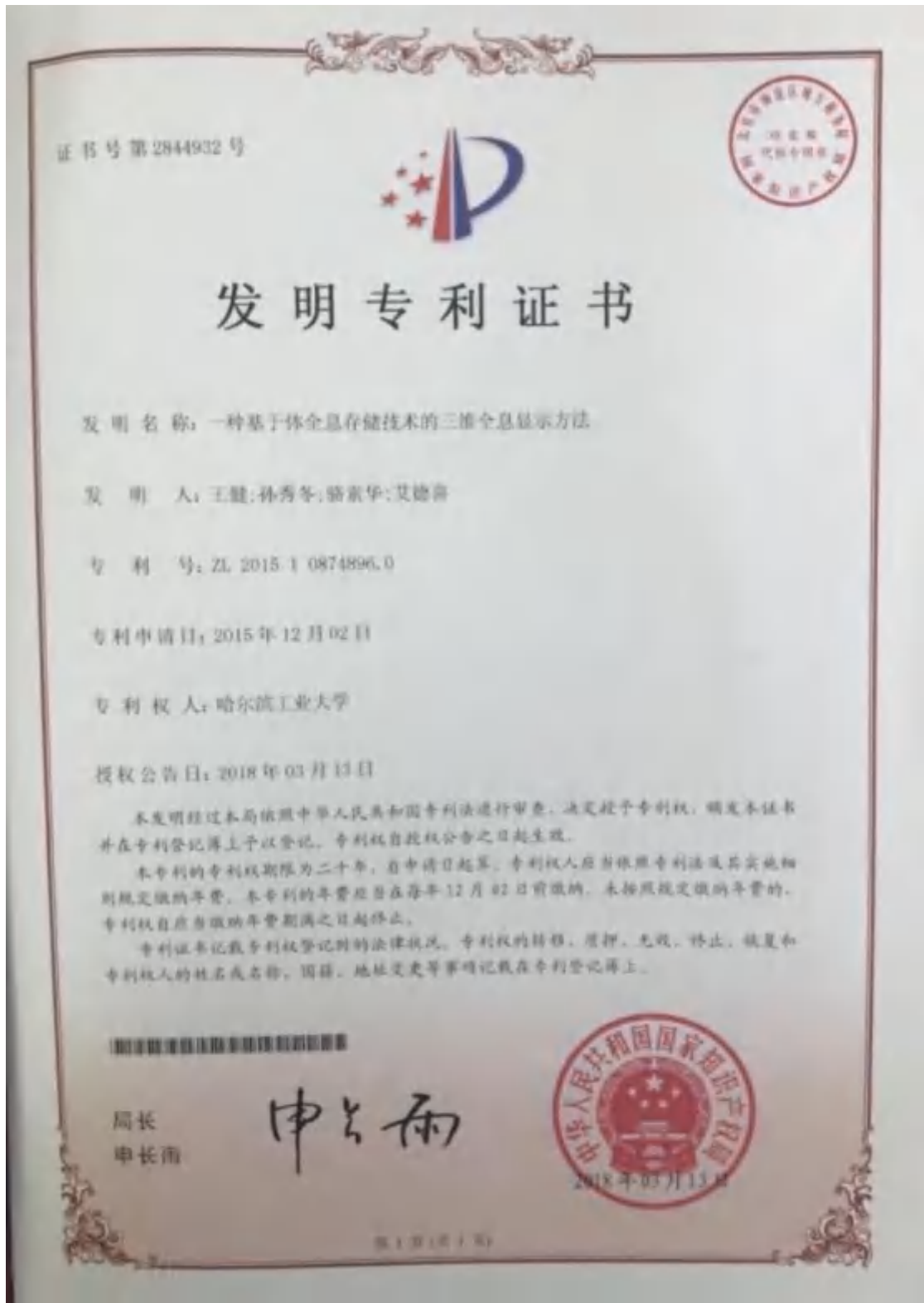
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| 1 | 一种用于二维材料薄膜厚度测量的高效光学测量方法 | 发明专利 | ZL201811548329.6 | 2020 |
| 2 | 一种基于体全息存储技术的三维全息显示方法 | 发明专利 | ZL201510874896.0 | 2018 |
| 3 | 电控连续渐变折射率电光晶体偏转器 | 发明专利 | ZL201410338950.5 | 2017 |
| 4 | 采用基于短腔腔增强关联光谱技术测量气体浓度的装置测量气体浓度的方法 | 发明专利 | ZL201310738166.9 | 2016 |
| 5 | 多通池内气体吸收光谱和吸收光程同步测量的装置和方法 | 发明专利 | ZL201810503921.8 | 2021 |
| 6 | 一种基于光学捷变频的快速气体吸收光谱测量装置及方法 | 发明专利 | ZL201910288076.1 | 2021 |
| 7 | 气体分子高分辨吸收光谱实验仪 | 实用新型 | ZL202022671068.6 | 2021 |
| 8 | 一种采用气体放电驱动天线收发信号的装置 | 发明专利 | ZL201811504774.2 | 2020 |
| 9 | 一种微波探测系统及方法 | 发明专利 | ZL201610716125.3 | 2019 |
| 10 | 一种基于霍尔效应精确定位和测量转速的实验教学装置 | 实用新型 | ZL201820959101.5 | 2019 |

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| 11 | 反射式光学偏转器 | 发明专利 | ZL202010929231.6 | 2022 |
| 12 | 基于聚合物透明电极的太赫兹波段电控液晶相移器的制备方法 | 发明专利 | ZL201610251983.5 | 2018 |
| 13 | 一种非侵入红外复合吸收精测血糖变化的方法 | 发明专利 | ZL201610614860.3 | 2019 |
| 14 | 基于波导与谐振腔耦合的微纳光学马达及其驱动方法 | 发明专利 | ZL201810132603.5 | 2019 |
| 15 | 一种高效样品微粒汇集装置及光学汇集方法 | 发明专利 | ZL202011155192.5 | 2021 |

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(3) 电控连续渐变折射率电光晶体偏转器



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(5) 多通池内气体吸收光谱和吸收光程同步测量的装置和方法



(6) 一种基于光学捷变频的快速气体吸收光谱测量装置及方法



(7) 气体分子高分辨吸收光谱实验仪

证书号第 13450203 号



实用新型专利证书

实用新型名称：气体分子高分辨吸收光谱实验仪

发 明 人：娄秀涛;赵海发

专 利 号：ZL 2020 2 2671068.6

专利申请日：2020 年 11 月 18 日

专 利 权 人：哈尔滨工业大学

地 址：150001 黑龙江省哈尔滨市南岗区西大直街 92 号

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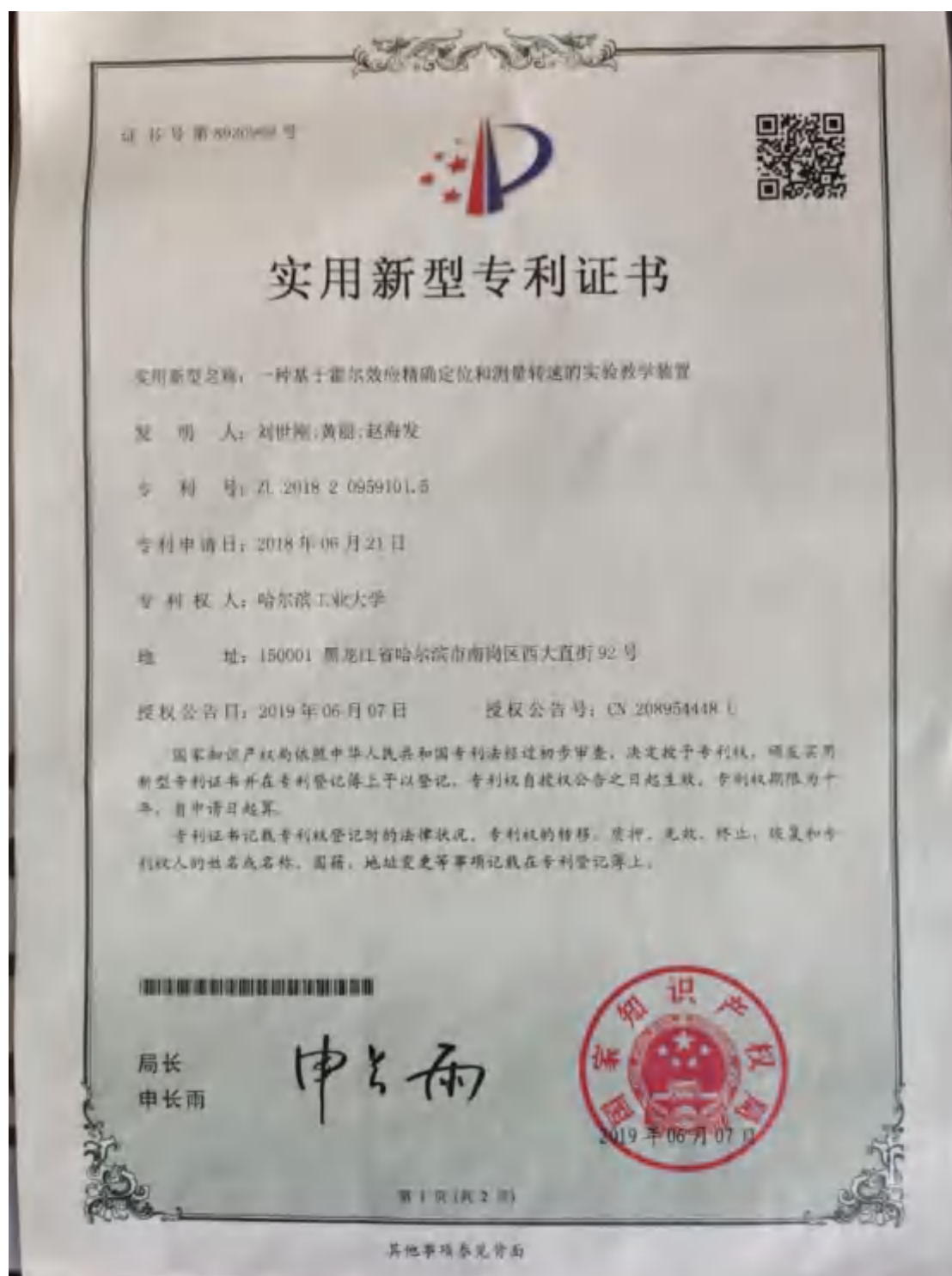
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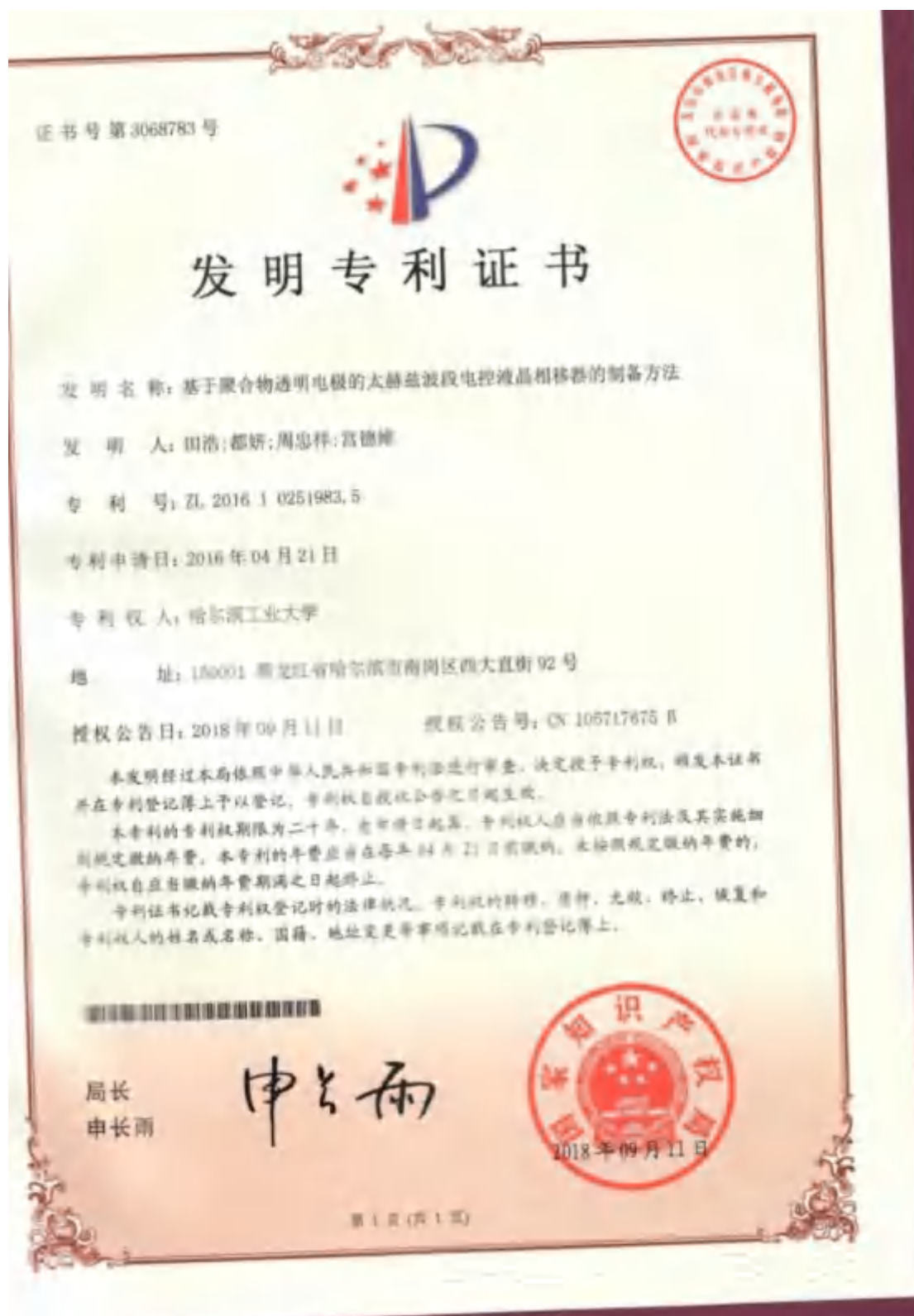
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(6) 学术竞赛等奖项佐证材料

| 序号 | 获奖名称 | 级别 | 获奖等级 | 学生姓名 | 年份 |
|----|------------------------------|-----|--------|---------------------|------|
| 1 | 2019年度中国光学学会光学优秀博士学位论文提名优秀论文 | 国家级 | 提名优秀论文 | 朱彤彤 | 2021 |
| 2 | 第十二届中国大学生物理学术竞赛 | 国家级 | 特等奖 | 骆全鑫、程湛博、厉彦榛、黄昕诺、岳琮沅 | 2021 |
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| 11 | 第七届全国大学生物理实验竞赛 | 国家级 | 一等奖 | 李磊、任睿溥、李俊杰 | 2021 |
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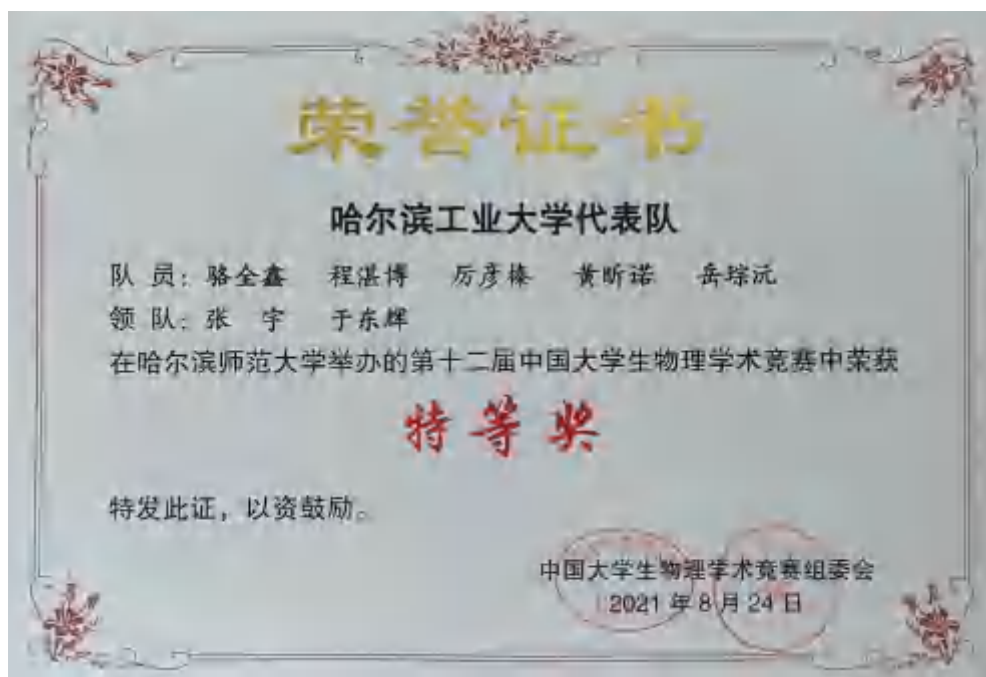
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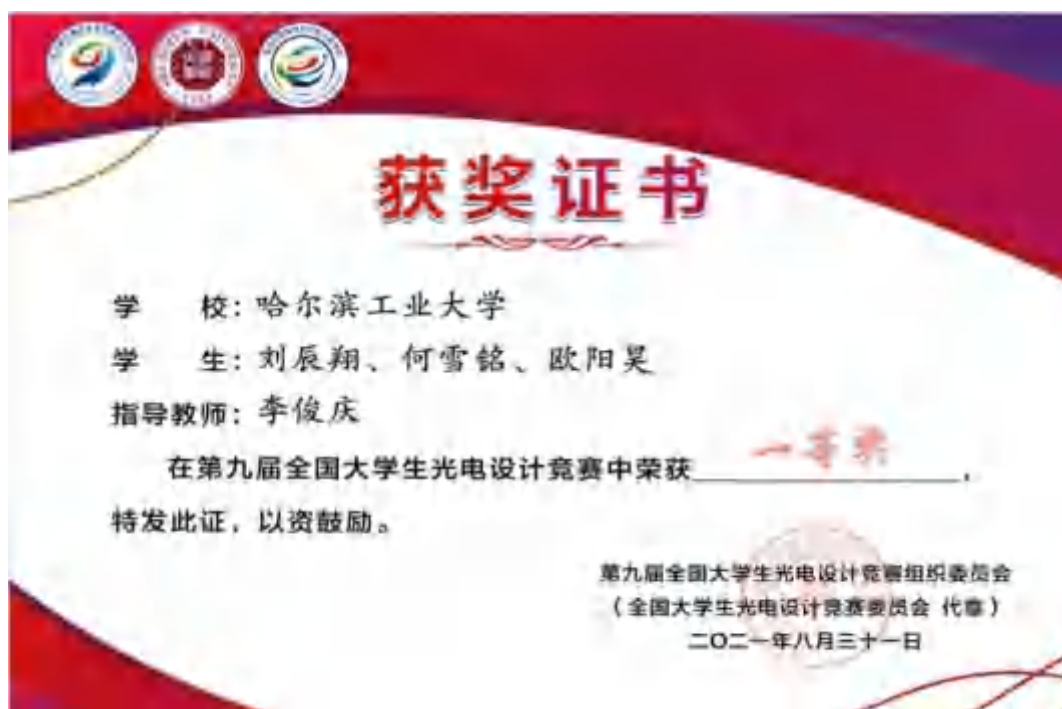
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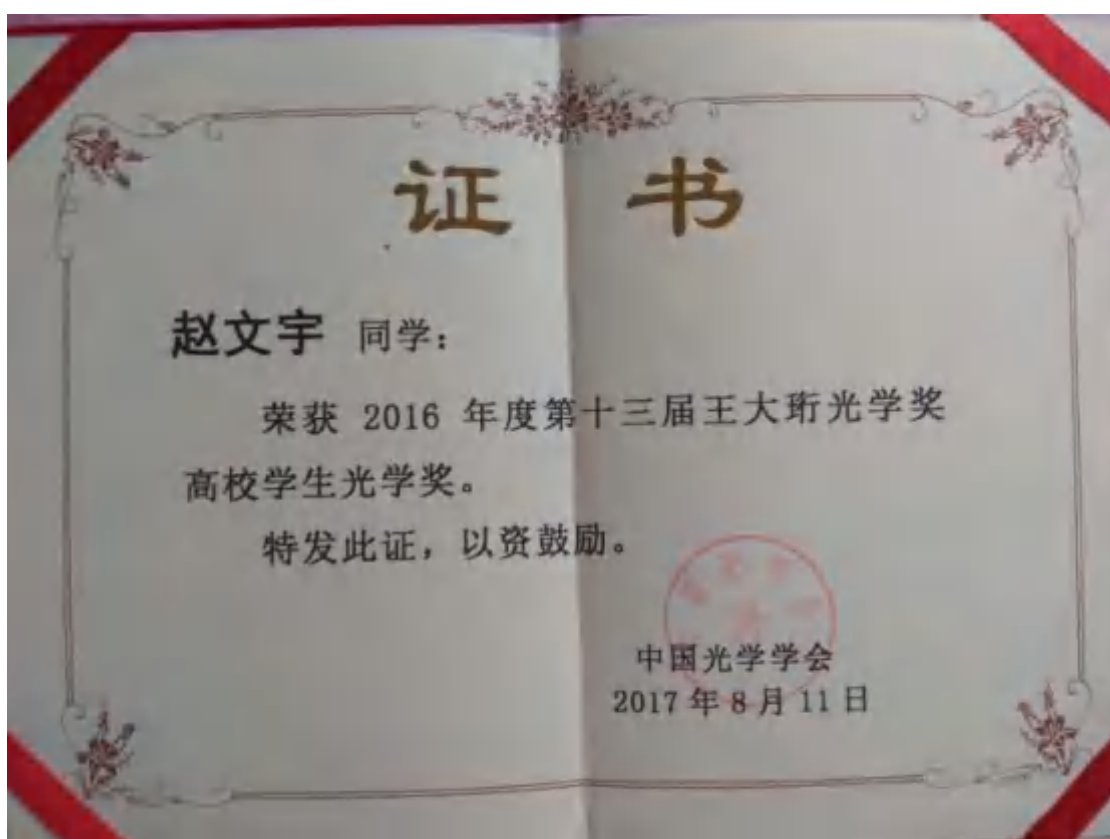
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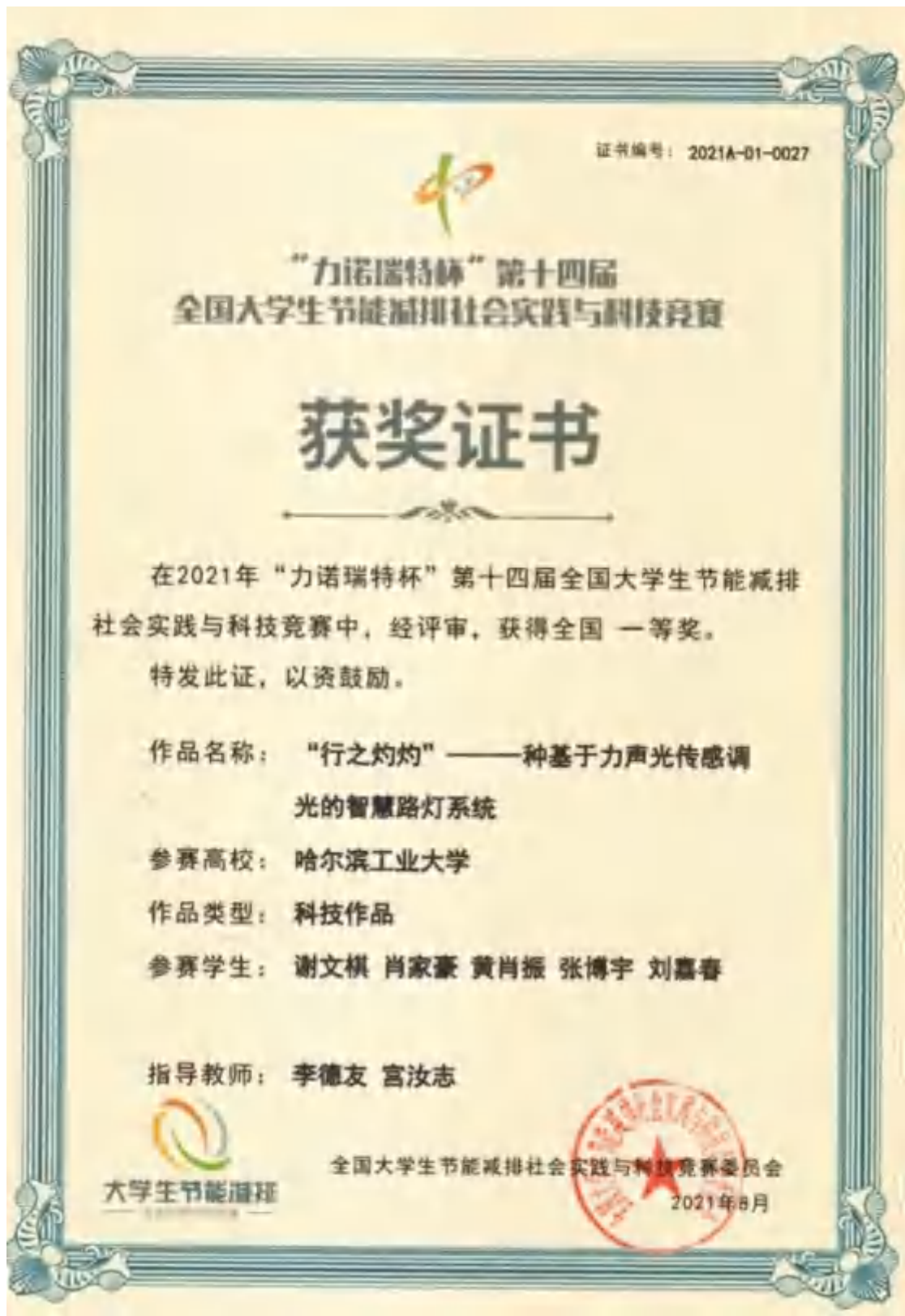
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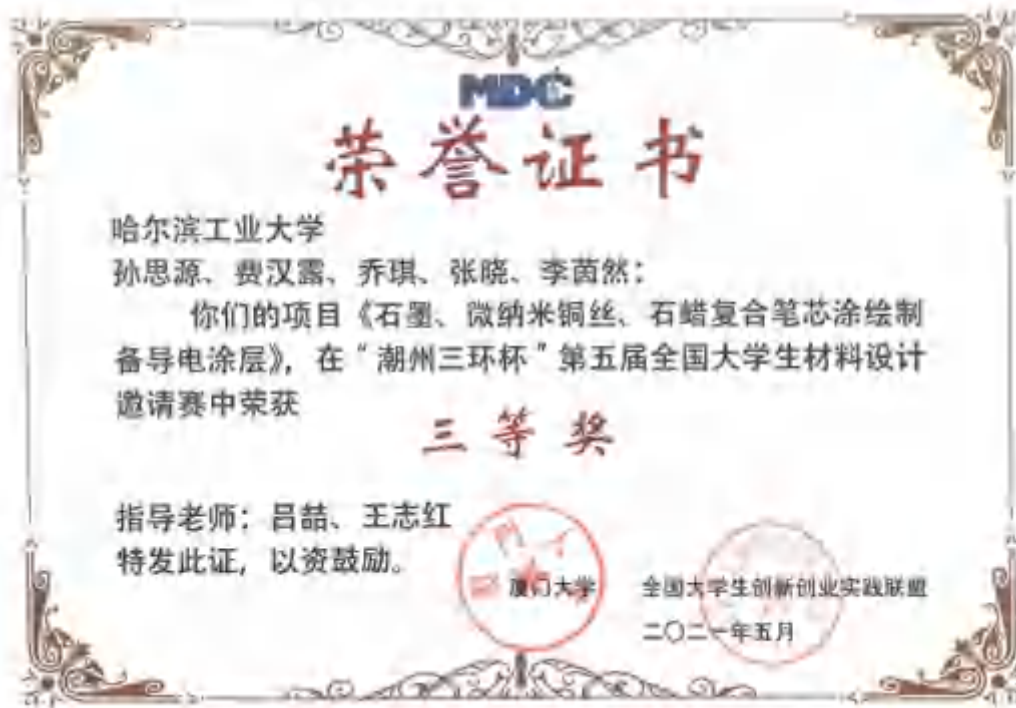
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(28) 第五届全国大学生材料设计邀请赛三等奖



(29) 全国光学与光学工程博士生学术联赛获奖证明





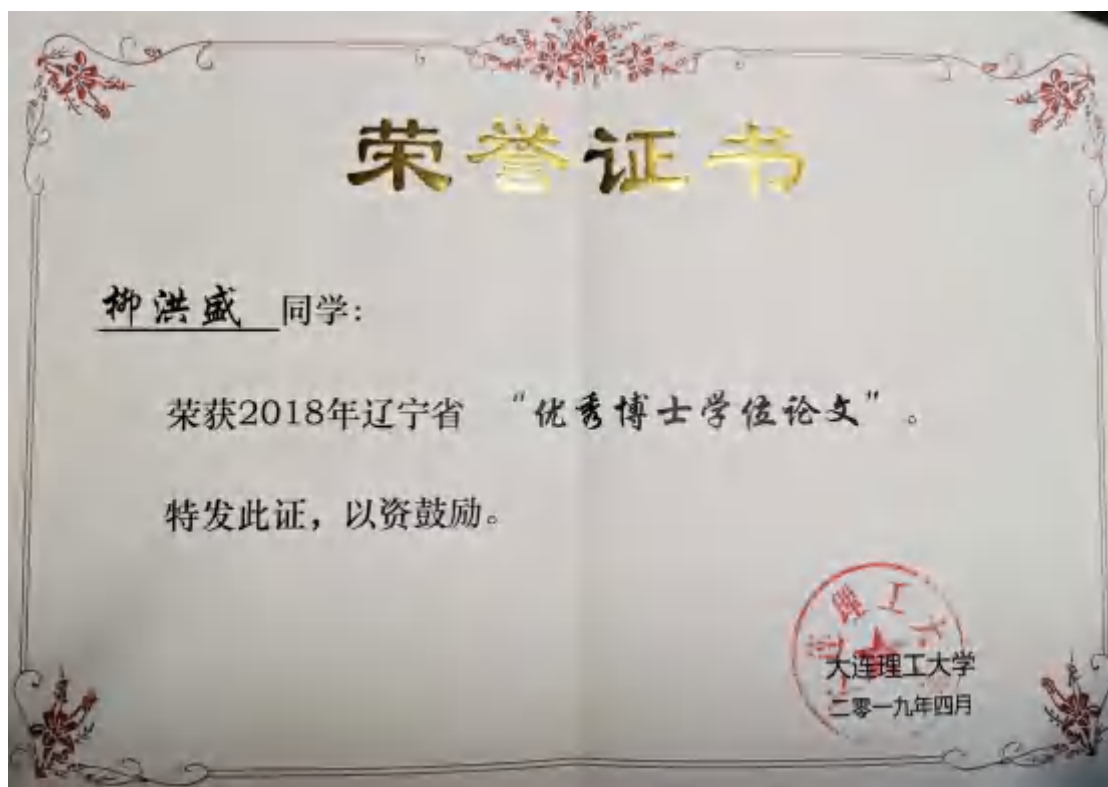
(30) 全国光学与光学工程博士生学术联赛获奖证明



(31) 辽宁省优秀博士学位论文证明



(32) 辽宁省优秀博士学位论文证明



(33) 辽宁省优秀博士学位论文证明



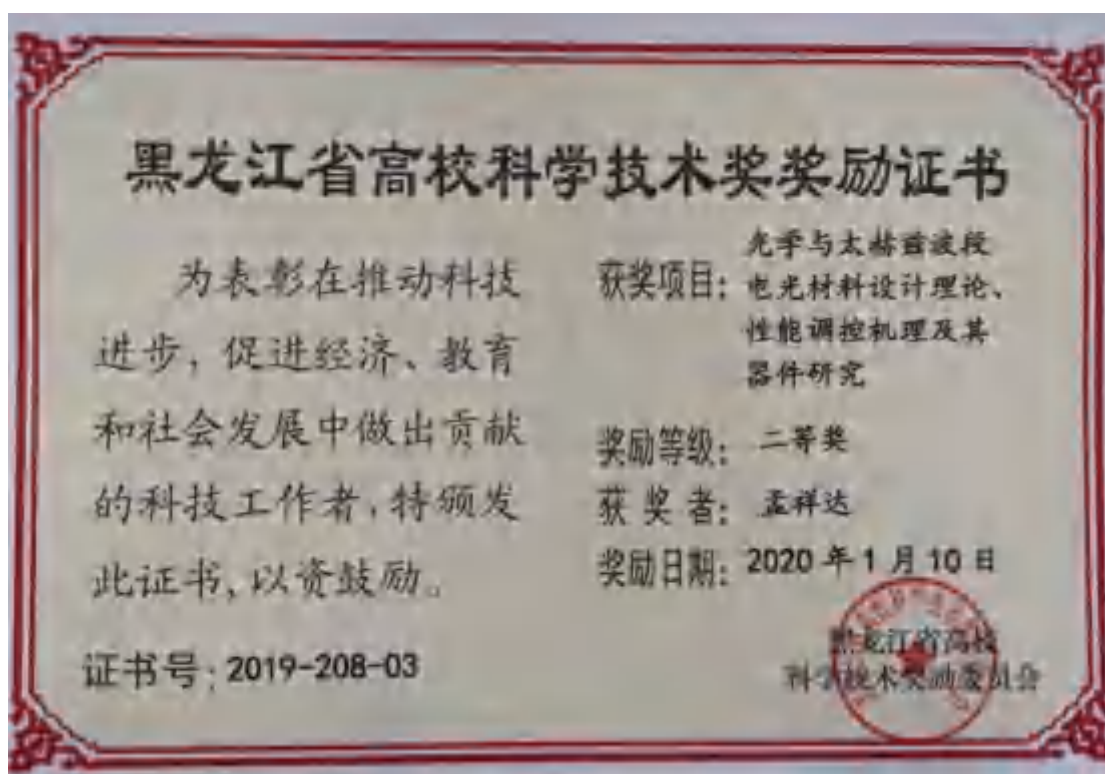
(34) 辽宁省优秀博士学位论文证明



(35) 辽宁省优秀博士学位论文证明



(36) 黑龙江省高校科学技术奖



(37) 第七届中国国际“互联网+”大学生创新创业大赛获奖证明



(38) 第十二届工银融e联“挑战杯”黑龙江省大学生创业计划竞赛



(39) 研究生国家奖学金（博士 42 人次）



(40) 研究生国家奖学金（硕士 26 人次）证明



(41) 辽宁省优秀毕业生证明



(42) 第五届研究生十佳团队证明



(7) 教学项目佐证材料

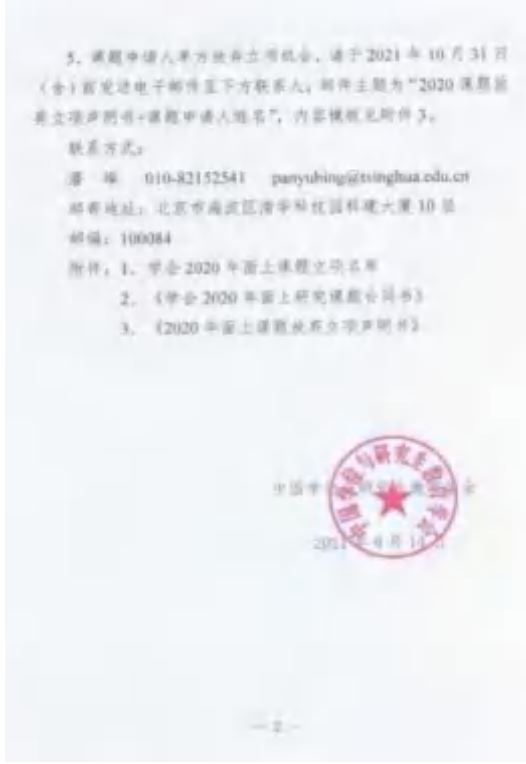
| 序号 | 项目名称 | 项目来源 | 负责人 | 年份 |
|----|-------------------------------------|-------------------------|-----|------|
| 1 | 物理学科研究生教学与科研融合研究与实践 | 中国学位与研究生教育学会 | 田浩 | 2020 |
| 2 | “一带一路”背景下哈工大对俄罗斯高校基础学科前沿科学研究提升计划 | 2021 年度教育部“一带一路”教育合作项目 | 李均 | 2019 |
| 3 | 依托国家级研究生联合培养实践基地高质量培养体系及方法研究 | 中国学位与研究生教育学会面上项目 | 李德友 | 2021 |
| 4 | 理论力学“一流课程”课程思政建设的探索与实践 | 教育部高等学校教学指导委员会 | 任延宇 | 2021 |
| 5 | 现代光电信息专业实验教学中心建设 | 教育部产学研合作协同育人项目 | 王健 | 2020 |
| 6 | 数学物理方法教学内容完善及教学体系优化设计 | 教育部高等学校物理学类专业教学指导委员会 | 姜向前 | 2016 |
| 7 | 创建以学生为主体的力学课程教学规范 | 教育部高等学校教学指导委员会 | 任延宇 | 2016 |
| 8 | 应用物理专业“科教融合”人才培养模式研究与实践 | 省高等教育教学改革一般研究项目 | 吕喆 | 2018 |
| 9 | 光电信息科学与工程专业资源共享课程建设—《光电信息功能材料》的教材建设 | 黑龙江省高等教育学会“十三五”高等教育科研课题 | 周忠祥 | 2016 |
| 10 | 基于人才个性化培养的基础物理分层次教学研究 | 黑龙江省高等教育学会“十三五”高等教育科研课题 | 孟庆鑫 | 2016 |
| 11 | 物理学科拔尖研究生的原始创新能力培养路径与培养实践 | 黑龙江省教育厅 | 丁卫强 | 2021 |
| 12 | 基于对俄合作提升研究生培养国际化水平 | 黑龙江省教育厅 | 袁承勋 | 2021 |
| 13 | 辽宁省一流线下课程：光波导技术 | 辽宁省教育厅 | 李淑凤 | 2022 |
| 14 | 国家本科教学工程项目：普通物理 | 中华人民共和国教育部 | 李淑凤 | 2021 |
| 15 | 以研究性实验为依托的物理拔尖人才创新能力培养研究与实验 | 辽宁省教育厅 | 赵纪军 | 2021 |
| 16 | 国家本科教学质量工程项目“基础拔尖人才”：王大珩班物理拔尖学生培养 | 中华人民共和国教育部 | 赵纪军 | 2020 |
| 17 | 国家本科教学质量工程项目强基计划：应用物理学 | 中华人民共和国教育部 | 赵纪军 | 2020 |
| 18 | 辽宁省线下一流课程：现代固 | 辽宁省教育厅 | 赵纪军 | 2020 |

| | | | | |
|----|---|----------------|-----|------|
| | 体物理学 | | | |
| 19 | 国家一流线上课程：大学物理 | 中华人民共和国教育部 | 李淑凤 | 2020 |
| 20 | 国家一流专业建设：应用物理学 | 中华人民共和国教育部 | 赵纪军 | 2019 |
| 21 | 中央高校教育教学改革专项，《固体物理学》特色课程群建设 | 中华人民共和国教育部 | 赵纪军 | 2019 |
| 22 | 大学物理 Mooc 练习活动资源建设及与传统教学融合模式的研究与实践 | 教育部高等学校教学指导委员会 | 李淑凤 | 2016 |
| 23 | 国家综合改革重点建设专业 | 辽宁省教育厅 | 赵纪军 | 2014 |
| 24 | 辽宁省教改项目“王大珩物理科学班”科研创新体系的全面建设 | 辽宁省教育厅 | 赵纪军 | 2016 |
| 25 | 辽宁省精品资源共享课：固体物理学 | 辽宁省教育厅 | 赵纪军 | 2015 |
| 26 | 辽宁省普通高等教育本科教学改革研究项目，“王大珩物理科学班”科研创新体系的全面建设 | 辽宁省教育厅 | 赵纪军 | 2015 |
| 27 | 辽宁省质量工程项目“应用物理学-优势特色专业” | 辽宁省教育厅 | 赵纪军 | 2014 |
| 28 | 辽宁省教改项目，《大学物理》跨校修读课程资源与教学模式优化共享的实践研究 | 辽宁省教育厅 | 李淑凤 | 2021 |
| 29 | 大学物理跨校修读课程教学模式的研究与实践 | 辽宁省教育厅 | 李淑凤 | 2018 |
| 30 | 校研究生教改重点项目“物理学研究生课程思政体系建设” | 大连理工大学 | 蒋雪 | 2021 |
| 31 | 校教改重点项目：中白“同窗友情”模式下深化“融合创新”合作办学的实践 | 大连理工大学 | 赵纪军 | 2022 |
| 32 | 校课程思政示范课：学术规范与科技英语写作 | 大连理工大学 | 赵纪军 | 2019 |
| 33 | 《计算凝聚态物理》课程教学与多元化育人模式探索 | 大连理工大学 | 蒋雪 | 2020 |

(1) 物理学科研究生教学与科研融合研究与实践



| 序号 | 课题名称 | 立项负责人 | 课题所属单位 |
|----|-------------|-------|--------|
| 1 | 基于...的...研究 | 张三 | 清华大学 |
| 2 | 基于...的...研究 | 李四 | 北京大学 |
| 3 | 基于...的...研究 | 王五 | 复旦大学 |
| 4 | 基于...的...研究 | 赵六 | 上海交通大学 |
| 5 | 基于...的...研究 | 孙七 | 浙江大学 |
| 6 | 基于...的...研究 | 周八 | 中国科技大学 |
| 7 | 基于...的...研究 | 吴九 | 南开大学 |
| 8 | 基于...的...研究 | 郑十 | 天津大学 |
| 9 | 基于...的...研究 | 冯十一 | 武汉大学 |
| 10 | 基于...的...研究 | 陈十二 | 华中师范大学 |
| 11 | 基于...的...研究 | 褚十三 | 湖南大学 |
| 12 | 基于...的...研究 | 褚十四 | 湖南大学 |
| 13 | 基于...的...研究 | 褚十五 | 湖南大学 |
| 14 | 基于...的...研究 | 褚十六 | 湖南大学 |
| 15 | 基于...的...研究 | 褚十七 | 湖南大学 |
| 16 | 基于...的...研究 | 褚十八 | 湖南大学 |
| 17 | 基于...的...研究 | 褚十九 | 湖南大学 |
| 18 | 基于...的...研究 | 褚二十 | 湖南大学 |
| 19 | 基于...的...研究 | 褚二十一 | 湖南大学 |
| 20 | 基于...的...研究 | 褚二十二 | 湖南大学 |
| 21 | 基于...的...研究 | 褚二十三 | 湖南大学 |
| 22 | 基于...的...研究 | 褚二十四 | 湖南大学 |
| 23 | 基于...的...研究 | 褚二十五 | 湖南大学 |
| 24 | 基于...的...研究 | 褚二十六 | 湖南大学 |
| 25 | 基于...的...研究 | 褚二十七 | 湖南大学 |
| 26 | 基于...的...研究 | 褚二十八 | 湖南大学 |
| 27 | 基于...的...研究 | 褚二十九 | 湖南大学 |
| 28 | 基于...的...研究 | 褚三十 | 湖南大学 |
| 29 | 基于...的...研究 | 褚三十一 | 湖南大学 |
| 30 | 基于...的...研究 | 褚三十二 | 湖南大学 |
| 31 | 基于...的...研究 | 褚三十三 | 湖南大学 |
| 32 | 基于...的...研究 | 褚三十四 | 湖南大学 |
| 33 | 基于...的...研究 | 褚三十五 | 湖南大学 |
| 34 | 基于...的...研究 | 褚三十六 | 湖南大学 |
| 35 | 基于...的...研究 | 褚三十七 | 湖南大学 |
| 36 | 基于...的...研究 | 褚三十八 | 湖南大学 |
| 37 | 基于...的...研究 | 褚三十九 | 湖南大学 |
| 38 | 基于...的...研究 | 褚四十 | 湖南大学 |
| 39 | 基于...的...研究 | 褚四十一 | 湖南大学 |
| 40 | 基于...的...研究 | 褚四十二 | 湖南大学 |
| 41 | 基于...的...研究 | 褚四十三 | 湖南大学 |
| 42 | 基于...的...研究 | 褚四十四 | 湖南大学 |
| 43 | 基于...的...研究 | 褚四十五 | 湖南大学 |
| 44 | 基于...的...研究 | 褚四十六 | 湖南大学 |
| 45 | 基于...的...研究 | 褚四十七 | 湖南大学 |
| 46 | 基于...的...研究 | 褚四十八 | 湖南大学 |
| 47 | 基于...的...研究 | 褚四十九 | 湖南大学 |
| 48 | 基于...的...研究 | 褚五十 | 湖南大学 |
| 49 | 基于...的...研究 | 褚五十一 | 湖南大学 |
| 50 | 基于...的...研究 | 褚五十二 | 湖南大学 |
| 51 | 基于...的...研究 | 褚五十三 | 湖南大学 |
| 52 | 基于...的...研究 | 褚五十四 | 湖南大学 |
| 53 | 基于...的...研究 | 褚五十五 | 湖南大学 |
| 54 | 基于...的...研究 | 褚五十六 | 湖南大学 |
| 55 | 基于...的...研究 | 褚五十七 | 湖南大学 |
| 56 | 基于...的...研究 | 褚五十八 | 湖南大学 |
| 57 | 基于...的...研究 | 褚五十九 | 湖南大学 |
| 58 | 基于...的...研究 | 褚六十 | 湖南大学 |
| 59 | 基于...的...研究 | 褚六十一 | 湖南大学 |
| 60 | 基于...的...研究 | 褚六十二 | 湖南大学 |
| 61 | 基于...的...研究 | 褚六十三 | 湖南大学 |
| 62 | 基于...的...研究 | 褚六十四 | 湖南大学 |
| 63 | 基于...的...研究 | 褚六十五 | 湖南大学 |
| 64 | 基于...的...研究 | 褚六十六 | 湖南大学 |
| 65 | 基于...的...研究 | 褚六十七 | 湖南大学 |
| 66 | 基于...的...研究 | 褚六十八 | 湖南大学 |
| 67 | 基于...的...研究 | 褚六十九 | 湖南大学 |
| 68 | 基于...的...研究 | 褚七十 | 湖南大学 |
| 69 | 基于...的...研究 | 褚七十一 | 湖南大学 |
| 70 | 基于...的...研究 | 褚七十二 | 湖南大学 |
| 71 | 基于...的...研究 | 褚七十三 | 湖南大学 |
| 72 | 基于...的...研究 | 褚七十四 | 湖南大学 |
| 73 | 基于...的...研究 | 褚七十五 | 湖南大学 |
| 74 | 基于...的...研究 | 褚七十六 | 湖南大学 |
| 75 | 基于...的...研究 | 褚七十七 | 湖南大学 |
| 76 | 基于...的...研究 | 褚七十八 | 湖南大学 |
| 77 | 基于...的...研究 | 褚七十九 | 湖南大学 |
| 78 | 基于...的...研究 | 褚八十 | 湖南大学 |
| 79 | 基于...的...研究 | 褚八十一 | 湖南大学 |
| 80 | 基于...的...研究 | 褚八十二 | 湖南大学 |
| 81 | 基于...的...研究 | 褚八十三 | 湖南大学 |
| 82 | 基于...的...研究 | 褚八十四 | 湖南大学 |
| 83 | 基于...的...研究 | 褚八十五 | 湖南大学 |
| 84 | 基于...的...研究 | 褚八十六 | 湖南大学 |
| 85 | 基于...的...研究 | 褚八十七 | 湖南大学 |
| 86 | 基于...的...研究 | 褚八十八 | 湖南大学 |
| 87 | 基于...的...研究 | 褚八十九 | 湖南大学 |
| 88 | 基于...的...研究 | 褚九十 | 湖南大学 |
| 89 | 基于...的...研究 | 褚九十一 | 湖南大学 |
| 90 | 基于...的...研究 | 褚九十二 | 湖南大学 |
| 91 | 基于...的...研究 | 褚九十三 | 湖南大学 |
| 92 | 基于...的...研究 | 褚九十四 | 湖南大学 |
| 93 | 基于...的...研究 | 褚九十五 | 湖南大学 |
| 94 | 基于...的...研究 | 褚九十六 | 湖南大学 |
| 95 | 基于...的...研究 | 褚九十七 | 湖南大学 |
| 96 | 基于...的...研究 | 褚九十八 | 湖南大学 |
| 97 | 基于...的...研究 | 褚九十九 | 湖南大学 |
| 98 | 基于...的...研究 | 褚一百 | 湖南大学 |



(2) “一带一路”背景下哈工大对俄罗斯高校基础学科前沿科学研究提升计划

2021年度“一带一路”教育合作项目清单

| 序号 | 项目名称 | 项目类别 | 实施单位 |
|----|------------|------|---------------------------|
| 1 | 北京交通大学 | 北京 | 中国工程院北京研究所 |
| 2 | 山东理工大学 | 山东 | 俄罗斯科学院远东分院俄国家科学院远东分院 |
| 3 | 山东理工大学 | 山东 | 俄罗斯科学院远东分院俄国家科学院远东分院 |
| 4 | 中国矿业大学(北京) | 北京 | 俄罗斯科学院远东分院俄国家科学院远东分院 |
| 5 | 山东理工大学 | 山东 | 俄罗斯科学院远东分院俄国家科学院远东分院 |
| 6 | 哈尔滨工业大学 | 广西 | “一带一路”沿线国家智库建设联盟(俄罗斯研究中心) |
| 7 | 广西民族大学 | 广西 | “一带一路”沿线国家智库建设联盟(俄罗斯研究中心) |
| 8 | 广西民族大学 | 广西 | “一带一路”沿线国家智库建设联盟(俄罗斯研究中心) |
| 9 | 哈尔滨工业大学 | 黑龙江 | “一带一路”沿线国家智库建设联盟(俄罗斯研究中心) |
| 10 | 黑龙江中医药大学 | 黑龙江 | 俄罗斯科学院远东分院俄国家科学院远东分院 |
| 11 | 哈尔滨工业大学 | 黑龙江 | 俄罗斯科学院远东分院俄国家科学院远东分院 |
| 12 | 哈尔滨工业大学 | 黑龙江 | “一带一路”沿线国家智库建设联盟(俄罗斯研究中心) |
| 13 | 哈尔滨工业大学 | 黑龙江 | 俄罗斯科学院远东分院俄国家科学院远东分院 |

(3) 依托国家级研究生联合培养实践基地高质量培养体系及方法研究

学位与研究生教育研究课题
合同书

课题编号: 2020RKA200

课题名称: 依托国家级研究生联合培养实践基地高质量培养体系及方法研究

课题委托方(甲方): 中国学位与研究生教育学会

课题受托单位(乙方): 哈尔滨工业大学

项目负责人: 李德友

课题实施单位:

起止日期: 2021年11月01日至2022年10月31日

中国学位与研究生教育学会
二〇二一年九月

七、合同签署

甲方:



 法定代表人(盖章)

乙方:



 课题受托单位(盖章)

(4) 理论力学“一流课程”课程思政建设的探索与实践

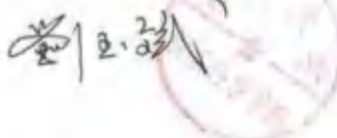
教育部高等学校物理学类专业教学指导委员会
2021 年度高等学校理论力学课程教学研究项目
立项通知

任延宇老师，您好：

您 2021 年申报的理论力学课程教学研究项目《理论力学“一流课程”课程思政建设的探索与实践》，经过答辩及专家讨论后已经立项，项目编号为 JZW-21-LL-06，请按照项目申报书和理论力学课程研究会要求认真执行，并及时提交中期报告和结题报告。

课程研究会感谢您的支持！

理事长：



教育部高等学校物理学类专业教学指导委员会
高等学校理论力学课程研究会
南开大学物理学院（挂靠单位（代章））

2021 年 8 月 30 日

(5) 现代光电信息专业实验教学中心建设

教育部高等教育司关于公布2019年第二批产学合作协同育人项目立项名单的通知

教高函函〔2020〕6号

各省、自治区、直辖市教育厅（教委），新疆生产建设兵团教育局，有关高等学校，有关企业：

为深入贯彻《国务院办公厅关于深化产教融合的若干意见》（国办发〔2017〕95号）精神，落实《教育部 工业和信息化部 中国工程院 关于加快建设新型工科 实施卓越工程师教育培养计划2.0的意见》（教高〔2018〕3号）要求，深化产教融合，精准合作，教育部会同有关企业和高校深入实施产学合作协同育人项目。

《教育部高等教育司关于公布有关企业支持的产学合作协同育人项目立项名单（2019年第二批）的通知》发布后，有关高校积极组织师生向企业提交项目申请，有关企业组织专家开展项目论证工作，并将经双方达成合作意向的项目向社会公示。经教育部产学合作协同育人项目专家组审定，现将立项项目名单予以公布（见附件）。

根据《教育部办公厅关于印发《教育部产学合作协同育人项目管理办法》的通知》（教高厅〔2020〕1号）要求，有关高校要加强项目后续的管理和跟踪，项目负责人要与相关企业加强联系，按照要求高质量推进项目实施，有关企业要确保资金及软硬件投入按时到位，规范项目管理，保证项目顺利实施。通过政府搭台、企业支持、高校对接、共赢共享，深化产教融合，促进教育链、人才链与产业链、创新链有机衔接，以产品和服务发展的最新需求驱动人才培养改革，培养支撑引领经济社会发展的新型专门人才。

附件：1.2019年第二批产学合作协同育人项目立项名单（按企业排序）

2.2019年第二批产学合作协同育人项目立项名单（按高校排序）

教育部高等教育司

2020年6月5日

附件1

2019年第二批产学合作协同育人项目立项名单（按企业排序）

| 项目编号 | 公司名称 | 项目类型 | 承担学校 | 项目名称 | 项目负责人 |
|-------------|-----------------|------------|------------|----------------------|-------|
| 20190201001 | 艾海康（北京）光电科技有限公司 | 线上课程类 | 中国传媒大学 | 通信与光电与微纳器件联合培养课程 | 李继军 |
| 20190201002 | 宇地星（北京）光电科技有限公司 | 线上课程类 | 浙江传媒学院 | 精工类课程线上教学可拓展实践教学平台构建 | 李继军 |
| 20190202046 | 武汉光通光电科技有限公司 | 虚拟仿真类课程类建设 | 湖南邮电职业技术学院 | 虚拟仿真实训课程建设及教学资源建设 | 李继军 |
| 20190202047 | 武汉光通光电科技有限公司 | 虚拟仿真类课程类建设 | 天津工业大学 | 虚拟仿真课程与教学资源建设及实训平台建设 | 武英新 |
| 20190202048 | 武汉光通光电科技有限公司 | 虚拟仿真类课程类建设 | 辽宁传媒大学 | 课程及线上虚拟仿真实训课程建设 | 李继军 |
| 20190202049 | 武汉光通光电科技有限公司 | 虚拟仿真类课程类建设 | 肇庆学院 | 课程及线上虚拟仿真实训课程建设 | 李继军 |
| 20190202050 | 武汉光通光电科技有限公司 | 虚拟仿真类课程类建设 | 山东财经大学 | 课程及线上虚拟仿真实训课程建设 | 李继军 |

(6) 数学物理方法教学内容完善及教学体系优化设计

关于高等学校数学物理方法课程
教学研究项目立项的通知

哈尔滨工业大学：

根据《关于高等学校数学物理方法课程开展教学研究项目立项的通知》(高物专教指字[2016]05号),你单位申报的项目《数学物理方法教学内容完善及教学体系优化设计》(课题负责人:姜向前),经专家委员会审定,同意立为高等学校数学物理方法课程2016年度资助项目,项目编号:JZW-16-SL-01。项目负责人接到本通知后,应立即开始启动项目,具体项目建设要求以《项目申请书》为准,并切实开展课题研究,按时结题。

教育部高等学校
物理学类专业教学指导委员会
(北京大学代章)  
2016年11月26日

(7) 创建以学生为主体的力学课程教学规范

关于高等学校力学课程 教学研究项目立项的通知

哈尔滨工业大学：

根据《关于高等学校力学课程开展教学研究项目立项的通知》(高物专教指字[2016]07号), 你单位申报的项目 《创建以学生为主体的力学课程教学规范》 (课题负责人: 任延宇), 经专家委员会审定, 同意立为高等学校力学课程 2016 年度资助项目, 项目编号: JZW-16-LX-09。项目负责人接到本通知后, 应立即开始启动项目, 具体项目建设要求以《项目申请书》为准, 并切实开展课题研究, 按时结题。

教育部高等学校
物理学类专业教学指导委员会
(北京大学代章) 
2016 年 11 月 26 日

(8) 应用物理专业“科教融合”人才培养模式模式研究与实践

黑龙江省教育厅

黑教高函〔2018〕608号

省教育厅关于公布2018年度高等教育教学改革研究项目立项结果的通知

各普通本科高校：

按照《省教育厅关于开展2018年度高等教育教学改革研究项目申报与结题的通知》（黑教高函〔2018〕281号）要求，我厅组织开展了2018年度高等教育教学改革研究项目立项工作，现将结果予以公布。其中，重点委托项目立项73项，一般项目备案586项。

请各高校认真按照《关于印发省教育厅关于进一步加强普通高等学校教育教学改革与研究的指导意见的通知》（黑教发〔2017〕9号）要求，加强对立项项目的支持和管理，对立项项目按照不低于1:1的标准足额落实配套经费，并指导项目主持人按期完成研究工作，培育更多优质教学成果。我厅将根据高校项目整体完成和推广情况，本着激励原则，动态调整下一年度立项指标。

| | | | | | | | | | |
|-----|---------|------------------------------|-------------|-----|-----|-----|-----|-----|--|
| 169 | 哈尔滨理工大学 | 基于400h“卓越计划”应用型人才培养模式创新及评价研究 | 40020180101 | 理华 | 李煜 | 孙明 | | | |
| 170 | 南京理工大学 | 基于学科交叉融合的“双一流”建设模式研究 | 10020180150 | 孙林江 | 何国栋 | 孙明 | 孙明方 | 吴建强 | |
| 171 | 哈尔滨理工大学 | 基于“双一流”精英人才培养的“双一流”建设模式研究 | 20020180151 | 孙林 | 孔强 | 刘国栋 | 孙明 | 孙明方 | |
| 182 | 哈尔滨理工大学 | “双一流”建设人才培养模式研究 | 50020180152 | 孙晓亮 | 孙林 | 杨威 | 孙林方 | 吴建强 | |
| 183 | 哈尔滨理工大学 | 应用型专业“科教融合”人才培养模式研究 | 50020180153 | 孙林 | 孙林 | 孙明 | 孙明方 | 吴建强 | |
| 184 | 哈尔滨理工大学 | 科教融合，协同育人，推进应用型人才培养模式研究 | 50020180154 | 孙林 | 孙林 | 孙明 | 孙明方 | 吴建强 | |
| 435 | 哈尔滨理工大学 | 应用型人才培养模式研究 | 50020180155 | 孙林 | 孙林 | 孙明 | 孙明方 | 吴建强 | |
| 436 | 哈尔滨理工大学 | 基于“双一流”建设人才培养模式研究 | 50020180156 | 孙林 | 孙林 | 孙明 | 孙明方 | 吴建强 | |

(9) 光电信息科学与工程专业资源共享课程建设—《光电信息功能材料》的教材建设

黑龙江省高等教育学会 文 件

黑高教会[2016]3号

关于下达高等教育科学研究“十三五”规划课题的 通 知

哈尔滨工业大学：

经学校初评推荐，省高等教育学会聘请专家组评审，学会秘书处审核，省高教学会批准贵校周忠祥同志申报的课题立项为黑龙江省高等教育学会高等教育科学研究“十三五”规划课题。现通知如下：

课题名称：光电信息科学与工程专业资源共享课程建设——《光电信息功能材料》的教材建设

课题类别：规划课题

课题编号：16G013

负责人：周忠祥

参加人：孟庆鑫、田浩、宫德维、申艳青、王莹

请贵校对课题加强管理，请课题负责人收到通知三个月内组织开题，按计划完成课题研究任务，争取出高质量的科研成果，为高等教育改革与发展做出贡献。

黑龙江省高等教育学会

2016年8月10日

(10) 基于人才个性化培养的基础物理分层次教学研究

黑龙江省高等教育学会 文 件

黑高教会[2016]3号

关于下达高等教育科学研究“十三五”规划课题的 通 知

哈尔滨工业大学:

经学校初评推荐,省高等教育学会聘请专家组评审,学会秘书处审核,省高教会批准贵校孟庆鑫同志申报的课题立项为黑龙江省高等教育学会高等教育科学研究“十三五”规划课题。现通知如下:

课题名称:基于人才个性化培养的基础物理分层次教学研究

课题类别:规划课题

课题编号:16G014

负责人:孟庆鑫

参加人:张宇、李均、骆素华、张伶俐、袁承勋、吴琦、韩权

请贵校对课题加强管理,请课题负责人收到通知三个月内组织开题,按计划完成课题研究任务,争取出高质量的科研成果,为高等教育改革与发展做出贡献。

黑龙江省高等教育学会

2016年8月10日

(11) 物理学科拔尖研究生的原始创新能力培养路径与培养实践

哈尔滨工业大学推荐申报2021年度黑龙江省高等教育教学改革研究项目立项名单

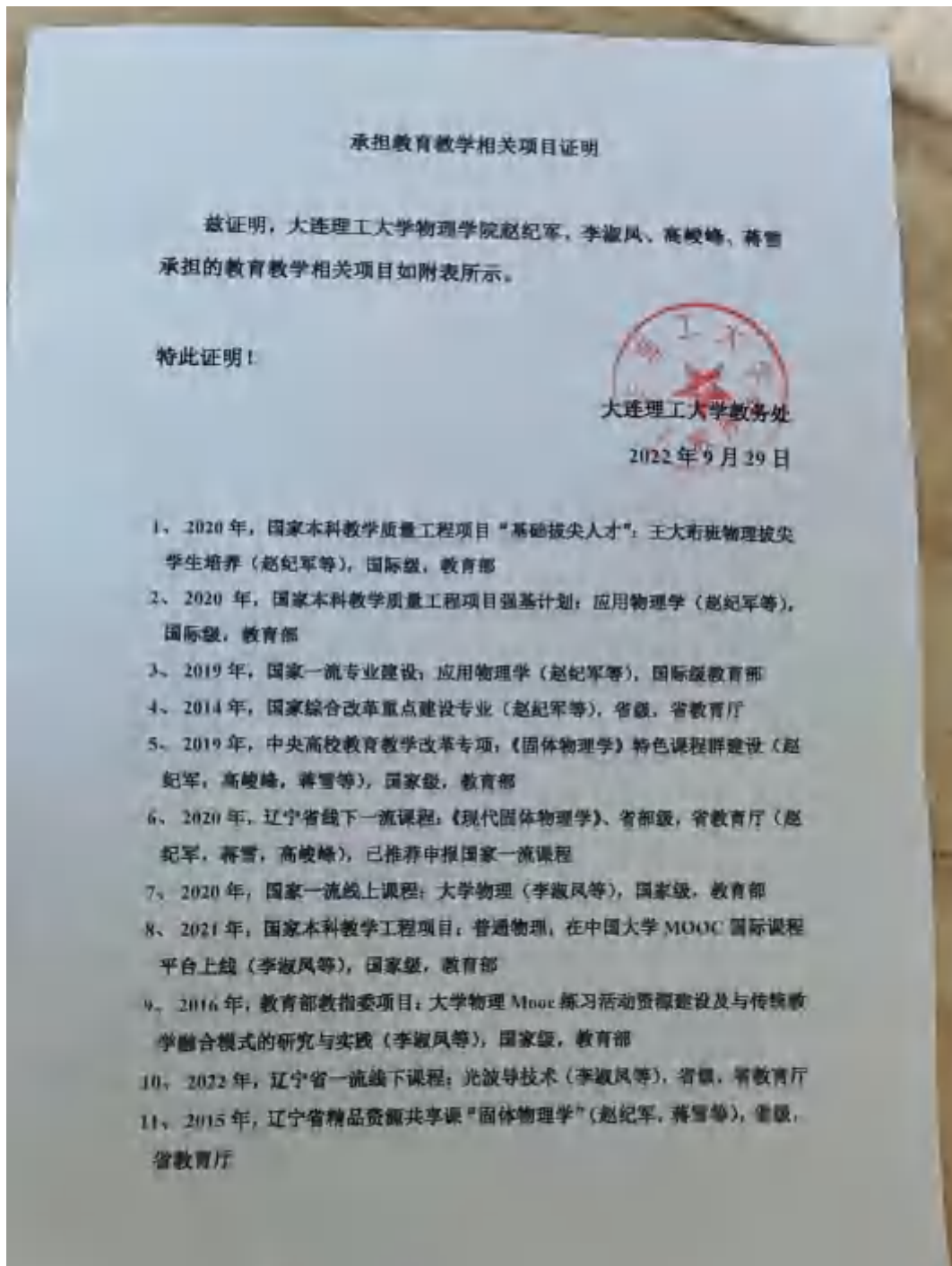
| 序号 | 学院名称 | 项目名称 | 项目负责人 | 项目参与人(共4人,按照排序填) | | | | 备注 |
|----|------|--------------------------------|-------|------------------|-----|-----|-----|--------|
| | | | | 1 | 2 | 3 | 4 | |
| 1 | 机电学院 | 基于融合-沉浸交互体验的可穿戴设备实训教学与实践 | 董明杰 | 付乃明 | 赵子祥 | 赵斌 | | 一般项目 |
| 2 | 机电学院 | 人工智能交叉学科类高层次人才工程训练创新能力培养的探索与实践 | 张磊 | 付强 | 何晓峰 | 曹华 | 李洪丹 | 一般项目 |
| 3 | 电信学院 | 基于物联网技术的中专子信类创新型人才培养模式探索与实践 | 刘国生 | 陈云 | 魏国元 | 李林 | 于雷 | 重点委托项目 |
| 4 | 机电学院 | 研究生创新能力培养新模式探索与实践 | 刘厚松 | 潘力 | 孙永军 | 张斌 | | 一般项目 |
| 5 | 材料学院 | 产学研协同育人实训基地建设运行机制 | 刘野 | 李峰 | 黄松波 | 谢立江 | 董志勇 | 重点委托项目 |
| 6 | 材料学院 | 面向理工科的课程与课程思政与思政课程建设与实践 | 刘国鑫 | 潘健 | 潘建群 | 韩立伟 | 李强 | 一般项目 |
| 7 | 机电学院 | “双中制”科技教育背景下高职院校研究生人才培养模式研究 | 马学忠 | 赵立军 | 董建民 | 肖国辉 | 李伟伟 | 一般项目 |
| 8 | 电气学院 | 三层次八阶段电气工程学科类一流人才培养能力培养与实践 | 白立刚 | 陈云 | 王福康 | 刘勇 | 胡军 | 一般项目 |
| 9 | 数学学院 | 数学学科高层次人才培养模式探索与实践 | 孙志安 | 洪恩英 | 胡宇明 | 傅法青 | 孙彬 | 一般项目 |
| 10 | 物理学院 | 基于过程化考核提升研究生培养国际化水平 | 姜泽旭 | 李宝 | 李功 | 王莹 | 姜海峰 | 一般项目 |
| 11 | 物理学院 | 物理学科拔尖研究生的原始创新能力培养路径与实践 | 丁立强 | 曹永超 | 孙宇翔 | 王博华 | 冯雷 | 一般项目 |


(12) 基于对俄合作提升研究生培养国际化水平

哈尔滨工业大学推荐申报2021年度黑龙江省高等教育教学改革研究项目立项名单

| 序号 | 学院名称 | 项目名称 | 项目负责人 | 项目参与人(共4人,按照排序填) | | | | 备注 |
|----|------|--------------------------------|-------|------------------|-----|-----|-----|--------|
| | | | | 1 | 2 | 3 | 4 | |
| 1 | 机电学院 | 基于融合-沉浸交互体验的可穿戴设备实训教学与实践 | 董明杰 | 付乃明 | 赵子祥 | 赵斌 | | 一般项目 |
| 2 | 机电学院 | 人工智能交叉学科类高层次人才工程训练创新能力培养的探索与实践 | 张磊 | 付强 | 何晓峰 | 曹华 | 李洪丹 | 一般项目 |
| 3 | 电信学院 | 基于物联网技术的中专子信类创新型人才培养模式探索与实践 | 刘国生 | 陈云 | 魏国元 | 李林 | 于雷 | 重点委托项目 |
| 4 | 机电学院 | 研究生创新能力培养新模式探索与实践 | 刘厚松 | 潘力 | 孙永军 | 张斌 | | 一般项目 |
| 5 | 材料学院 | 产学研协同育人实训基地建设运行机制 | 刘野 | 李峰 | 黄松波 | 谢立江 | 董志勇 | 重点委托项目 |
| 6 | 材料学院 | 面向理工科的课程与课程思政与思政课程建设与实践 | 刘国鑫 | 潘健 | 潘建群 | 韩立伟 | 李强 | 一般项目 |
| 7 | 机电学院 | “双中制”科技教育背景下高职院校研究生人才培养模式研究 | 马学忠 | 赵立军 | 董建民 | 肖国辉 | 李伟伟 | 一般项目 |
| 8 | 电气学院 | 三层次八阶段电气工程学科类一流人才培养能力培养与实践 | 白立刚 | 陈云 | 王福康 | 刘勇 | 胡军 | 一般项目 |
| 9 | 数学学院 | 数学学科高层次人才培养模式探索与实践 | 孙志安 | 洪恩英 | 胡宇明 | 傅法青 | 孙彬 | 一般项目 |
| 10 | 物理学院 | 基于过程化考核提升研究生培养国际化水平 | 姜泽旭 | 李宝 | 李功 | 王莹 | 姜海峰 | 一般项目 |
| 11 | 物理学院 | 物理学科拔尖研究生的原始创新能力培养路径与实践 | 丁立强 | 曹永超 | 孙宇翔 | 王博华 | 冯雷 | 一般项目 |

(13) 剩余部分教改项目证明材料



- 
- 12、2016年,辽宁省教改项目“王大野物理科学班”科研创新体系的全面建设(赵纪军等),省级,省教育厅
 - 13、2015年,辽宁省普通高等教育本科教学改革研究项目:“王大野物理科学班”科研创新体系的全面建设(赵纪军等),省级,省教育厅
 - 14、2014年,辽宁省质量工程项目:“应用物理学-优势特色专业”(赵纪军等),省级,省教育厅
 - 15、2021年,省教改项目:“《大学物理》跨校修读课程资源与教学模式优化共享的实践研究”(李淑凤等),省级,省教育厅
 - 16、2018年,省教改项目:大学物理跨校修读课程教学模式的研究与实践(李淑凤等),省级,省教育厅
 - 17、2021年,省教改项目:以研究性实验为依托的物理拔尖人才创新能力培养研究与实践,省级,省教育厅
 - 18、2019年,校课程思政示范课:学术规范与科技英语写作(赵纪军,蒋雪等),校级,大连理工大学
 - 19、2022年,校教改重点项目:中白“同窗友情”模式下深化“融合创新”合作办学的实践(高峻峰等),校级,大连理工大学
 - 20、2021年,校教改重点项目:物理学研究生课程思政体系建设(赵纪军,蒋雪),校级,大连理工大学
 - 21、2020年,校课程思政先导课:《计算凝聚态物理》课程教学与多元化育人模式探索(蒋雪,赵纪军),校级,大连理工大学

(8) 教学论文发表佐证材料

| 序号 | 文章名称 | 刊物名称及级别 /会议名称 | 第一作者 | 出版 年份 |
|----|--|---|------|----------|
| 1 | Geometric explanation of conic-section interference fringes in a Michelson interferometer. | American Journal of Physics (SCI 期刊) | 方光宇 | 2013 |
| 2 | Naked-eye observation of visible spectra using a transmission-grating-based spectrometer. | European Journal of Physics (SCI 期刊) | 林珊 | 2021 |
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(1) Geometric explanation of conic-section interference fringes in a Michelson interferometer., 2013 American Journal of Physics



Geometric explanation of conic-section interference fringes in a Michelson interferometer

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A simple geometric method based on wave-front analysis is developed to provide a rigorous explanation for the various interference fringes in a Michelson interferometer without the compensator plate. In view of the fact that a homogeneous plane of rays from a point light source becomes spherical as it passes obliquely through a glass plate, the wave-front deviation from a spherical one is obtained by calculating the spherical focal distance of the central ray. If the compensator plate is removed, the two central rays along the interfering paths have different spherical focal distances (SFDs). Therefore, the optical path difference can never be compensated with the movement of the reflection mirror. The wave-front difference or the optical path difference is determined to the accuracy of second order by comparing the two pairs of principal curvatures radii and the focal point-by-point subtraction of the optical path length difference (OPLD) is revealed. Numerical results support the theoretical analysis with great accuracy. *PACS: 01.40.+g, 42.15.Hw, 42.30.Fv*
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1. INTRODUCTION

The Michelson interferometer (MI) is a standard topic in undergraduate physics courses and laboratories. It is also a research device which can readily be modified to measure such quantities as refractive index and thickness.¹⁻⁷

The most commonly observed interference fringes in the MI are circular bands and parallel straight lines. If the central or noncentralized, the formation of which is analyzed in most textbooks.⁸ These fringes are formed in a standard MI setup with the compensator plate (CP) in place, the role of which is also well explained in texts.⁹ However, it is not an uncommon phenomenon among students that for an MI without the CP illuminated by a monochromatic light source, the optical path length difference could be compensated by the movement of either one of the two reflective mirrors. Therefore, students do not expect to see interference fringes with shapes of conic sections such as ellipses, hyperbolas, or parallel straight lines rather than circles when the compensator plate is removed, although this is a well known phenomenon among professionals.¹⁰

The formation of these conic-section fringes differs in principle from hyperbolic fringes produced with tilted reflective mirrors. Numerical calculations based on geometric optics can simulate the phenomena, which can also be verified with the help of commercial software such as *excel*.¹¹ Theoretical equations describing conditions for the appearance of different fringe shapes have also been derived in Ref. 5. Although the equations are correct and applicable, they are complicated and lack clear physical significance and consequently do not appeal to undergraduate students.

In this paper, we present a concise explanation of conic-section fringes by comparing the wave-fronts of the interfering waves and calculating their difference. The calculation is completed by representing each of the wave-fronts locally with an elliptic paraboloidal surface, whose principal curvatures

of maximum coincide with the primary and secondary foci of the central ray. Because the positions of the foci have simple analytical expressions, which are already introduced in high-school geometric optics, the method is easier for students to understand. In addition to its educational purpose, the geometric method can also be adapted for applications in areas such as imaging, thickness and refractive index measurement, wave-front sensing, and optical metrology.

II. THEORETICAL ANALYSIS

In a standard Michelson interferometer, as shown in Fig. 1, the compensator plate CP is a duplicate of the beam splitter BS, which guarantees that the beam reflected from the movable mirror M_1 and the beam from the stationary mirror M_2 pass through the same thickness of glass.

In its operation, without losing generality, we assume the light is emitted by a monochromatic point source. Three types of light sources and the corresponding interference fringes, such as an extended light source and localized fringes of equal inclination or equal thickness, can be analyzed in the same way.

The point light source S_0 emits a spherical wave, which is divided in amplitude into two parts by the BS. The two waves are then reflected by the mirrors and recombined by the BS, and interference fringes appear on the screen. If we remove the CP, the waves reflected by the BS pass on through the BS three times, while the transmitted wave passes through the BS only once. The two different optical paths are depicted in Fig. 2. The virtual point sources in the two arms are denoted by S_{10} and S_{20} , respectively.

In a homogeneous isotropic medium, each ray r - z raylight line and normal to the wave-front. The properties of the wavefront can be studied by observing the behavior of the ray normal to it.¹² The ray r - z raylight line has the spherical wavefront from S_0 , is affected by its passage through the glass

(2) Naked-eye observation of visible spectra using a transmission-grating-based spectrometer. 2021 European Journal of Physics

Naked-eye observations of visible spectra using a transmission-grating-based spectrometer

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Abstract

Naked-eye observations of visible spectra using a transmission-grating-based spectrometer, with the spectral resolution of 0.2 nm @ 550 nm, are reported in this article. The resultant accuracy of the wavelength readout with the naked-eye is better than 1 nm. The diffraction angles are precisely confirmed when the virtual diffraction image of the entrance slit (line spectrum) are collinear with two collimating slits. Furthermore, a converging lens is used to magnify the tiny light source (or narrow slit), and receive more radiant flux, with no effect on the accuracy of the spectral measurement. The spectrometer provides a good training aid for college students because (1) it is a low cost and handy demo experiment; (2) it has no optoelectronic sensor, it is completely naked-eye observation; (3) it helps students intuitively understand the diffracting and imaging processes; (4) the resolving power meets the requirements for teaching purposes.

Keywords: naked-eye observation, transmission-grating-based spectrometer, diffraction image, demonstrations experiment

(Some figures may appear in colour only in the online journal)

1. Introduction

The operating principle of widely used blazed reflection grating based spectrometer (Czerny–Turner design) is as follows [1]: the target light rays, after the entrance slit, are reflected and collimated by a concave mirror, and then decomposed by the grating. The deflected parallel

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(3) 工程类专业学位研究生培养的治理改革——基于政策供需协调视角的多案例比较研究

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工程类专业学位研究生培养的治理改革^{*} ——基于政策供需协调视角的多案例比较研究

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摘 要:大国竞争时代的科技人才需求强化了工程类专业学位研究生教育的重要性与迫切性,但培养制度与社会需求的协调性仍然不足。如何实现供需协调是亟需思考的“瓶頸式”命题。首先基于供需侧“教学-科研-实践”三角模式和供给侧“劳动力、资本和技术”侧三角模式构建供需匹配的新治理模型;然后通过对案例比较从制度设计和制度执行两个层面构建实践新模式;最后进行理论模型与实践模式匹配,运用“供-需”和“理论-实践”二维框架解释供需协调的内在逻辑。研究表明,工程类专业学位研究生培养的治理改革可以围绕管理授权、导师队伍、实践基地、治理机制和经费方案五个方面进行实施,形成产教融合新格局,完善研究生教育数字化转型治理基础。

关键词:研究生教育;工程类专业学位;培养体系改革;供需协调;多案例研究

一、问题的提出

2020年以来美国陆续发布的《国防授权法案》均旨在通过天河超算科学技术基金会推动技术转化,增加研发经费,支持重点领域研究生科研活动,创新实践平台建设和技术商业化。中美战略竞争愈演愈烈,科技创新、技术产业化成为大国竞争的主要方式^①,是支撑关键科学技术全球领先的基本前提。注重工程类专业学位研究生培养,推动产学研融合,是科技创新与成果转化和人才供需协调的重要方式^②,将研究生培养与市场联系起来是创新变革的关键因素^③,是推动高水平科技自立自强的人才供给保障^④。

中国经济实力的突破需要优化科技创新需求和人力资本投入的协调性^⑤。从供给侧视角来看,中国科学技术已经逐渐走向前沿,模仿创新已经难以满足发展需求,市场中工程技术人才的需求结构更加复杂^⑥。为服务国家重大战略和重大社会挑战,增加专业学位研究生培养规模和提升培养质量是一种必然趋势。中国专业学位研究生的招生和毕业规模处于逐年增加的趋势,原有培养体系不能适应招生规模扩张后的培养压力,难以满足市场对人才质量的

要求,倒逼研究生教育数字化转型的有紧迫性。因此进行工程类专业学位研究生培养体系的改革已经成为一个重要议题。

1-1设计2:工程类专业学位研究生的培养基础

宏观视角下,国家发展战略和市场需求动态变化的,需求重心可能会随领域热点的变化而发生转移^⑦;微观视角下,不同专业、岗位和企业对研究生的职业素养需求也存在差异,因此培养制度的供给需要具有动态匹配特性^⑧。然而,供需要素的动态匹配需要有明确的需求定包作为方向引导。在工程教育领域中,由“教学-科研-实践”构成的三角模式能够概括工程类专业学位研究生的培养导向^⑨,每个具体的培养目标都可以依据“教学-科研-实践”三个方面进行调整,以适应不同的培养需求。

但随着研究生规模扩大和工程技术人才素质要求的提高,培养制度供给出现了无效供给、供给质量低、供需结构失调等问题^⑩。具体问题可从三角模式的三个方面展开。

其一,课程教学方面存在组织管理不独立,教育治理专业性不强,培养目标模糊的问题。工程专业

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(5) 半无界空间上函数的傅里叶-贝塞尔积分展开. 大学物理

半无界空间上函数的傅里叶-贝塞尔积分展开

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摘要: 在半无界空间中, 本征函数族只由本征函数族所规定, 因此可作为广义傅里叶级数展开的基础. 本文从定态薛定谔方程和 $f(\rho)$ 上函数的广义傅里叶级数展开出发, 利用贝塞尔函数的递推展开公式以及贝塞尔函数零点的递推公式, 给出了半无界空间上函数的傅里叶-贝塞尔积分展开问题, 得到了本征函数族为特选归一化公式. 当 μ_n 趋于无穷时, 不能使函数族收敛地表示, 得到了函数的傅里叶-贝塞尔积分展开公式.

关键词: 贝塞尔函数; 傅里叶-贝塞尔积分

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在轴对称中, 当给定柱坐标系边界条件解关于 ρ 的贝塞尔方程与边界条件构成本征值问题. 此时本征函数为贝塞尔函数 $J_n(\sqrt{\mu_n} \rho)$, 其中 μ_n 为本征值由边界条件确定, 本征函数族 $J_n(\sqrt{\mu_n} \rho)$ 构成完备正交系, 由 \cos - Laure 定理本征值问题的性质可知, 函数一定定义在有界区间 $[0, \mu_n]$ 上的分段光滑函数 $f(\rho)$ 可展开为本征函数族为^[1] 傅里叶级数, 这一结果可在任一标准的数学物理方法教材上找到^[2]. 由于 $f(\rho)$ 是定义在有界区间上的函数, 但关于半无界空间情况和 $\rho_n \rightarrow \infty$, $f(\rho)$ 是否仍可展开为广义傅里叶级数这个问题一般的教材并未提及. 在继续读完的《数学物理方法》中有提及^[3], 但只是给出了结果. 结果表明当 $\rho_n \rightarrow \infty$ 时傅里叶-贝塞尔级数展开变为傅里叶-贝塞尔积分, 傅里叶-贝塞尔积分展开的严格数学证明单从存在但过于复杂并不适合课堂上介绍, 是后存在一种简单但值得借鉴的由这一结果是本文将要解决的问题. 本文针对定义在半无界空间上的分段光滑函数, 利用贝塞尔函数的递推展开公式和贝塞尔函数零点的递推公式给出傅里叶-贝塞尔积分展开及其系数公式.

由贝塞尔函数的性质可知, 定义在有界区间 $[0, \mu_n]$ 上的分段光滑函数 $f(\rho)$ 可展开为广义傅里叶级数

$$f(\rho) = \sum_{n=1}^{\infty} A_n J_n(\sqrt{\mu_n} \rho) \quad (1)$$

其傅里叶系数 A_n 为

$$A_n = \frac{1}{N_n} \int_0^{\mu_n} f(\rho) J_n(\sqrt{\mu_n} \rho) \rho d\rho \quad (2)$$

其中 N_n 为归一化, 考虑 $\mu_n \rightarrow \infty$ 的情况, 则方程 (2) 可写为

$$A_n = \lim_{\mu_n \rightarrow \infty} \sum_{n=1}^{\infty} \left[\frac{1}{N_n} \int_0^{\mu_n} f(\rho) J_n(\sqrt{\mu_n} \rho) \rho d\rho \right] J_n(\sqrt{\mu_n} \rho) \quad (3)$$

贝塞尔函数的积分为

$$N_n = \frac{1}{2} \left(\rho J_n' \left(\frac{\mu_n}{\rho} \right) \right) \Big|_0^{\mu_n} = \frac{1}{2} \rho J_n(\sqrt{\mu_n} \rho) \quad (4)$$

对于贝塞尔函数其渐近公式为

$$J_n(x) \sim \sqrt{\frac{2}{\pi x}} \cos \left(x - \frac{n\pi}{2} - \frac{\pi}{4} \right) \quad (5)$$

由于渐近公式的 $O(x^{-1/2})$ 项可以舍弃, 故其求导后项舍 $O(x^{-3/2})$ 项亦可舍弃

$$J_n(x) \sim -\sqrt{\frac{2}{\pi x}} \sin \left(x - \frac{n\pi}{2} - \frac{\pi}{4} \right) \quad (6)$$

令 $x = \sqrt{\mu_n} \rho$, 则 $x_n = \sqrt{\mu_n} \rho_n$, 由方程 (5) 和 (6) 可得

$$J_n(x_n) \sim \sqrt{\frac{2}{\pi x_n}} \cos \left(x_n - \frac{n\pi}{2} - \frac{\pi}{4} \right) \quad (7)$$

$$J_n'(x_n) \sim -\sqrt{\frac{2}{\pi x_n}} \sin \left(x_n - \frac{n\pi}{2} - \frac{\pi}{4} \right) \quad (8)$$

将方程 (7) 和 (8) 代入方程 (4) 可得

$$A_n \sim \frac{\rho_n}{\pi x_n} \left[1 - \frac{n}{x_n} \cos \left(x_n - \frac{n\pi}{2} - \frac{\pi}{4} \right) \right] \quad (9)$$

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(6) 熵概念的延拓——从热熵到信息熵. 大学物理

熵概念的延拓——从热熵到信息熵

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摘要 本文首先介绍了热力学熵、信息熵的概念的建立过程,并基于熵在信息论中的应用实例找出热力学熵与信息熵两者之间的关联。最后通过对普朗克熵公式的推广,并结合现代实验验证过程论证了信息熵与热力学熵之间的联系。

关键词 热力学熵; 信息熵; 普朗克熵公式; 熵的关联

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在物理学中的热力学中,熵(Entropy)是描述热力学系统状态的物理量,这也是普朗克最初创造的目的。借鉴于物理学中熵的概念,香农为了定量描述一条消息的信息量而在通信领域引入信息熵。本文将探讨热力学第二定律的成果能否应用到信息领域,即热力学熵与信息熵的相关性。

1 热力学中的熵

1.1 热力学中重要的定律反定理——熵概念的状态

人们最初研究热力学的目的是为了应用热力学的原理制造某种机器代替人力对外做功,所以便有了人们对热机的研究^[1]。说到热机便不得不提 1824 年法国工程师卡诺在他的论文《关于热力的几点说明》中提出的一种理想热机——卡诺热机^[2],这种热机的工作循环由两个准静态等温过程和两个准静态绝热过程组成。卡诺同时也证明了这种热机的效率最大,为热力学效率指明了方向。由此引出的卡诺定理也为后来的熵概念的提出埋下了伏笔。

自然界中发生的所有热力学过程都必须遵循能量守恒和守恒定律,卡诺循环在不违背热力学第一定律的情况下出现了可逆与不可逆这两个关键词。这个问题研究的便是热力学过程发生的方向性问题了,解决这个问题的便是热力学第二定律。

在卡诺定理的基础上,1870 年克劳修斯在其论文《关于热动力》[3]1851 年开尔文在其论文《热动力

理论》中分别提出了热力学第二定律^[4]。

克劳修斯表述:不可能使热量从低温物体传到高温物体而不引起其他变化;

开尔文表述:不可能从单一热源吸取热量,将其完全变为有用的功而不引起其他变化。

1.2 克劳修斯熵

随着对热力学第二定律的深入研究,人们发现一切与热现象有关的实际过程都是不可逆的,一切自然过程都有其方向性问题,这就给热力学系统所进行的不可逆过程的初态和终态之间存在巨大的差异性^[5]。这种差异决定了热力学过程的方向性。为了研究这种差异性,我们更应该对热力学系统的状态进行定性或定量的描述。

克劳修斯首先达成了这个成就。在对卡诺定理及热力学第二定律的研究中,克劳修斯推出了克劳修斯等式^[6]。在可逆循环中,从一个状态到另一个状态的过程中, $\int \frac{\delta Q}{T} = 0$ (称为热温比), $\int \frac{\delta Q}{T}$ 的积分与过程无关,只与初末状态有关。这在物理意义上说明系统中存在一个状态函数,1855 年克劳修斯在其论文《力学的热理论的必要方程之便于应用的形式》命名这种状态函数为熵^[7],可逆循环中就有 $\int \frac{\delta Q}{T} = S_2 - S_1$ 。

这是熵在历史舞台上的首次露面,而且熵是系统状态函数这一基本点为后来所有的熵理论的延拓

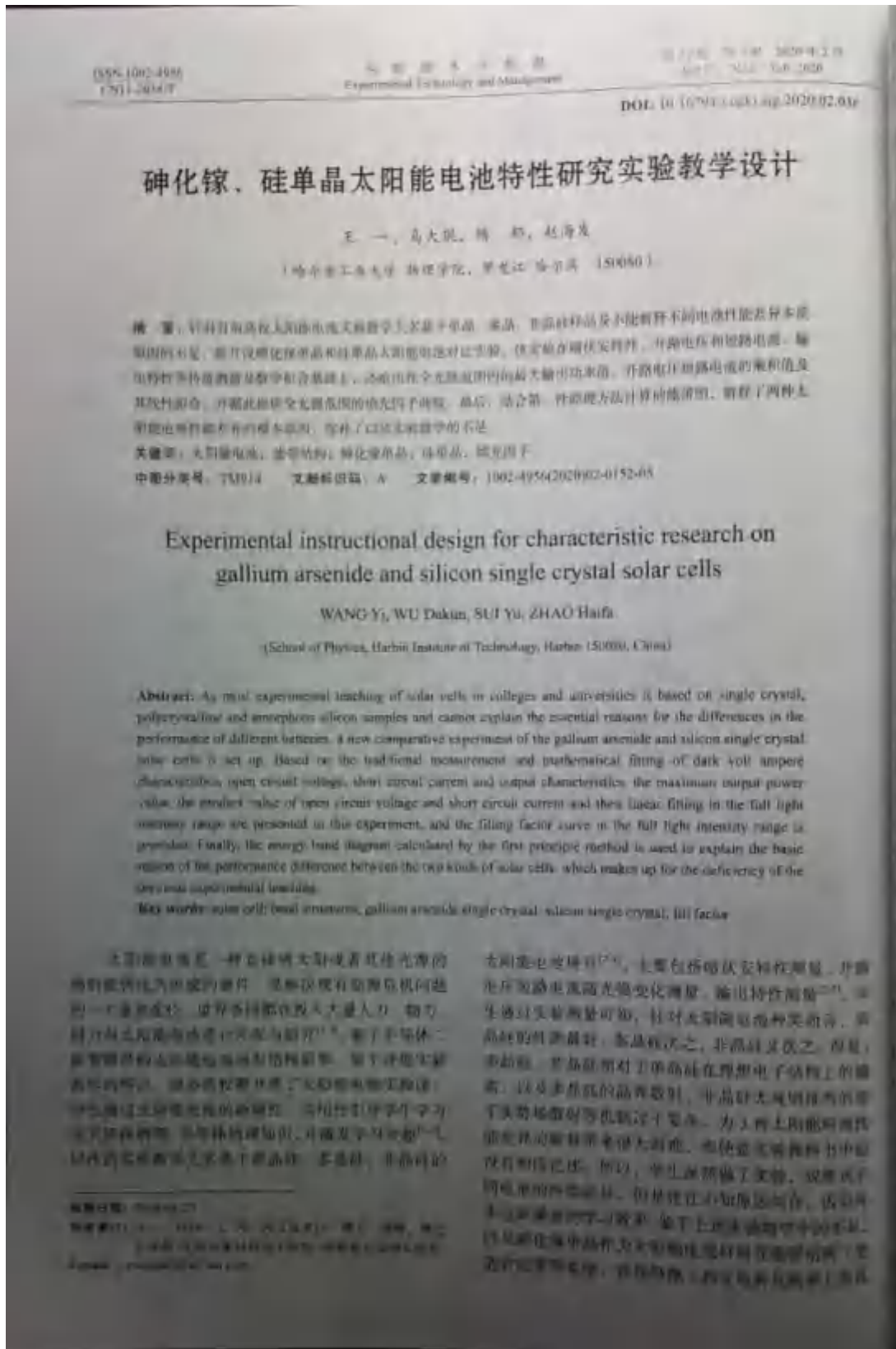
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(7) 砷化镓、硅单晶太阳能电池特性研究实验教学设计. 实验技术与 管理



(9) 用低功率可见二极管激光器进行倍频实验. 实验技术与管理

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虚拟仿真技术

用低功率可见二极管激光器进行倍频实验

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摘 要:针对用高功率近红外固体激光器进行激光倍频实验存在光路调节难度大、安全防护要求高的问题,提出一种基于低功率连续可见二极管激光器的倍频实验方案。使用功率小于 50 mW、中心波长 780 nm 的二极管激光光作为基频光,利用偏硼酸钡晶体生成了 4 W 激光功率的倍频紫外激光,并由光电倍增管探测。实验结果表明倍频光功率与基频光功率为二次函数关系,倍频效率与基频光功率为一次函数关系。实验结果与理论值高度吻合。该实验方案具有光路校准难度小和安全防护要求低的优点。

关键词: 二极管激光; 倍频实验; 实验教学

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Experiment of frequency doubling with low power visible diode laser

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Abstract: In view of the difficulty of adjusting the optical path and the high requirement of safety and protection in the laser frequency doubling experiment, a high-power near-infrared solid-state laser, as a source of the frequency doubling experiment based on the low-power continuous visible diode laser is proposed. By using a diode laser with a power of less than 50 mW and a central wavelength of 780 nm as the fundamental frequency light, a frequency-doubled ultraviolet laser with a power of 4 W is generated by BaB₂O₇ crystal, and is detected by an efficient photomultiplier tube. The experimental results show that the power of doubled frequency light is the quadratic function with the power of fundamental frequency light, and the efficiency of doubled frequency light is the quadratic function with the power of fundamental frequency light. The experimental results are in good agreement with the theoretical results. The proposed experimental scheme has the advantages of low difficulty in optical calibration and lower requirements for safety protection.

Key words: diode laser; frequency doubling experiment; experimental teaching

非线性光学是现代光学的一个重要分支,利用非线性效应获取短波光已成为激光技术研究的一个热点方向^[1]。在大学物理实验和近代物理实验等课程中开设的激光倍频实验^[2],多采用波长 1 064 nm 的高功率近红外固体激光器^[3]作为光源。然而,高功率

不可见光源在实验中存在光路调节难度大和安全防护要求高的问题。

本文提出一种采用低功率、连续可见二极管激光(LD)为光源的倍频实验方案。采用功率小于 50 mW 的 780 nm 绿色 LD 光源作为基频光^[4],使用偏硼酸钡(BBO)晶体生成 4 W 倍频紫外激光,并由光电倍增管(PMT)探测。基频光为低功率可见光,大大降低了倍频实验的光路校准难度和安全防护要求。同时,4 W 紫外光探测实验可进一步丰富学生关于微弱光信号检测方面的知识。

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浅谈利用薛定谔方程计算量子态的时间演化

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摘要:量子力学教学中,薛定谔方程是描述一个量子系统变化的核心部分.学生对薛定谔方程的学习,可以通过量子物理和经典物理的不同对比,在量子物理教学中,薛定谔方程的讲解是一个非常重要的内容.因为在量子力学中对于薛定谔方程的理解,需要用到初态薛定谔方程,而对于量子态随时间的变化部分并不清晰,因此我们引入概率幅模型.一个量子态在一个耦合的系统中,量子态随时间的变化部分并不清晰,在薛定谔方程中利用薛定谔方程描述量子态随时间的变化过程,给出几类随时间变化的薛定谔方程,从而更加直观的理解薛定谔方程在一个量子系统中的应用.

关键词:量子系统; 薛定谔方程

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作为物理学研究前沿的问题,在经典物理中存在许多尚未解决的科学问题,许多科研人员仍在研究基本的物理现象^[1],然而对于微观世界的规律,则需要用量子力学来描述^[2].量子力学和相对论是现代物理学的重要组成部分,从教学的角度而言,学习量子力学的那部分学生觉得量子力学是一个不重要的学科,甚至在未来的工程应用或科学研究中尝试而忽略带有量子的概念.随着科学技术的进展,更多的学科需要考虑到量子效应,而且通过引入量子效应,能取得非常大的进展.例如,利用量子系统实现的量子计算机^[3],能够解决经典计算机不能解决的问题.基于量子态的通信任务,能够实现完全安全的通信方案^[4].此外基于量子态,对目标进行探测,也可极大的提高探测的精度^[5].由此可见,量子物理会在未来的科学研究中占有越来越重要的地位^[6-11].

在大学物理的教学过程中,应结合技术与教育教学深度融合.本文主要针对薛定谔方程描述一个量子系统的演化问题.微观粒子的定态薛定谔方程已经有很特殊的例子,例如势阱中的微观粒子的计算.而学生对薛定谔方程中的时间演化部分,学习效果较差.因此,我们引入一个概率幅模型,该模型是量子力学中的一个重要的模型^[12].然而该模型涉及到量子力学相应的知识点,而本文只简要的推

导和表示,从简单的物理过程描述,建立量子态随时间变化的方程,从而让学生体会薛定谔方程主导的量子系统随时间的变化有什么特征.由于该系统是一个真实的实验系统,从另一个侧面让学生了解到已经有大量的实验从不同角度证实了薛定谔方程的正确性.讲解过程中对量子态的性质进行了详细的解释,让学生加深对量子态性质的理解.通过对该模型的量子态的性质讨论,引导学生考虑量子态的性质特征.

哈尔滨工业大学在 2018 年下半年开展了“新时代新工科大学物理教学体系”的改革,在原有的大学物理教学中增加了部分量子物理的比重.让学生掌握所学习的量子力学部分知识,在以后的工作和学习中都能能够学以致用,激发学生从事当前的最前沿内容的动力,从而认识到量子力学中的方程和假设是实验所证实的,进一步提高学生的科研能力.

1 增加内容特点与学生学习成效的分析

薛定谔方程时间演化教学内容特点:薛定谔方程是量子物理中最核心的内容之一,利用其描述物理系统和经典系统具有完全不同的性质,因此有必要建立物理模型,并进行方程的求解,进一步理解量子态的叠加性质,以及量子物理和经典物理之间的

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(12) 建立自主探究性实验教学模式提升实验教学的人才培养水平, 物理实验



建立自主探究性实验教学模式 提升实验教学的人才培养水平

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摘 要: 为了实现从“重视知识传授”向“重视能力培养”转变, 实现“以教为主”向“以学为主”转变, 在大学物理实验课堂教学中建立了自主探究性实验教学模式, 并研发出了与之配套的系列自主探究性实验项目。教学实践表明, 该模式培养学生实验能力和知识的应用能力方面取得了较好的效果。

关键词: 自主探究; 实验教学; IET

中图分类号: O11 文献标识码: B DOI: 10.13700/jyzz.2018.12.011

作为基础性实验课程, 大学物理实验在高校理工科各专业人才培养体系中占有重要的基础地位。但是传统的物理实验教学模式, 注重知识的学习和技能的培训, 在对学生能力的培养和创新意识的培养方面存在不足。针对该问题, 我们近年在教学模式上做了有针对性的改革, 从2014年起建立了自主探究性的实验教学新模式, 自主探究式的实验教学模式, 其主要着眼点是通过教学设计让学生深度参与实验探究的全过程, 学生通过自主设计, 自主搭建, 自主调试, 自主探索的方式, 建立实验系统并进行项目研究。在教师的指导下学生自主深度参与包括实验系统关键部件的设计和搭建的实验全过程, 对学生深入学习促进自身能力的提高产生了很大的推动力。实践表明, 这种实验教学模式激发了学生学习的主动性和勇于接受挑战的积极性, 改变了以往学生实验主要是利用现成仪器进行测量的实验学习模式, 在课堂教学中, 形成了科学探索和研究的氛围, 实验的过程使学生在知识、能力和素质三方面得到了协调发展。

1 自主探究性实验教学的要点和内容

为了实现从“重视知识传授”向“重视能力培养”转变, 实现“以教为主”向“以学为主”转变, 我们在大学物理实验课堂教学中建立了自主探究性 (Discovery, IET) 实验教学新模式, 并研发出

了与之配套的系列自主探究性 (IET) 实验新项目。

IET 自主探究性实验教学是基于问题导向教学法 (Problem based learning, PBL), 通过学生“自主设计, 自主搭建, 自主调试, 自主探索”的实验教学过程, 实现教与学的良性互动, 把科学研究的理念和方法贯穿在整个课堂教学过程中的一种教学模式。图1是 IET 探究性实验的教学过程框图。

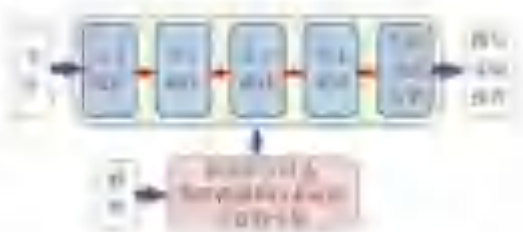


图1 IET 自主探究性实验的教学过程框图

在 IET 探究性实验项目选择上, 注重选择更具新颖性、技术性、挑战性的实验课题, 并注意和已经学过的物理实验内容相衔接。先后开发了 3 种 IET 探究性实验项目: IET 电磁混合耦合悬浮实验, IET 磁耦合共振式无线电力传输实验, IET 自旋光平台微镜、臂陀螺实验。

以“EIS 电磁混合耦合悬浮实验”的课堂教学为例, 指导教师首先组织学生研讨实现电磁混合耦合的物理机理和技术途径, 明确实验任务和

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(13) 基于光致瞬态光栅测量液体内的声速, 物理实验

基于光致瞬态光栅测量液体内的声速

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摘 要: 光栅通过调出与声波光栅, 其折射光的强度发生周期性变化, 变化频率即为介质传播声波的频率, 由此可以测量介质中的声速。利用光致瞬态光栅测量了常温、高压条件下磺基苯酚乙醇中的声速, 分别为 1.457 m/s 和 1.129 m/s 。与文献报道结果相符。光声干涉模式的声速测量方法, 便于与其他测量设备使用, 可用于探测高温、高压等极端条件下液体中的声速。

关键词: 光致瞬态光栅; 声速; 干涉模式测量; 液体

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声速是物理学、生物学、天文学以及海洋学等许多研究领域都需要的重要参量, 能够为探索介质内部微观结构、分子运动和分子间相互作用等提供重要信息, 因此对气态、液态及固态介质中声速的测量都具有重要意义。在大学物理实验中通常开设空气中声速测量^[1]和声光衍射方法测量液体中的声速^[2]等综合性实验, 这些实验普遍采用超声换能器在介质中产生固定频率的超声波, 再利用压电效应或光栅衍射原理对超声波的波长进行探测, 通常需要将介质与超声换能器直接接触, 属于接触式测量。本文介绍基于光致瞬态光栅测量介质中声速的方法, 该方法具有非接触测量的优点, 可用于观测高温、高压等极端条件下材料中的声速。

1 光致瞬态光栅测量介质中声速的原理

利用光致瞬态光栅技术测量介质中声速的原理如图 1 所示^[3]。来自同一光源波长为 λ 的两个反交秒激光脉冲以夹角 θ 入射到介质中, 当两个光脉冲实现时间、空间重合时, 在介质中产生干涉, 沿其形成的干涉条纹间距为 d 。如果介质在激光光的波长处有吸收, 则会跃迁至激发态, 形成激发态分子布居数光栅。介质分子通过无辐射弛豫从激发态跃迁回基态的过程中释放出的热能将

会使干涉条纹的亮纹区域被瞬时加热, 从而在介质中产生瞬态的热光栅。此后, 热扩散作用将会驱动高温区域(即干涉条纹的亮纹区域)的分子向低温区域(即干涉条纹的暗纹区域)移动, 导致介质中某一确定位置处的密度发生周期性振荡, 密度的周期性变化对应介质中双向传播的超声波, 超声波的波长即为干涉条纹间距 d , 频率即为介质密度振荡的频率 f , 与声速 v 的关系为

$$v = f \cdot d \quad (1)$$

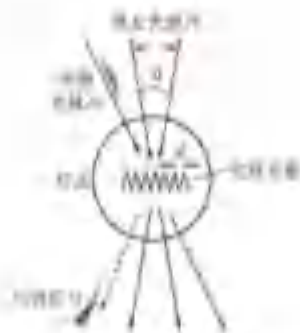


图 1 光致瞬态光栅技术测量声速的原理图

密度的周期性变化会在介质中形成周期性变化的折射率分布, 即瞬态布拉格光栅。让另一束光

“第 13 届全国高等学校物理实验教学研讨会”论文

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(14) 用外腔半导体激光获取强荧光物质的拉曼特征峰值. 物理实验

用外腔半导体激光获取强荧光物质的拉曼特征峰值

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【摘 要】利用外腔半导体激光作为激发光源, 采用移频激光法获得了强荧光物质三环唑的拉曼特征峰值. 实验结果发现, 在 230~400 nm 范围内激发条件下采集的三环唑样品拉曼光谱信号严重重叠, 在 310~330 nm 处采集到的拉曼光谱(300~330 nm) 和期望的拉曼特征峰, 与实际误差一致.

【关键词】拉曼光谱; 外腔半导体激光; 激光物理; 移频激光

【中图分类号】O434; O437.3; O438.3; O438.4; O438.5; O438.6; O438.7

光学类实验不仅是在普通物理实验教学中占有一定的比重, 而且也是近代物理实验课程的主要内容. 激光拉曼光谱是分子光谱学的重要分支, 是一种重要的物质结构和物质成分分析手段, 也是普通物理实验中吸收光谱实验内容的扩展, 因而有很多高校将其纳入了近代物理实验教学内容^[1]. 目前大多数高校所开设的拉曼光谱实验采用的样品都是四氯化硅, 它具有丰富的振动模式, 因而具有典型的拉曼特征峰. 为了拓展实验内容, 可以增加实验系统和实验参数的优化环节^[2], 或者引入新的测量样品^[3]. 在实际的生产生活应用中, 很多物质的拉曼特征峰并不十分明显, 甚至会被淹没在强烈的荧光背景中. 开展对强荧光物质拉曼特征峰的测量可以进一步丰富学生对于激光拉曼光谱的认识. 本文以水稻种植中广泛使用的农药三环唑为测量样品, 采用自制的可调谐外腔半导体激光器 (External cavity diode laser, ECDFL) 为激发光源, 应用移频激光法获取三环唑的拉曼特征峰值. 将可调谐外腔半导体激光器引入物理实验教学已见诸报道^[4], 可进一步丰富和拓展近代物理实验中激光特性和能量测量实验的有关内容.

1 实验原理

1.1 移频激光激发光源

在拉曼光谱检测中, 测量的拉曼信号往往是

叠加在大的荧光背景之上, 前者常常比后者大几个数量级, 给拉曼光谱分析带来很大困难. 对于分子的某一振动模式其拉曼位移是一定的, 因此拉曼散射光的频率会随入射激光的频率变化而等量变化. 但是荧光的发射频率却对入射激光频率的变化不敏感, 因此可利用拉曼散射光和荧光的这一特性差异, 通过改变激光光的波长来抑制荧光背景. 图 1 展示了移频激光法的主要过程和原理. 从图中可以看出, 若干个不同波长激光激发产生的拉曼光谱相减, 可以有效抑制荧光, 从而获得较为“纯粹”的拉曼散射光谱.



图 1 移频激光拉曼光谱原理示意图

1.2 外腔半导体激光器

一般来讲, ECDFL 是通过外部的光栅将部分激光能量反馈给 LD 芯片使得定频率的模式在竞争中胜出, 从而达到选频的目的. 光外腔的物理

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和大学物理实验的教学工作, 研究方向为激光与光学.



(15) 波尔共振实验中基于旋转矢量的相位差分析, 大学物理

波尔共振实验中基于旋转矢量的相位差分析

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摘要: 在波尔共振实验中, 通过引入两个旋转矢量分别描述受迫振动和驱动力, 解决了实验中因无法直接测量受迫振动位移而带来的困难. 针对有初速度的杆上质点与固定支点的连线形成角度和转角的情况, 由旋转矢量法分别给出了受迫振动与驱动力间的相位差. 研究表明, 无论初始角位移是否为零, 受迫振动与驱动力间的相位差恒为 $-\pi/2$ 与电机轴转动角位移恒相差 $\pi/2$ 之差.

关键词: 波尔共振; 旋转矢量; 相位差

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波尔共振是大学物理的一个基本实验, 通过这个实验可以很清晰地得到阻尼振动、受迫振动的振动规律, 特别是对于受迫振动可以验证其幅度和相位特性. 通过该实验可以使学生对理论课程阻尼振动、受迫振动内容有更深刻的理解. 波尔共振实验的一个重点是使用闪光灯测受迫振动与驱动力的相位差. 现有的文献 [1-4] 给出了闪光的瞬间旋转机械转盘上自转对应的角度即为相位差这一结论, 教学过程中学生因对这一结论不明其原理导致经常对此感到困惑. 本文将利用描述简谐振动的旋转矢量严格推导出受迫振动与驱动力的相位差.

1 受迫振动与驱动力的旋转矢量表示

波尔共振实验中, 电机匀速转动带动与电机轴偏心连接的连杆, 使其在平衡位置做简谐振动. 连杆又与波尔摆相连, 连杆的周期性运动相当于给波尔摆施加一个周期性的驱动力, 最终可以得到一个与驱动力同频的简谐振动. 为了研究受迫振动的相位特性需要测量受迫振动与驱动力这两个简谐振动的相位差. 如果分别知道任一时刻两个简谐振动的相位, 两者之差即可. 旋转矢量法是描述简谐振动的一个巧妙方法, 我们引入两个旋转矢量分别描述驱动力与受迫振动. 对于受迫振动令其振动方程为

$$x_1 = A \cos(\omega t + \varphi_1) \quad (1)$$

其中 A 表示振幅, ω 表示振动频率, φ_1 表示初相位. 对于驱动力, 令其满足的振动方程为

$$x_2 = B \cos(\omega t + \varphi_2) \quad (2)$$

其中 B 表示振幅, ω 表示振动频率, φ_2 表示初相位. 由波尔共振仪轴杆的设计可知, 波尔摆杆在平衡位置时连杆处在距离水平方向达到最大位移, 这意味着波尔摆杆在平衡位置时驱动力对应的简谐振动的位移为零. 不管电机从何位置启动, 将波尔摆首次通过平衡位置的时刻定义为初始时刻, 则可得在初始时刻波尔摆和驱动力的位移均为零, 即 $x_1 = x_2 = 0$. 由 $t = 0$ 时, $x_1 = x_2 = 0$, 可得 $\sin(\varphi_1) = \sin(\varphi_2) = 0$. 由 $t = 0$ 时, $x_1 = x_2 = 0$, 可得 $\sin(\varphi_1) = \sin(\varphi_2) = 0$. 考虑 φ_1 得 $\varphi_1 = 0$ 或 $\varphi_1 = \pi$. 考虑 φ_2 得 $\varphi_2 = 0$ 或 $\varphi_2 = \pi$. 经过时间 t , 得到与驱动力同频的简谐振动 $x_1 = A \cos(\omega t + \varphi_1)$ 与两个振动对应的旋转矢量都由初始位置 O 开始位置 P_1 分别如图 1、图 2 所示. 用旋转矢量 \vec{OP}_1 表示波尔摆的受迫振动, 用旋转矢量 \vec{OP}_2 表示驱动力.



图 1 受迫振动的旋转矢量表示

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(16) 工科院校物理演示实验体系的实践与探索. 实验研究与探索

工科院校物理演示实验体系的实践与探索

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摘 要:传统的大学物理课堂演示实验主要是通过讲解帮助学生理解物理知识的过程。然而,这种演示实验形式上过于简单直观,无法引起学生的深入思考和探究,导致学生不能主动接受,无法参与物理演示实验的设计、制作、演示等环节。本文结合自身特色并且围绕培养学生从知识学习能力的培养到科研创新能力提升和工程实践能力培养的目标,围绕重点物理实验教学改革并开展了相应的教学方法改革。多年的教学实践表明了这种演示实验体系能够较好地调动学生获取知识的主观性和锻炼动手能力,在培养学生综合素养和科研创新能力方面发挥着重要的作用。

关键词:演示实验; 大学物理; 教学方法改革; 工科院校

中图分类号: O4-35; G 420 **文献标志码:** A

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Practice and Explore on Physics Demonstration Experiment System of Engineering Colleges

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Abstract: In universities, traditional physics classroom demonstration experiments mainly help students to acquire knowledge of physics. However, these demonstration experiments are too simple and intuitive in form to make students the simply consideration and exploration of physical systems. The students passively accept knowledge, unable to participate in the design, production, presentation of these demonstration experiments. Considering the characteristics and students' conversion from the ability to learn knowledge to the research innovation capacity and engineering practice capacity, we launch a teaching reform and put it into practice around this new interactive demonstration experimental system. Years of teaching practice shows that such a demonstration experiment system based interactive teaching can greatly enable students to gain knowledge and exercise thinking ability. Thus, it plays an extremely important role in enhancing students' comprehensive quality and research innovation capacity.

Key words: demonstration experiment; university physics; teaching methods reform; engineering colleges

0 引言

物理学是实验的科学,这决定了物理演示实验在大学物理的教学中占有十分重要的地位。物理演示实

验是教师在传授物理知识的过程中为配合教学而将演示给学生看,在操作的同时引导学生观察、思考、思考和分析的一种物理实验教学方式^[1]。随着现代物理学的发展,物理现象不断地展示给学生,引导学生观察和思考物理现象^[2]。近40年,美国大学陆续开始推行大学物理演示实验^[3]。1980年,美国大学陆续开始推行大学物理演示实验^[4]。1980年,美国大学陆续开始推行大学物理演示实验^[5]。这些大学非常重视演示实验教学,力求通过多种形式包括实物演示、计算机模拟、动画电影等

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(17) 核物理多功能教学实验仪教学应用与功能开发. 物理实验



核物理多功能教学实验仪的 教学应用与功能开发

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摘 要: 根据教学需求自主设计了核物理多功能教学实验仪, 该仪器集高压、低压、放大、单道脉冲幅度分析器、IVC、VFC、定时、定标、控制与显示于一体, 充分利用主机公共资源, 通过切换不同种类的探测器, 选配不同种类、不同量程的放射源和不同几何配置的探测样品, 实现多功能的教学应用任务。

关键词: 核物理; 放射源; 探测器

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核物理实验是核物理专业的中心教学环节, 随着科学技术的进步和实验手段的提高, 近年来出现了一些新的高性能实验仪器设备^[1], 特别是计算机在核物理实验中的应用, 为核物理实验增添了无限的生机, 推动了核科学技术变革, 开创了核技术应用的新时代。核仪器的小型化、微机化、智能化, 都促进了原子核物理不断开拓新的研究和应用领域。核物理实验是理论与实践密切结合的基础课程, 体现了多种学科实验技术的综合应用。核物理实验教学可以培养学生理论联系实际, 善于发现问题、分析问题和解决实际问题的能力, 使学生初步掌握核物理的研究方法, 核物理实验的特点, 一般流程和规律, 培养严谨的科学态度、实事求是和坚韧不拔的工作作风。

基础性的核物理实验主要是强度测量, 其次是能量测量, 测量系统往往由高压电源、低压电源、线性脉冲放大器^[2]、单道脉冲幅度分析器、定标器等仪器组成。测量系统连接繁琐且所占空间很大。随着 NIM 系统的出现, 组合插件构成的测量系统可使仪器结构紧凑, 连接方便, 但 NIM 系统成本较高, 功能也较为单一。

为满足核物理专业的教学要求, 哈尔滨工业大学物理系自主设计了核物理多功能教学实验仪器, 如图 1 所示。自制的核物理多功能教学实验仪器是集高压、低压、放大、单道脉冲幅度分析器、



图 1 核物理多功能教学实验仪实物图

IVC、VFC、定时、定标、控制与显示于一体, 在充分利用主机公共资源的基础上, 通过前面板上的插插件切换不同量程的探测器, 选配不同种类、不同量程的放射源和探测测量及不同几何参数的探测样品, 可实现多样性配置, 从而实现多样性的功能。

核物理多功能教学实验仪的电气组成

图 2 为核物理多功能教学实验仪的电气组成原理框图。高压电源采用线性电源, 固定在机壳地板上, $\pm 15\text{ V}$ 为线性脉冲放大器、鉴别及分割器, $\pm 15\text{ V}$ 为 I/V 和 V/P 变换及高压电源供电, $\pm 15\text{ V}$ 为定时、定标、控制和显示供电, $\pm 24\text{ V}$

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(18) 等厚干涉实验中缺陷检测的数值模拟, 物理实验



等厚干涉实验中缺陷检测的数值模拟

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摘 要: 利用数值方法研究了等厚干涉实验中平面玻璃存在凸凹缺陷时的干涉条纹, 给出了不同缺陷尺寸、高度或曲率情况下的数值变化规律, 讨论了利用干涉条纹进行缺陷检测的方法。该研究给出等厚干涉实验中的典型问题的数值求解办法, 给理论和实验提供参考。

关键词: 等厚干涉; 数值仿真; 缺陷检测

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1 引言

等厚干涉实验是物理实验课程中的经典实验项目, 其中牛顿环和劈尖干涉实验最为典型, 在大学物理实验课堂上也是最为常见^[1]。学生可以通过实验来观察干涉现象, 同时也学习了利用等厚干涉进行曲面曲率和薄膜厚度测量的方法。实际上, 等厚干涉实验除了用来做曲率和厚度测量外^[2], 另外一个典型的应用就是玻璃表面缺陷的检测^[3]。然而, 带有特定尺寸和形状的缺陷的平面玻璃不易加工, 这也影响了这类实验的开设。目前, 多数物理实验教材只把这个问题作为讨论题供学生思考, 然而, 仅靠理论分析, 无法验证分析的准确性, 也不能得到直观的认识。

事实上, 当牛顿环或劈尖干涉装置中的玻璃表面存在的凸起或凹陷形状比较规则, 且变化比较平缓时, 可以利用计算机程序来模拟其产生的干涉现象, 并给出干涉条纹的形状和分布, 进而研究干涉条纹随结构参数变化的规律。在本次研究中, 选择劈尖干涉装置, 研究玻璃存在球面凸起或凹陷时, 不同缺陷尺寸、高度和曲率对干涉条纹的影响。

2 原理与公式

图 1 为劈尖装置示意图, 1 块平板玻璃水平

放置, 另一块玻璃与水平玻璃存在一定夹角, 下侧玻璃上表面存在 1 个球面凸起缺陷, 2 块玻璃之间形成一定厚度的空气间隙。当一单色平面波从上方垂直照射到劈尖装置上时, 其中一部分光在 A11 表面反射并与从 A22 面反射的光发生干涉。假设干涉的 2 束光的强度均为 I_0 , 则距离 O 点水平距离为 x 处的干涉条纹的强度可以写为^[4]:

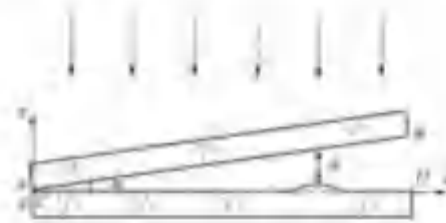


图 1 存在缺陷的劈尖结构图

$$I(x) = 4I_0 \sin^2(\pi \Delta d / \lambda) \quad (1)$$

其中 λ 为入射光在真空中的波长, Δd 为 x 位置处 2 个玻璃表面间的距离, 如图 1 所示。在无缺陷的位置, 空气间隙的厚度为线性变化, 即:

$$\Delta d(x) = x \tan \alpha \quad (2)$$

其中 α 为劈尖的夹角, 在缺陷位置, 空气间隙的厚度不再是线性变化, 考虑到缺陷的形状为球形

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(19) 物理学拓展课程提升学生综合素质，物理与工程

物理学拓展课程提升学生综合素质

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摘 要 在大学物理课程教学中,受到学时及自身体系的限制,教师很难将物理学的前沿问题、最新发展以及对其他学科的影响更多地传授给学生,这在一定程度上影响了学生综合素质的提升.本文通过“改变世界的物理学”课程的教学实践介绍物理学拓展课程在创新人才培养以及提升学生综合素质方面的作用.

关键词 物理学拓展;综合素质;物理学思维方法

PHYSICS CURRICULUM DEVELOPMENT PROMOTES STUDENT'S COMPREHENSIVE QUALITY

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Abstract Restricted by the teaching hours and systems, in the teaching of university physics courses, teachers can hardly instruct frontier physics problems, its recent development and its influence on other sciences to students, which affects the improvement of students' comprehensive quality to some extent. This paper introduces the role of physics curriculum development in the cultivation of innovative talents and improving students' comprehensive quality through the teaching practice of the course "physics changes the world".

Key words physics development; comprehensive quality; thinking methods of physics

学生综合素质的提升有赖于知识面是否广泛,思维方法是否正确.物理学拓展课程的综合素质提升不可替代的作用.这是因为在物理学发展过程中的每一次重大突破不仅带来物理学新领域、新方向以及交叉学科和新技术学科的崛起,而且改变了人们的思维方式.“物理学为科学进步和发明的利用,提供所需训练有素的人才,物理学还培养化学家、工程师、计算机科学家,以及其他物理科学和生物医学科学工作者的教育中,是一个重要的组成成分”.物理学的思维方法在现代信息化战争中也发挥着重要作用.

但受到大学物理课程学时及自身体系的限制,教师很难将物理学的前沿问题、物理学的最新发展以及对其他学科的影响更多地传授给学生,这在一定程度上影响了学生对物理学重要性的认识,也影响了学生学习与物理的兴趣.所以有必要开设物理学拓展课程,将大学物理课程中涉猎不到

的前沿问题、与其他专业融合、渗透以及物理学思维方法对高新技术发展的推动作用介绍给学生.

我校于2014年秋季学期尝试以“改变世界的物理学”为题为本科生开设了全校任选的物理学拓展课程.自开课以来,一直受到学生们的喜爱,平均每年的选课人数在300人左右,选课人数位于全校选修课程前列.“改变世界的物理学”课程荣获了2019年哈尔滨工业大学优秀课程称号.本文基于“改变世界物理学”课程的教学理念和教学实践介绍物理学拓展课程在创新人才培养以及在提升学生综合素质中的作用.

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(21) 关于“光的偏振”教学的一些尝试. 物理与工程

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关于“光的偏振”教学的一些尝试

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摘 要 在大学物理教学中, 让学生更好地接受知识, 是每一位教师的心愿。物理课程是理工科学科的重点, 课堂教学中应多开展演示和实际应用实验, 按照这一想法, 笔者在2018年秋季学期对“光的偏振”部分做了新的教学尝试。课堂开始以演示实验开始, 结合实验现象来讲解物理知识, 分析实验条件和计划对现象的影响, 并通过不同现象来加深对物理知识的理解, 实践表明, 教学效果良好且可推广, 大部分学生对这种教学模式持肯定态度。本文介绍新教学模式的指导思想, 实施计划和教学流程。

关键词 大学物理教学; 光的偏振; 演示实验

AN ATTEMPT ON TEACHING “POLARIZATION OF LIGHT”

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Abstract In college physics teaching, students' better acquisition of knowledge has been the responsibility of every teacher. According to the characteristic of experimental source, college physics teaching should focus more on demonstration experiments and practical application. Based on the idea, a new teaching method on “polarization of light” was tried in fall semester of 2018. The class begins with demonstration experiments, and relative knowledge is conveyed accompanying experimental phenomena. Then the influence of experimental conditions on phenomena is discussed, and applications are given for better understanding of physical knowledge. Practice shows that the teaching effect is pretty good, and the majority of students have positive attitude towards this teaching method. This paper introduces the guiding ideology, implementation plan and student feedback of the new teaching method.

Key words college physics teaching; polarization of light; demonstration experiment

教育的目标是“授之以渔”, 而非“授之以鱼”, 也就是说, 育人是教育的最终目的。在教育过程中, 学生是学习的主体, 教育目标的达成, 离不开学生的主动学习。因此, 调动学生的学习积极性, 让学生能够主动地学习, 在课堂教学中至关重要。

在高等教育中更是如此。

我教的大学物理课程是在本科一年级的春季学期和二年级的秋季学期讲授的。经过高考前的紧张复习和磨炼, 加上高考后漫长的暑假, 一部分学生在学习上比较懒散, 甚至出现厌学情绪, 在课

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(22) 关于“光的干涉”部分教学模式的探讨. 物理通报

关于“光的干涉”部分教学模式的探讨^{*}

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摘 要:对大学物理“光的干涉”部分教学模式进行探讨,通过物理过程引领数学的展开,理论与实验相结合,建立“理论数学引领型”,以此作为贯穿“光的干涉”部分的主线,引入物理过程相对应的工程技术应用,提高学生的学习兴趣,让课堂变得生动有趣,提高课堂上学生的抬头率。

关键词:大学物理;光的干涉;教学模式

1 引言

高等教育应以人才培养为根本,大力提升本科生教学质量,以改革创新为引领,不断增强科技核心竞争力,大学物理作为理工科大学的一门公共基础课,在人才培养方面发挥着重要作用,但是,随着社会的发展,大学物理教学面临的问题越来越多,互联网的发展导致学生学习知识多元化,但是掌握知识的方法简单化,碎片化问题也越发严重,大学物理传统的教学模式是老师讲,学生被动地听,学生兴趣缺失,课堂参与度不够,最终导致的结果是学生学大学物理,只是为了应付考试,能力并未得到真正的提高^[1]。要想改变这些问题,大学物理教师应在课堂上吸引住学生的注意力,提高学生的兴趣,如何做这一点,就需要正确处理“教”与“学”的问题;需要大学物理教师在教育理念、教学方法、教学内容、教学模式等多方面下功夫^[2]。本文主要针对“光的干涉”部分的教学模式展开探讨。学生在高中时对光学部分的知识掌握主要在于几何光学,如光的直线传播原理、折射定律、反射定律等,而对于杨氏双缝干涉、薄膜干涉等干涉等“光的干涉”内容还比较陌生,光的相干叠加,理论性强,不利于学生理解和掌握。为了让学生真正掌握“光的干涉”部分的知识,本文将从以下教学过程安排探讨新的教学模式。

2 教学过程

本文提出的新的教学模式需要引入物理过程,每节课对应一个物理过程,具体教学过程如图1所示。



图1 基于物理过程的教学过程

(1)首先介绍本节课的知识点和能力培养,大学物理教学应始终以学生为中心^[3],每一节课都应该让学生在开始本节课之前就了解本节课的知识点,以及本节课主要培养自己哪方面的能力,例如,杨氏双缝干涉的主要知识点为杨氏双缝干涉的光程差、条纹位置和条纹间距等,主要培养的能力有近似处理问题的能力以及多角度思考问题的能力。

^{*} 黑龙江省教育厅“十三五”规划课题“新时期工科大学物理课程体系建设与实施”(项目编号:1310101000)和哈尔滨工业大学教育基金会项目“物理教学改革创新”项目(项目编号:1310101000)资助。

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(24) “工作室物理”教学中的演示实验. 物理与工程

“工作室物理”中的演示实验

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摘要 “工作室物理”将学生的探索研究融入到课堂教学中, 通过演示实验和教学软件激发学生主动学习, 也叫作 TEAL (Technology Enabled Active Learning); TEAL 充分体现了“以学生为中心”的教学理念, 充分发挥学生自主学习的积极性, 而课堂演示实验直观、真实、可信度高, 学生可直接参与其中, 起到了教与学的互动, 由学生自己做演示实验, 分析实验现象, 总结实验规律, 使学生从被动获取知识到主动获取知识, 从而锻炼了思维能力, 是培养学生创新思维、创造能力以及提升学生实践能力的最好的手段和途径。

关键词 工作室物理; 演示实验; 互动式教学

在传统“灌输式”教学模式中, 学生是被动的知识者、接受者, 他们的注意力主要是放在对知识要点的掌握上。这在一定程度上剥夺了学生自主学习、主动获取知识的行为, 限制了学生的创新与创造性思维的发展, 同时也违背了“以学生为中心”的教学理念。

“工作室物理” (Studio Physics) 教学模式^[1]提倡导师和引导, 学生主动获取知识, 主动探索物理规律, 学生之间、师生之间互动, 互动式的教学模式, 是充分发挥学生自主学习, 实现“以学生为中心”的教学理念的最佳教学模式。这种教学模式需要在专任教师中进行, 主要是营造一个以学生为中心的主动学习环境 (Student-Centered Active Learning Environment for Undergraduate Programs, 缩写: SCALE-UP)。这种教学模式通过演示实验和教学软件激发学生主动学习, 所以通常也叫作 TEAL (Technology Enabled Active Learning)。这种教学模式培养学生从三维角度思考问题。由于学生有通常接触的多为平面的屏幕和平面的书刊, 因此思维方式被限定在二维的范围内, 而对于物理现象的理解常需要从三维的角度思考问题, 尤其在力学、电磁场和光学领域。

目前, 世界上已有近百所高校针对其不同需求, 采用了“工作室物理”教学模式。如美国的麻省理工学院 (MIT)、哈佛大学、耶鲁大学、台湾的中正大学, 等等^[2-4]。并且, 美国一些大学已经开始在化学、工程类各门课程中实施这种教学模式。我校 (哈尔滨工业大学) 从 2014 年开始试行这种教学模式^[5], 取得了一定的成果, 今年春季学期学校建成了 4 个“工作室物理”教室, 营造了 SCALE-UP 学习环境, 如图 1 所示, 从最初的一个试点增加到 4 个所。但由于条件所限, 我们只是基于课堂演示实验引导学生自己动手做演示实验, 仔细观察、总结实验规律, 从而分析、总结出物理规律, 激发学生自主学习的兴趣; 同时也培养了学生的科研能力、创造能力, 提升了学生的科学素养。

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(25) 刚体定点转动代数讲法的教学实践. 物理与工程

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教学经验交流

刚体定点转动代数讲法的教学实践

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摘 要 代数是人们描述自然规律的一个极为重要的数学工具,但它自身的抽象性又使得初学者很难掌握这一数学工具.理论力学是物理系本科生的第一门理论课程,将运用数学工具处理物理问题的训练明显较普通物理加强.我们尝试在理论力学课的刚体力学一章采用代数讲法,从而使抽象的数学理论与具体的物理刚体有了直接的对应,并为后续专业类课程学习打下了良好的基础.

关键词 理论力学;刚体定点转动;线性代数

USING THE ALGEBRAIC METHOD TO TEACH THE FIXED POINT ROTATION OF RIGID BODY

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Abstract Algebra is a very important mathematical tool in physics research and description. But it is also very difficult for beginners to learn due to its abstraction. Theoretical mechanics is the first theoretical course of physics undergraduates, and its training of using mathematical tools to deal with physical problems is more important than that of general physics. We try to use algebraic method to teach the fixed point rotation of rigid body in theoretical mechanics course, to make a connection between the abstract mathematical theory and specific physical issues, and lay a solid foundation for the follow-up courses.

Key words theoretical mechanics; the fixed point rotation of rigid body; algebraic method

代数无论在物理学研究还是在应用领域都是一个极为重要的数学工具.初学者由于缺乏实际应用方面的训练很难掌握这一数学工具.这很大程度上是由于线性代数通常在大一上学期学习,而大量运用代数知识的专业类课程一般都安排在大三下学期或者大四上学期.基于这样的情况,我们尝试在大二上学期理论力学课的刚体一章引入代数讲法.这样做的原因在于,刚体的定点转动可以用固定在刚体上的刚体坐标系(刚体系)

来描述.这样对刚体的描述就转化为对两个坐标系间相互关系的描述,从而可以用代数理论严格表示出来.此外,刚体转动是同学们在日常生活中常见的现象,通过将转动与代数理论的结合可以帮助同学们更好地复习和理解大一刚刚学过的线性代数的很多抽象公式.经过我们多年的教学尝试,这一教学思路已经取得了较好的教学效果.现在给出这一讲法的主要步骤:

令 O, O' 是转动前的刚体系, O', O'' 是转

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工作室物理教学中的互动反馈教学设计*

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摘 要: 将互动反馈教学系统应用到工作室物理教学中, 进行了基于互动的互动反馈教学设计. 工作室物理教学以学生主动获取知识, 探索物理规律, 互动反馈系统提供多种数据反馈以创造学习环境为中心开展课堂. 提高了学生主动获取知识, 探索物理规律, 互动反馈系统提供多种数据反馈以创造学习环境为中心开展课堂. 提高了学生主动获取知识, 探索物理规律, 互动反馈系统提供多种数据反馈以创造学习环境为中心开展课堂.

关键词: 工作室物理教学; 互动反馈; 大学物理

工作室物理^① Studio Physics^② 教学模式是由美国约翰斯·霍普金斯大学创立的一种集成了理论和实验的教学方式. 这种新的教学模式的教学理念是引导学生主动获取知识, 探索物理规律, 其特点优势是实验和理论课程有机结合. 目前, 这种工作室物理教学模式已在美国、中国台湾地区以及一些欧洲国家的大学进行了将近 2.6 年的时间^[1-3], 并取得了良好的教学效果. 哈尔滨工业大学从 2014 年也开始在大学本科课程中进行工作室物理试点班 (32 人) 教学尝试, 深受学生的欢迎, 取得良好的教学效果^[4]. 在实践的探索中, 我们发现在进行工作室物理教学过程中, 教师对学生课堂活动的引导是提高这种教学模式教学效率的关键. 我们下一步将在正常的教学班级 (一般人数在 150 人左右) 中进行这种教学模式, 希望有更多的学生受益. 在考虑如何将工作室物理教学模式应用于正常教学班级 (一般为 150 人左右的教学班) 时, 我们面临的首要问题是: 如何提高教师和学生课堂参与度, 在学生分组进行实验探索时引导学生, 积极引导学生的学习活动. 经过多方面的调研和尝试, 我们认为将互动反馈教学系统运用到工作室物理教学课堂中, 是一种非常行之有效的办法.

从教学角度, 促进学生发展是学校永恒的主题, 为了保证教学质量, 一个核心的问题就是了解学生对知识的掌握程度, 合理安排教学内容并调整教学进度, 在教学中加强和促进教学过程中师生之间

的沟通、交流与互动是很必要的. 然而, 在现实的教学环境下, 由于物理课程上课学时等客观因素的限制, 要想在教学过程中充分体现师生之间的交流与互动并非易事. 而“互动反馈教学系统”, 正好弥补了这个问题. 互动反馈教学系统是一种提高课堂互动, 鼓励学生参与到教学活动中的管理工具. 它是基于“互动式学习理论”设计的一种允许学生在课堂上通过手持设备以选择或数字来回答问题的学生互动反馈系统^[5]. 教师上课时根据教学的内容提出标准化的问题 (例如选择题或者数字化的问题), 学生利用手持设备选择答案, 系统后台具备数据记录、统计、分析等管理功能, 可以根据学生的回答答案实时形成统计报表显示在屏幕上, 教师可以随时查看, 并根据报表掌握学生的掌握程度, 即时调整进度. 除了口头出题, 系统还支持应答测试、投票游戏、测试比赛、拖拽发言、问卷调查、点名考勤等多种形式的互动方式, 活跃了课堂, 激发了学生参与课堂学习的积极性, 高效地促进了教师的课堂教学活动^[6]. 在大班工作室教学模式中, 借助于互动反馈教学系统, 可以实现多种形式的师生互动, 能够有效控制课堂, 促进师生互动, 解决了课堂反馈教学实践过程中存在的问题.

1 实践成效以及发现的问题

从 2014 年, 我校开始每学期都从英才学院挑选一个班级进行工作室物理教学尝试, 已经进行了 3

* 黑龙江省教育厅教育科学规划项目 (项目编号: 125214001176) 资助项目. 黑龙江省教育科学规划项目 (项目编号: 142044) 黑龙江省高等学校“十二五”规划课题. 项目编号: 140016. 作者简介: 张伶俐 (1990—), 女, 硕士, 助教. 从事大学物理教学-实验室教学研究.

(27) 利用通识课程《物理学史》来传递物理文化和培养物理思维—— 通识课程《物理学史》建设及改革初探， 物理通报



利用通识课程《物理学史》来传递 物理文化和培养物理思维

—— 通识课程《物理学史》建设及改革初探 ——

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摘 要: 通过问卷调查深入了解了本科《物理学史》课程的需求,分析了将《物理学史》课程设为本科通识课程的可行性,帮助学生形成正确的概念,对物理学史课程的教学理念、教学内容和考核手段等方面做了改革,实践证明,物理学史课程的教学可以实现传递物理文化,培养物理思维模式的教学目标。

关键词: 通识课程 《物理学史》 教学理念 教学手段

1 引言

物理学史是人类对自然界各种物理现象的认识史,主要研究物理学发生和发展的基本规律,包括物理学概念、思想的发展和变革,物理学发展与社会背景关系,各个分支学科间联系以及与其他学科的联系。近年来,人们对物理学史的教育功能及其在大学物理、大学物理实验教学中的作用进行了深入的探讨和教学实践研究,很多高校的物理专业,都将物理学史设为通识性的基础必修课或选修课程,可以说物理学史的教育地位已经被广大物理教师和教育学家认可,但是,作为通识课程的物理学史应该具有怎样的教学理念?究竟应该讲什么怎么讲?这些内容是否对学生有吸引力?如何能让这门既有科学性又有文化性的课程担负起其重要教育的责任?这些都是需要我们去深入探讨的。

2 问卷调查

为了解决这些问题,我们面向哈尔滨工业大学2011年入学的本科生发放了调查问卷,共发放了100份问卷,收回有效问卷81份。问卷内容涉及本课程的选修意愿,课程内容和学习方式等几个方面,目的是为了了解学生对物理学史的需求情况。问卷涉

及的问题以及调查结果如表1所示。

表1 调查问卷

| 问题 | 选项 | 所占比例/% |
|---------------------------|--------------------|--------|
| 1. 是否愿意选修《物理学史》通识课程 | A. 愿意 | 75.2 |
| | B. 不愿意 | 24.8 |
| | C. 无所谓 | 18.4 |
| 2. 能否在课余时间完成《物理学史》课程 | A. 能 | 12.3 |
| | B. 不能但不愿意自学 | 25.9 |
| | C. 可以自学一部分 | 36.2 |
| | D. 愿意通过 Web 进行学习 | 25.6 |
| 3. 你希望能通过《物理学史》了解哪些内容(多选) | E. 要自学,但需要老师给出明确指导 | 25.6 |
| | A. 物理学史史学的前沿进展 | 43.2 |
| | B. 历史对物理的学习兴趣 | 43.2 |
| | C. 物理学家的思维习惯、物理思想 | 43.2 |
| | D. 了解物理学的前沿 | 43.2 |
| | E. 物理学家的研究方法 | 43.2 |

调查结果表明,多数同学对《物理学史》这门通识课程有兴趣,会选择选修,还有一部分同学没有明确的反对,这就是说,针对全校开设这样一门通识选修课程具有可行性。在对该课程教学手段的调查中,结果表明,认为学习过程不需要老师参与的同学(选择A的)约为16.8%,而37.6%的学生选择

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(28) 分离变量法教学内容优化及本征值问题引入方式研究. 物理通报

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大学物理教学

分离变量法教学内容优化及本征值问题引入方式研究

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(收稿日期:2017-05-04)

摘要:现有数学物理方法教学体系中,分离变量法在前,本征值问题在后.而在分离变量过程中,又涉及到本征值问题,这样的安排导致学生在学习分离变量法过程中,不能很好地理解本征值问题是分离变量法的基础,不利于学生严密逻辑思维能力的培养.针对这个问题我们开展了分离变量法教学内容优化的研究,提出一种更有利于学生严密逻辑思维能力的培养的分离变量法教学方案:将本征值问题提前,将其置于定解问题之前,分离变量法之前.进而,为避免直接引入 Sturm-Liouville 方程而导致的突兀性问题,给出了分离变量法教学顺序调整后的 Sturm-Liouville 方程的引出方案.

关键词:分离变量法 本征值问题 教学内容优化

数学物理方法是物理系的主干课程之一,是一门利用数学语言描述并解决自然科学及工程技术中所遇到的一些问题,建立定解问题,求解定解问题的专业基础课.课程的教学目的是通过课程的学习进

一步提高学生的数学素养和应用数学知识解决复杂物理问题的能力;为理论物理课程的学习打下牢固的基础.但现有的教材内容体系存在一些问题,不利于学生严密思维能力的培养.我们知道,求解定解问

Rethinking on the Adiabatic Process of Ideal Gas from Microcosmic Perspective

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Abstract: The paper "Deducing adiabatic equation of ideal gas via kinetic theory of gas" (College Physics (Major: 2017) mentioned a non-rigorous way to calculate the ideal gas adiabatic formula: the physical nature of this derivation is the collective effect of the gas molecules on the piston wall, that is, the microscopic interpretation of the gas pressure. Based on this point, this paper gives the corresponding derivation from the perspective of the external work of gas molecules, and finds a more intuitive solution in its given model and correspond more directly to the first law of thermodynamics. Thus more directly links itself to the micro-process and microscopic ideal gas state equation.

Key words: ideal gas; adiabatic process; microscopic work

(29) 压缩传感在物理演示实验中的应用. 物理与工程

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物理实验



压缩传感在物理演示实验中的应用

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摘要 将科学前沿技术融入物理演示实验中有助于大学生科研兴趣和创新能力提高, 解决近几年兴起的科研热点问题——压缩传感技术融入物理演示实验中, 研制了一套值得推广的单像素成像演示设备, 该设备符合物理学实验应具有的基本操作性和便捷性, 还具有新颖原理和深远深刻的特点, 可满足不同层次涉及物理学和光学、统计物理、热力学等相近领域, 同时还涵盖了信息技术、通信技术、计算机技术等各个交叉领域的知识, 本实验的开发对通过物理演示实验来提高大学生创新能力科研兴趣方面具有参考价值。

关键词 演示实验; 科技前沿; 压缩传感; 单像素成像

AN APPLICATION OF COMPRESSIVE SENSING TECHNOLOGY IN PHYSICS DEMONSTRATION EXPERIMENT

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Abstract Merging the cutting-edge science and technology into physical demonstration experiment will help students improve research interests and creative ability. By integrating the compressive sensing technology, which is a hot topic rising in recent years, into demonstration experiments, we developed a set of single pixel imaging demonstration equipment which is worthy to be popularized. The device meets the requirement of physical demonstration experiment for its ease and portability. It also has some characteristics of evidencing visible phenomenon and revealing profound principle. The contents of our device are not only related to physics knowledge including optics, statistical physics and optoelectronic devices, but also cover the basics of a plurality of cross-cutting areas of information technology, communications technology and computer science. This article will be very valuable for enhancing students' creativity and research interests through physics demonstration experiments.

Key words demonstration experiment; technological frontier; compressive sensing; single pixel imaging

物理学是实验的科学,这决定了物理演示实验在大学物理的教学中占有十分重要的地位,然而,现有的演示实验目前多数侧重于通过简单直观的实验来辅助学生理解大学物理的课堂教学内容,而不能引起学生对科学问题的深入思考和有效地激发学生的科研兴趣^[1]。演示实验内容严重与现代科技脱钩,忽视了尖端科学技术在提高大学生科研兴趣和创新能力方面起到的重要作

用。本文以理工科大学的物理演示实验设计为背景,通过近几年兴起的科研热点之一——压缩传感融入到演示实验中,研制了一套值得推广的设备

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(30) 学科思维导图与物理知识网络的构建

学科思维导图与物理知识网络的构建

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摘 要 大学物理是理工科学生的一门基础课程, 在高校教学中的重要性和难度显著。学习大学物理, 不仅要掌握自然界的客观规律、分析和研究方法, 更重要的是从整体上认识和把握物理学并应用于实际。物理学中诸多物理量有着内在的必然的联系, 通过绘制学科思维导图或概念图, 来探究这些相关联的物理量的关系和本质, 物理规律的应用范围, 可以更透彻地理解和掌握物理知识。本文将思维导图工具引入大学物理的教学中, 利用思维导图帮助学生对物理量间的关联, 构建知识网络, 激发学生的求知热情, 提高学习效率和教学效果。

关键词 大学物理; 思维导图; 学科思维导图; 概念图

0 引言

作为基础学科的大学物理课程, 是高校为理工科学生提供的全面系统地掌握物理学基础的重要课程。大学物理涉及力学(质点运动学、质点动力学、刚体、转动、波动、相对论基础)、热学(气体动理论、热力学)、电磁学(静电、稳恒磁场、电磁感应)、光学(干涉、衍射、偏振)、量子力学(波粒二象性、薛定谔方程等)等五个主要部分, 包含了物质的基本结构、物质运动、物质的相互作用与转化的普遍规律, 其基本概念、基本定理、定律被广泛信用于所有的学科领域。因为涵盖内容广, 概念、物理量、公式、定理定律繁多, 学生在大学物理课程的学习中, 不仅要掌握自然界的客观规律, 研究方法和解题技能, 还要从整体上认识和掌握物理学, 因此有相当大的难度, 也让许多学生感到畏俱和排斥。

思维导图, 是一种发散性的、图像式思维的工具, 是美国著名心理学家、教育学家 Tony Buzan 发明的一种有效的思维模式^[1]。思维导图运用图文并重的技巧, 把各级主题的关系用相互隶属与相关的层级图表现出来, 把主题关键词与图像、颜色等建立记忆链接, 充分运用左右脑的机能, 在科学、逻辑、判断与艺术、图像、感觉等之间平衡发展。思维导图既可以优化人们的思维和记忆, 提高工作效率, 又能激发人的创造力。

近年来, 思维导图这一功能强大的思维工具, 已被越来越多的人掌握和使用, 广泛地应用于商业、学习以及日常生活中。特别是随着众多的思维导图制作软件的推出和不断更新, 如 Mindmanager、XMIND、mindmap、FreeMind、NovaMind、Mind Vector 等, 更多的人喜欢上这一工具。哈佛大学、剑桥大学、伦敦经济学院等知名学府正在使用和教授思维导图, 华中科技大学 2013 年在国内率先推出了思维导图公开课, 得到学生热烈欢迎, 把学科思维导图的图、文、色、关联引入大学物理教学, 特别是后学习阶段, 对学生整合物理知识, 提高学习兴趣和学习效率大有帮助, 已有老师在进行这方

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(32) 静电场中的相关能量

静电场中的相关能量

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摘 要 从点电荷到带电体, 对静电场中引入的几个相关能量, 即电势能、电场能、电相互作用能、自能、固有能, 互能等做了辨析和探究。电势能是带电体间或带电体与电场间的相互作用能, 具有相对性和共有性。在指定电荷相距无限远电势能为零的条件下, 解析法给出的电场能本质上就是电势能, 自能即固有能和互能分别是对单一带电体的电场能和多带电体间的相互作用能的另一种表述。论文还进一步对点电荷电场能发散问题做了讨论, 指出了球形带电体场能公式的适用极限。

关键词 电场能; 电势能; 自能; 互能; 点电荷自能

静电场中, 根据电场力对电荷做功与路径无关, 引入了电势能的概念。电势能是电荷与场的电荷或其所激发的电场之间的相互作用能, 因此为系统所共有, 但其具有相对性的属性。另一方面, 在静电场中移动电荷, 电场力能够做功又说明静电场蕴藏一定的能量, 这就是静电能。电容器放电有发光、发热等现象产生, 就是电容器储存的静电能转换为其他形式能量的结果。对于静电场中这两个能量-电势能和静电能, 一般教材中少有对其关系做进一步的说明, 容易让学生产生疑惑, 那么电场能与电势能是否是两种不同性质的能量, 有什么区别和联系呢? 关于这方面的教学研究也有发表^[1,2], 本论文从静电学中电势能和电场能两个概念的引入出发, 探究它们的实质和内在关联, 指出电场能的本质就是电势能, 文中也对自能和互能做了讨论, 并从点电荷的电场能积分公式发散, 说明点电荷作为一个理想化的模型, 在非理想情况下其自能公式的适用极限。

1 点电荷系的电势能

最简单的带电体就是点电荷。当两个点电荷 q_1 和 q_2 从相距无限远移至相距 r_{12} 时, 电场力做功 $A = -\frac{q_1 q_2}{4\pi\epsilon_0 r_{12}}$, 设两点电荷相距无限远时系统电势能为零, 由保守力做功等于势能的减少量, 则两电荷相距 r_{12} 时系统的电势能就是

$$W_{12} = \frac{q_1 q_2}{4\pi\epsilon_0 r_{12}} \quad (1)$$

这一势能也是移动电荷过程中外界克服电荷间相互作用的库仑力所做的功, 因此也称为电荷间的相互作用能。式 (1) 写成另一种形式为

$$W_{12} = \frac{1}{2}(U_1 q_1 + U_2 q_2) \quad (2)$$

其中, U_1 是点电荷 q_1 产生的电场在 q_2 所在位置的电势; U_2 是点电荷 q_2 产生的场在 q_1 所在位置的电势, 后面的讨论如无特别说明均设定无限远点为电势零点, 如果继续在系统内移入第三个点电荷 q_3 ,

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(33) 基于教学的洛伦兹变换推导的探讨

基于教学的洛伦兹变换推导的探讨

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摘 要 分析了国内外大学物理教材中的几种洛伦兹变换的推导方法, 从教学出发, 对相对论洛伦兹变换推导处理加以探讨。

关键词 相对论教学; 洛伦兹变换; 推导

狭义相对论是大学物理教学中的难点, 教师在讲授过程中, 能够提供的实例有限, 而对于经典时空观已经根深蒂固的学生, 在学习的过程中更是有许多难以理解, 感觉不可思议的问题, 让学生欣然接受相对论理论并学懂学透, 对教师也是一个挑战。

众所周知, 爱因斯坦的两条假设是狭义相对论的根本, 由此得到的洛伦兹变换又是新时空理论, 描述物理量的变换以及所有物理规律的协变性的基础。关于洛伦兹变换, 在国内外使用的各类教材^[1-5]及一些文章^[6-8]中, 有多种推导方法, 但部分推导或者因为后果没有交代清楚, 或者推导过程中关系式的给出令人信息, 或者又引入许多新假设, 这样的推导不但教师不自信, 学生也会有抵触心理。因此让学生一开始就能接受洛伦兹变换, 无疑在后面的教学中会收到事半功倍的效果。

1 基于物理教学的几种推导方法

在下文的推导中涉及的两个惯性参考系 S (时空坐标用 (x, y, z, t) 表示) 和 S' (时空坐标用 (x', y', z', t') 表示) 均按常规的约定, 即 x 轴和 x' 轴重合, y, z 轴和 y', z' 轴相互平行, 并且在 $t=t'=0$ 时, 原点 O 和 O' 重合, S' 相对 S 以速度 u 沿 x/x' 轴作匀速直线运动。在这样的约定下, 绝大多数的推导都直接给出洛伦兹变换中的两个关系式 $y=y', z=z'$, 其余两个关系式的推导大致可总结为如下 3 种。

推导 1: 考虑到运动的相对性, S' 相对 S 以速度 u 沿 x/x' 轴作匀速直线运动, 则有 S 相对 S' 以速度 $-u$ 沿 x' 轴作匀速直线运动, 文献[1]的将其称为爱因斯坦条件。

测量 O 点的坐标, 在 S 系总有 $y=0$, 在 S' 系则有 $y'=-ut'$, 即 $y'+ut'=0$, 对此特殊点二者同时为零, 可认为在任意时刻, 对任意点, 二者有比例关系

$$x=k(x'+ut') \quad (1)$$

同理测量 O' 点的坐标时, 可得

$$x'=k'(x-ut) \quad (2)$$

根据狭义相对论的相对性原理, 两惯性系等价, 有 $k=k'$, 假设在 $t=0$ 时刻在原点发出一个光信号, 利用光速不变原理, $x=ct, x'=ct'$, 结合式 (1) 和式 (2) 可确定出系数 $k=(1-\beta^2)^{-1/2}$, 不再赘述。

推导 2: 首先说明要推导的变换是线性的, 于是将变换直接写成如下形式

$$x'=\alpha x+\beta t \quad (3)$$

$$t'=yx+\delta t \quad (4)$$

并由此计算出 O' 在 S 系中的速度 $dx/dt=u$ 和 O 在 S' 系中的速度 $dx'/dt'=-u$, 其次是根据空间的

(9) 教师获奖佐证材料

| 序号 | 成果名称/获奖名称 | 级别 | 获奖等级 | 负责人/教师姓名 | 年份 |
|-----|--------------------|-----|------|-------------------|------|
| 1. | 中国政府友谊奖 | 国家级 | | A. A. Kudryavtsev | 2021 |
| 2. | 国家技术发明奖 | 国家级 | 一等奖 | 马晶 | 2014 |
| 3. | 国防技术发明奖 | 国家级 | 特等奖 | 马晶 | 2013 |
| 4. | 国防科技进步奖 | 国家级 | 一等奖 | 马晶 | 2018 |
| 5. | 教育部长江学者特聘教授 | 国家级 | | 赵纪军 | 2017 |
| 6. | 国家自然科学二等奖 | 国家级 | 二等奖 | 赵纪军 | 2010 |
| 7. | 国家技术发明奖 | 国家级 | 二等奖 | 马晶 | 2009 |
| 8. | 国防科技进步奖 | 国家级 | 二等奖 | 马晶 | 2010 |
| 9. | 国防科技进步奖 | 国家级 | 二等奖 | 马晶 | 2008 |
| 10. | 国防技术发明奖 | 国家级 | 二等奖 | 马晶 | 2008 |
| 11. | 国防科学技术奖 | 国家级 | 二等奖 | 马晶 | 2006 |
| 12. | 国防科技进步奖 | 国家级 | 三等奖 | 马晶 | 2002 |
| 13. | 2017年度国防科技工业十大创新人物 | 国家级 | | 马晶 | 2017 |

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|-----|---|-----|---------|-----|------|
| 14. | CCTV 科技盛典“2017 年度科技创新人物” | 国家级 | | 马晶 | 2017 |
| 15. | 全国“最美教师”荣誉称号 | 国家级 | | 马晶 | 2016 |
| 16. | 全国模范教师 | 国家级 | | 马晶 | 2014 |
| 17. | 国防科技创新团队奖 | 国家级 | | 马晶 | 2015 |
| 18. | 国家万人计划领军人才 | 国家级 | | 赵纪军 | 2017 |
| 19. | 国家千人青年项目 | 国家级 | | 高峻峰 | 2018 |
| 20. | 国家自然科学基金优秀青年基金 | 国家级 | | 周思 | 2022 |
| 21. | 全国高等学校教师自制实验设备创新大赛 | 国家级 | 一等奖 | 王健 | 2021 |
| 22. | 第一届“高等教育杯”全国高等学校物理基础课程青年教师讲课比赛 | 国家级 | 一等奖（冠军） | 孟庆鑫 | 2014 |
| 23. | 2018 年全国高校教师教学创新大赛——第五届全国高等学校教师自制实验教学仪器设备创新大赛 | 国家级 | 银奖 | 刘世刚 | 2018 |
| 24. | 第六届“高等教育杯”全国高等学校物理基础课程青年教师讲课比赛 | 国家级 | 二等奖 | 张伶俐 | 2021 |
| 25. | 全国高等学校教师自制实验设备创新大赛 | 国家级 | 三等奖 | 娄秀涛 | 2021 |
| 26. | 全国大学生创新创业实践联盟年会优秀论文 | 国家级 | 三等奖 | 高波 | 2018 |
| 27. | 国务院特殊津贴 | 国家级 | | 赵纪军 | 2016 |

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|-----|--------------------------------|------|-------|-----|------|
| 28. | 国务院特殊津贴 | 国家级 | | 马晶 | 2015 |
| 29. | 全国光学与光学工程博士生学术联赛 | 国家级 | 百强导师奖 | 田浩 | 2021 |
| 30. | 全国光学与光学工程博士生学术联赛 | 国家级 | 百强导师奖 | 丁卫强 | 2021 |
| 31. | 黑龙江省发明奖 | 省级 | 特等奖 | 马晶 | 2016 |
| 32. | 黑龙江省课程思政教学竞赛 | 省级 | 特等奖 | 任延宇 | 2021 |
| 33. | 中国高等学校十大科技进展 | 省部级 | | 谭立英 | 2017 |
| 34. | 中国高等学校十大科技进展 | 省部级 | | 马晶 | 2013 |
| 35. | 自然科学奖 | 省部级 | 二等奖 | 赵纪军 | 2017 |
| 36. | 辽宁省五一劳动奖章 | 省部级 | | 赵纪军 | 2021 |
| 37. | 第六届“高等教育杯”全国高等学校物理基础课程青年教师讲课比赛 | 东北地区 | 一等奖 | 张伶俐 | 2021 |
| 38. | 龙江学者支持计划-特聘教授 | 省级 | | 田浩 | 2019 |
| 39. | 黑龙江省五一劳动奖章 | 省级 | | 马晶 | 2017 |
| 40. | 黑龙江省青年教学名师 | | | 李均 | 2021 |
| 41. | “龙江楷模” | 省级 | | 谭立英 | 2018 |
| 42. | 先进技术应用转化大赛 | 省部级 | 金奖 | 马晶 | 2019 |

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|-----|--------------------------------|-----|------------|-----|------|
| 43. | 黑龙江省青年教师教学竞赛 | 省级 | 理科组 一等奖 | 任延宇 | 2020 |
| 44. | 黑龙江省高校微课教学竞赛 | 省级 | 一等奖 | 任延宇 | 2018 |
| 45. | 卓越大学联盟青年教师教学创新大赛 | 省级 | 一等奖 | 任延宇 | 2018 |
| 46. | 第三届“高等教育杯”全国高等学校物理基础课程青年教师讲课比赛 | 省级 | 一等奖 | 黄丽 | 2017 |
| 47. | 第二届“高等教育杯”全国高等学校物理基础课程青年教师讲课比赛 | 省级 | 一等奖 | 韩权 | 2016 |
| 48. | 黑龙江省高等学校课程思政示范课程 | 省级 | | 任延宇 | 2020 |
| 49. | 黑龙江省高等学校物理基础课程青年教师讲课比赛 | 省级 | 二等奖 | 李均 | 2018 |
| 50. | 教育部自然科学奖 | 省部级 | 二等奖 | 赵纪军 | 2016 |
| 51. | 河北省自然科学奖 | 省部级 | 二等奖 | 赵纪军 | 2018 |
| 52. | 辽宁省自然科学奖 | 省部级 | 二等奖 | 赵纪军 | 2019 |
| 53. | 山西省自然科学奖 | 省部级 | 二等奖 | 赵纪军 | 2021 |
| 54. | 哈尔滨工业大学第二届课程思政教学竞赛 | 校级 | 一等奖 | 周可雅 | 2021 |
| 55. | 哈尔滨工业大学优秀教育教学管理奖 | 校级 | | 张宇 | 2021 |
| 56. | 哈尔滨工业大学 2020 年教学贡献奖励 | 校级 | 一等奖 | 时红艳 | 2020 |
| 57. | 基于分类人才培养的物理分层次教学模式研究 | 校级 | 一等奖 | 孟庆鑫 | 2019 |

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|-----|----------------------|----|----------|-----|------|
| 58. | 固体光学课程研究型教学理论与实践 | 校级 | 二等奖 | 田浩 | 2017 |
| 59. | 首届哈尔滨工业大学“育人新星”青年导师 | 校级 | | 田浩 | 2021 |
| 60. | 哈尔滨工业大学 2021 年教学贡献奖励 | 校级 | 教学优秀奖一等奖 | 李均 | 2021 |
| 61. | 哈尔滨工业大学首届课程思政教学竞赛 | 校级 | 特等奖 | 任延宇 | 2019 |
| 62. | 第七届校青年教师教学基本功竞赛 | 校级 | 一等奖 | 李均 | 2018 |
| 63. | 在线教学“我最喜爱的课程”奖 | 校级 | | 李均 | 2020 |

(1) 中国政府友谊奖获奖报道

哈尔滨工业大学 物理学院

首页 学院概况 师资队伍 招生与培养 科研与教学 国际合作 科学研究 合作交流 校友工作 人才招聘

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哈工大物理学院俄罗斯专家荣获“中国政府友谊奖”

2021年9月30日，中国政府友谊奖颁奖仪式在北京举行，国务院总理李克强亲自主持大会，宣读2020和2021年度中国政府友谊奖获奖外国专家，李克强对获奖的外国专家表示了热烈祝贺，充分肯定了他们为中国改革开放和发展进步、中外友好交往所作作出的贡献，并通过他们向所有在华工作的外国专家和家属，以及关心支持中国发展的国际友人表示亲切的问候。

哈工大物理学院俄罗斯专家安纳托利教授及其在中俄科技合作、中俄联合办学等方面作出的贡献，荣获2021年度“中国政府友谊奖”。

安纳托利教授是国际知名的气体放电及等离子体应用专家，经我校推荐，袁忠林、袁承勋两国领导长期合作，他积极投身于中国的基础科学研究和基础学科人才培养，尤其是在等离子体基础研究、基础学科人才联合培养、中俄联合办学等方面做出突出贡献。

中国政府友谊奖是中国政府为表彰在中国现代化建设和改革开放事业中作出突出贡献的外国专家而设立的崇高奖项，安纳托利教授荣获国家外国专家局于1991年正式设立，并在每年国庆节前夕举行“友谊奖”的颁奖仪式，这是中国政府对贡献突出的外国专家给予的最高荣誉。

物理学院积极开展对外合作，先后与圣彼得堡国立大学、莫斯科国立大学、俄罗斯科学院南俄方单位开展了深入的学术交流和合作，先后组织了7届中俄学术研讨会，4届等离子体物理暑期学校和1届中俄物理学术研讨会，此外还组织不定期的学术互访、5年来，超过100人次来专家应邀来我校进行学术交流，举办学术讲座上百次，中俄双方合作取得了一系列重要成果，获批国家自然科学基金中俄合作交流项目、中青年人才项目等6项国际合作项目，并获批2000年以来创新型人才国际合作培养项目，全职引进俄罗斯专家4名，短期合作教授20余名，通过学术交流合作及人才交流，形成了稳定的合作研究方向，促进了物理学科的发展，推动了我校大学科工程的建设，为“中俄联合校园”的建设和发展奠定了基础。

(2) 国家技术发明一等奖获奖证明



(3) 国防技术发明特等奖获奖证明



(4) 国防科技进步一等奖获奖证明



(5) 教育部长江学者特聘教授证明



(6) 国家自然科学二等奖获奖证明



(7) 国家技术发明二等奖获奖证明



(8) 国防科技进步二等奖获奖证明



(9) 国防科技进步二等奖获奖证明



(10) 国防技术发明奖二等奖获奖证明



(11) 国防科学技术奖二等奖获奖证明



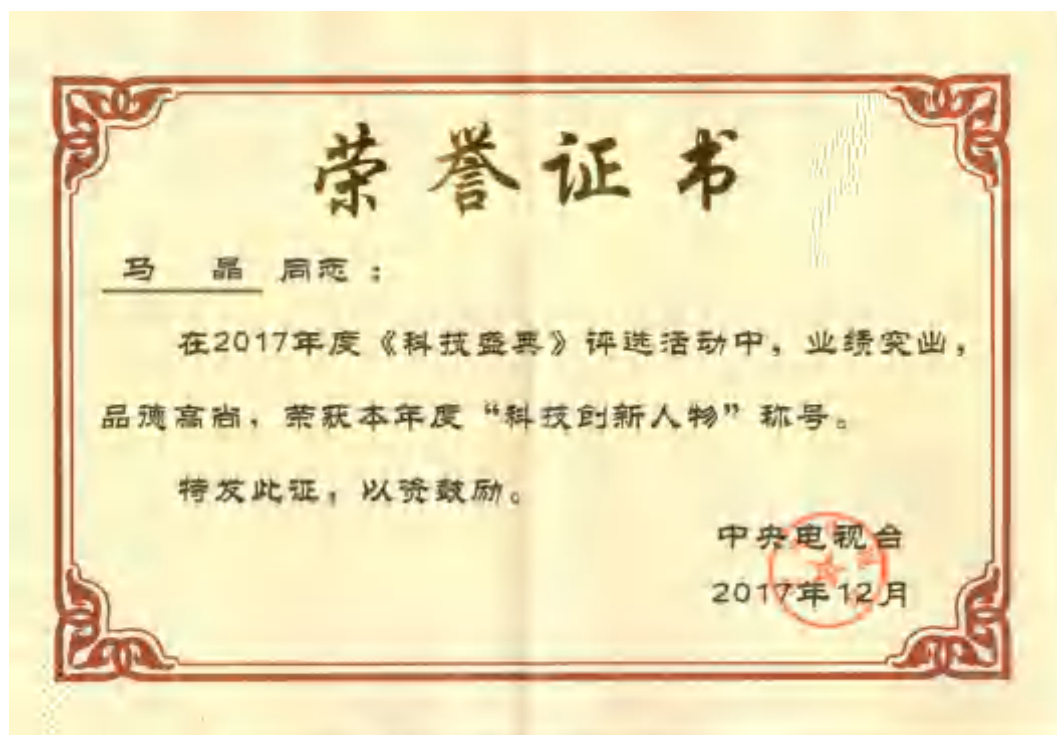
(12) 国防科技进步三等奖获奖证明



(13) 2017 年度国防科技工业十大创新人物



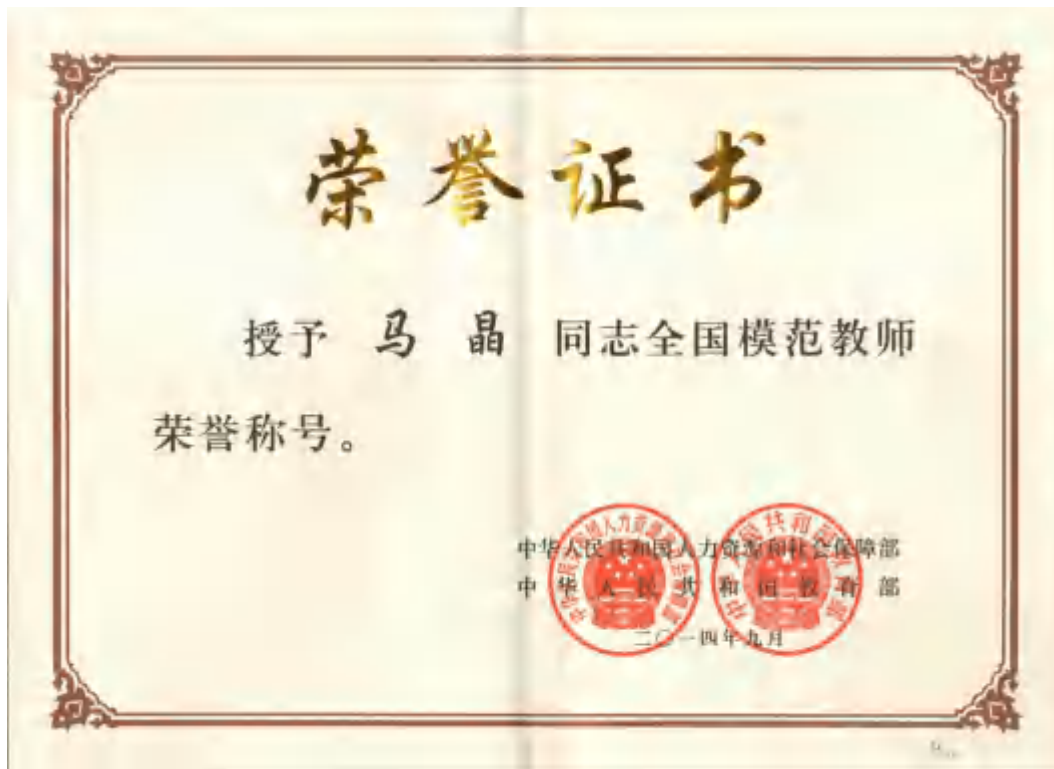
(14) CCTV 科技盛典 “2017 年度科技创新人物”



(15) 全国最美教师荣誉称号



(16) 全国模范教师



(17) 国防科技创新团队奖获奖证明



(18) 国家万人计划领军人才证明



(19) 国家千人青年项目证明

中共大连理工大学委员会人才办公室



关于大连理工大学第十四批国家“千人计划”入选人员名单的证明

根据中组部办公厅《关于印发第十四批国家“千人计划”入选人员名单的通知》（组厅字〔2018〕3号）精神和要求，大连理工大学 Teel Andrew Richard 入选外专短期项目，方蔚瑞、高峻峰、邵堃入选青年项目。

特此证明。

中共大连理工大学委员会人才办公室

2018年9月21日

人才办公室

(20) 国家自然科学基金优秀青年基金证明

国家自然科学基金资助项目批准通知

(包干制项目)

周思 先生/女士:

根据《国家自然科学基金条例》、相关项目管理办法规定和专家评审意见,国家自然科学基金委员会(以下简称自然科学基金委)决定资助您申请的项目。项目批准号:12222403,项目名称:固属物理及其交叉学科研究,资助经费:200.00万元,项目起止年月:2023年01月至2025年12月,有关项目的评审意见及修改意见附后。

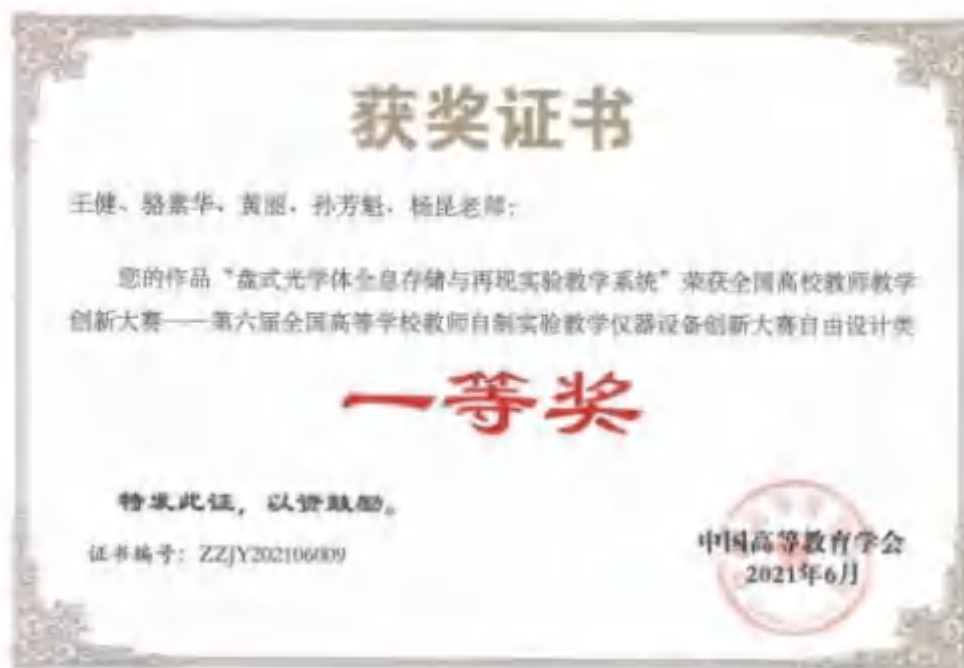
请您尽快登录科学基金网络信息系统(<https://isisn.nsf.gov.cn>),认真阅读《国家自然科学基金资助项目计划书填报说明》并按要求填写《国家自然科学基金资助项目计划书》(以下简称计划书)。对于有修改意见的项目,请您按修改意见及时调整计划书相关内容;如您对修改意见有异议,请在电子版计划书报送截止日期前向相关科学处提出。

请您将电子版计划书通过科学基金网络信息系统(<https://isisn.nsf.gov.cn>)提交,由依托单位审核后提交至自然科学基金委。自然科学基金委审核未通过的,将退回的电子版计划书修改后再行提交;审核通过的,打印纸质版计划书(一式两份,双面打印)并在项目负责人承诺栏签字,由依托单位在承诺栏加盖依托单位公章,且将申请书纸质签字盖章页订在其中一份计划书之后,一并送至自然科学基金委项目材料接收工作组。纸质版计划书应当保证与审核通过的电子版计划书内容一致。自然科学基金委将对申请书纸质签字盖章页进行审核,对存在问题的,允许依托单位进行一次修改或补齐。

向自然科学基金委提交电子版计划书,报送纸质版计划书并补交申请书纸质签字盖章页截止时间节点如下:

1. 2022年8月26日16点:提交电子版计划书的截止时间(视为计划书正式提交时间);
2. 2022年9月2日16点:提交修改后电子版计划书的截止时间;
3. 2022年9月9日16点:报送纸质版计划书(其中一份包含申请书纸质签字盖章页)的截止时间;

(21) 全国高等学校教师自制实验设备创新大赛一等奖获奖证明



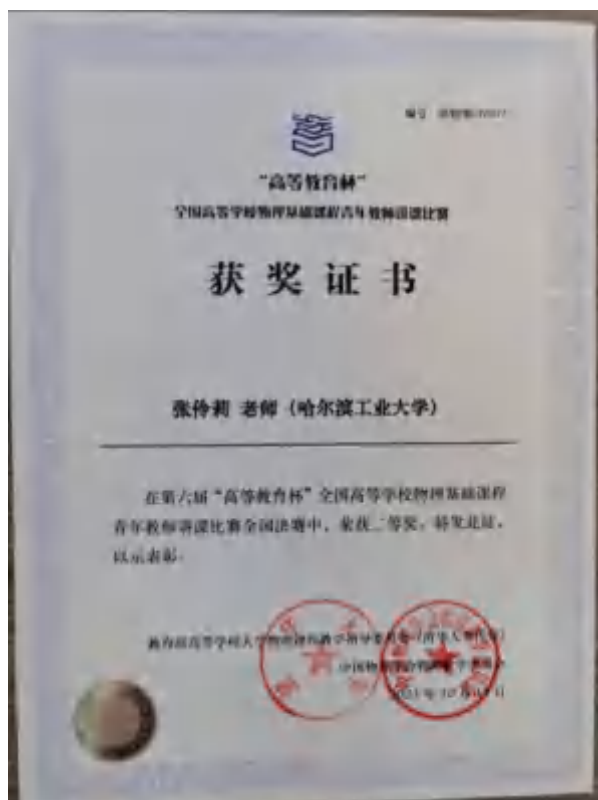
(22) 第一届“高等教育杯”全国高等学校物理基础课程青年教师讲课比赛国家级一等奖（冠军）证明



(23) 2018 年全国高校教师教学创新大赛——第五届全国高等学校教师自制实验教学仪器设备创新大赛证明



(24) 第六届“高等教育杯”全国高等学校物理基础课程青年教师讲课比赛国家级二等奖证明



(25) 全国高等学校教师自制实验设备创新大赛三等奖获奖证明



(26) 全国大学生创新创业实践联盟年会优秀论文获奖证明



(27) 国务院特殊津贴证明



(28) 国务院特殊津贴证明



(29) 全国光学与光学工程博士生学术联赛百强导师奖证明



(30) 全国光学与光学工程博士生学术联赛百强导师奖证明



(31) 黑龙江省科学技术奖二等奖获奖证书



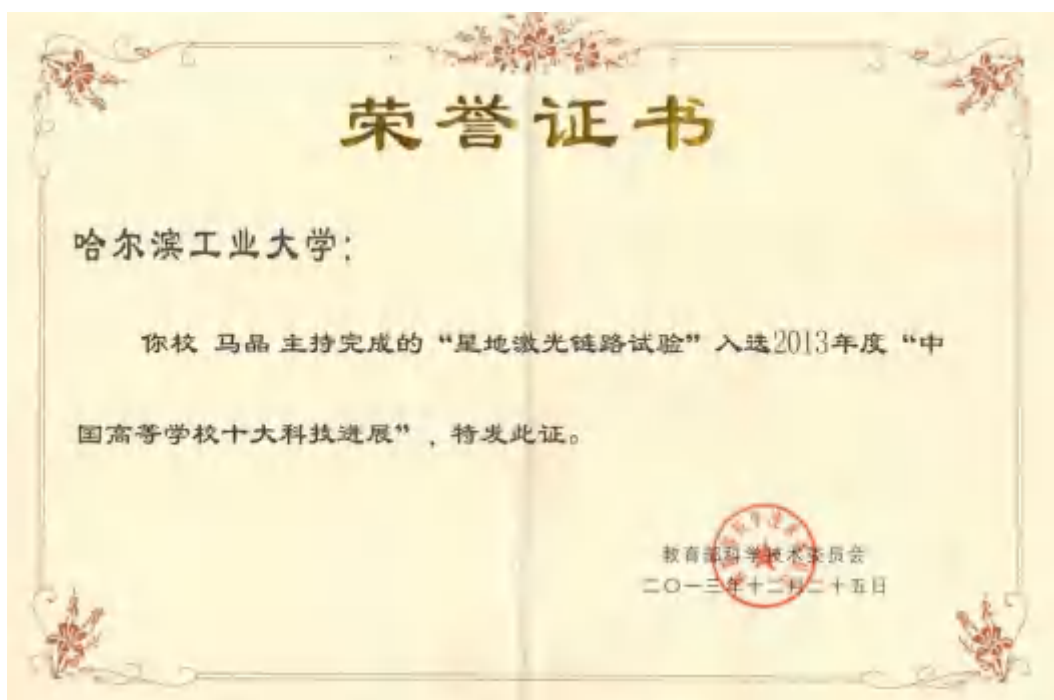
(32) 黑龙江省课程思政教学竞赛特等奖证明



(33) 2017 年度中国高等学校十大科技进展获奖证明



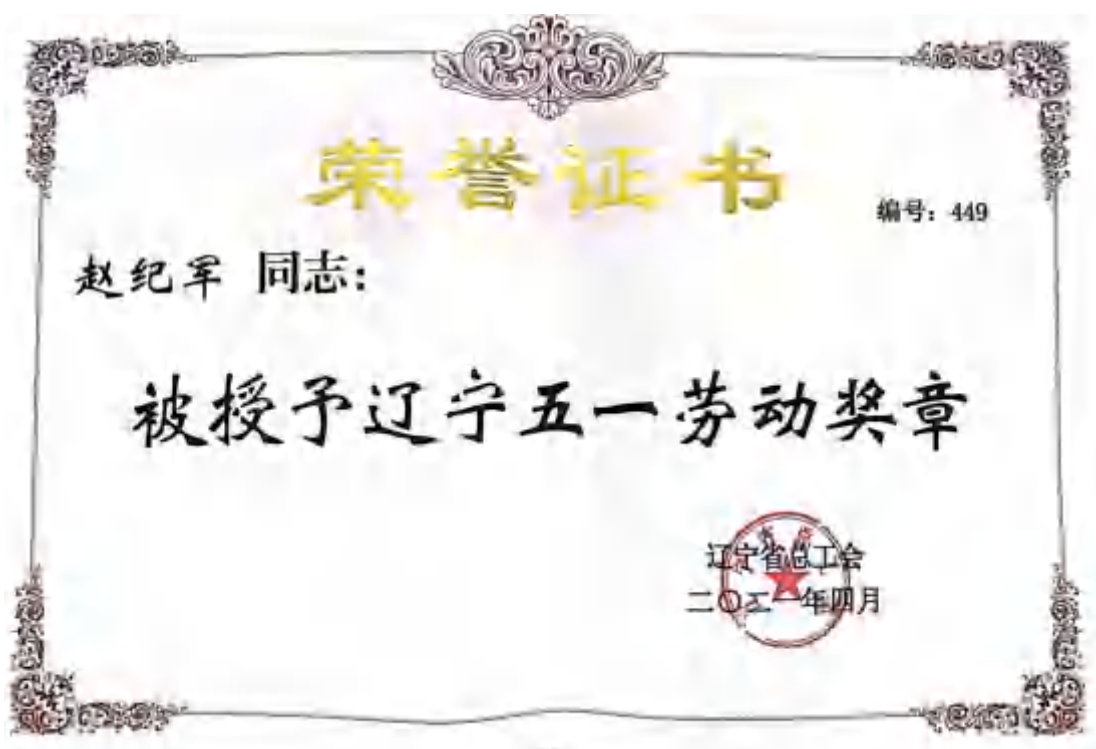
(34) 2013 年度中国高等学校十大科技进展获奖证明



(35) 自然科学奖二等奖获奖证明



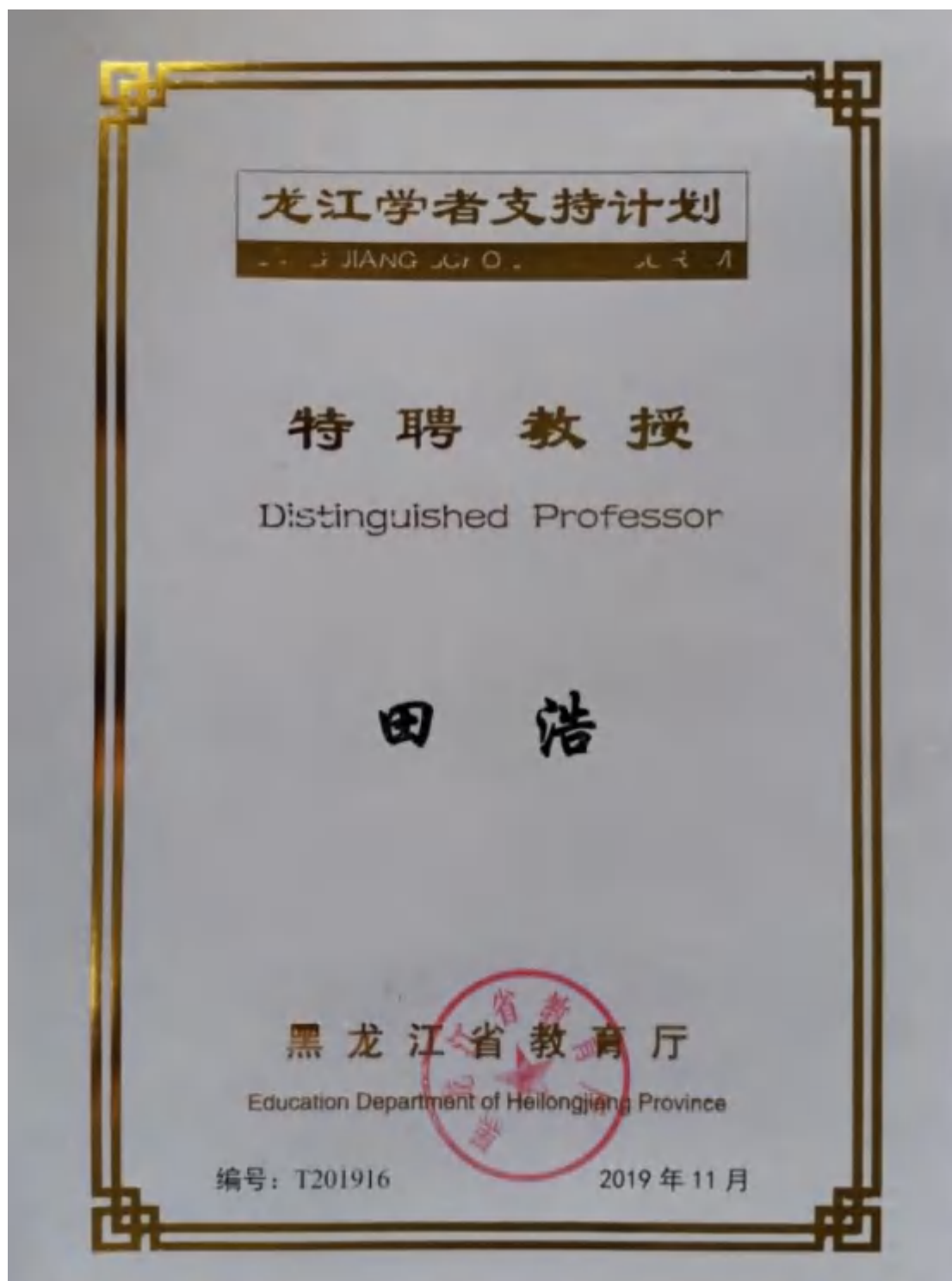
(36) 辽宁省五一劳动奖章证明



(37) 第六届“高等教育杯”全国高等学校物理基础课程青年教师讲课比赛东北赛区一等奖证明



(38) 龙江学者支持计划-特聘教授证明



(39) 黑龙江省五一劳动奖章



(40) 黑龙江省青年教学名师证明

2021年黑龙江省教学名师候选人名单

| 序号 | 候选人姓名 | 工作单位 | 地区 |
|---------------|-------|---------|------|
| 普通高等學校 (113名) | | | |
| 1 | 钱洋 | 哈尔滨理工大学 | 哈尔滨市 |
| 2 | 孙强 | 哈尔滨理工大学 | 哈尔滨市 |
| 3 | 刘立伟 | 哈尔滨理工大学 | 哈尔滨市 |
| 4 | 孙强 | 哈尔滨理工大学 | 哈尔滨市 |
| 5 | 张道新 | 哈尔滨理工大学 | 哈尔滨市 |
| 6 | 曹强松 | 哈尔滨理工大学 | 哈尔滨市 |
| 7 | 马树敏 | 哈尔滨理工大学 | 哈尔滨市 |
| 8 | 郭金梅 | 哈尔滨理工大学 | 哈尔滨市 |
| 9 | 张景华 | 哈尔滨理工大学 | 哈尔滨市 |
| 10 | 王磊 | 哈尔滨理工大学 | 哈尔滨市 |
| 11 | 孙强 | 哈尔滨理工大学 | 哈尔滨市 |
| 12 | 魏华 | 哈尔滨理工大学 | 哈尔滨市 |
| 13 | 李平山 | 哈尔滨理工大学 | 哈尔滨市 |
| 14 | 孙景波 | 哈尔滨理工大学 | 哈尔滨市 |

(41) “龙江楷模”荣誉称号



(42) 2019年先进技术应用转化大赛金奖



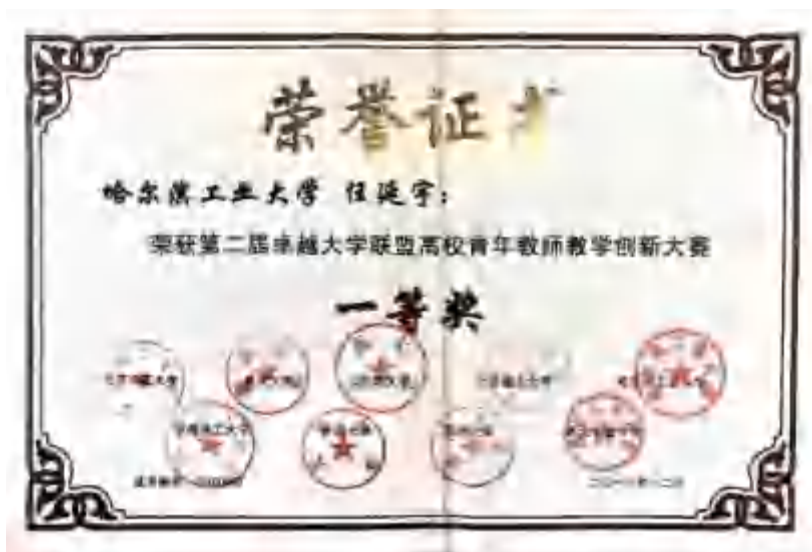
(43) 黑龙江省青年教师教学理科组一等奖获奖证明



(44) 黑龙江省高校微课教学竞赛一等奖证明



(45) 卓越大学联盟青年教师教学创新大赛一等奖证明



(46) 第三届“高等教育杯”全国高等学校物理基础课程青年教师讲课比赛省级一等奖证明



(47) 第二届“高等教育杯”全国高等学校物理基础课程青年教师讲课比赛省级一等奖证明



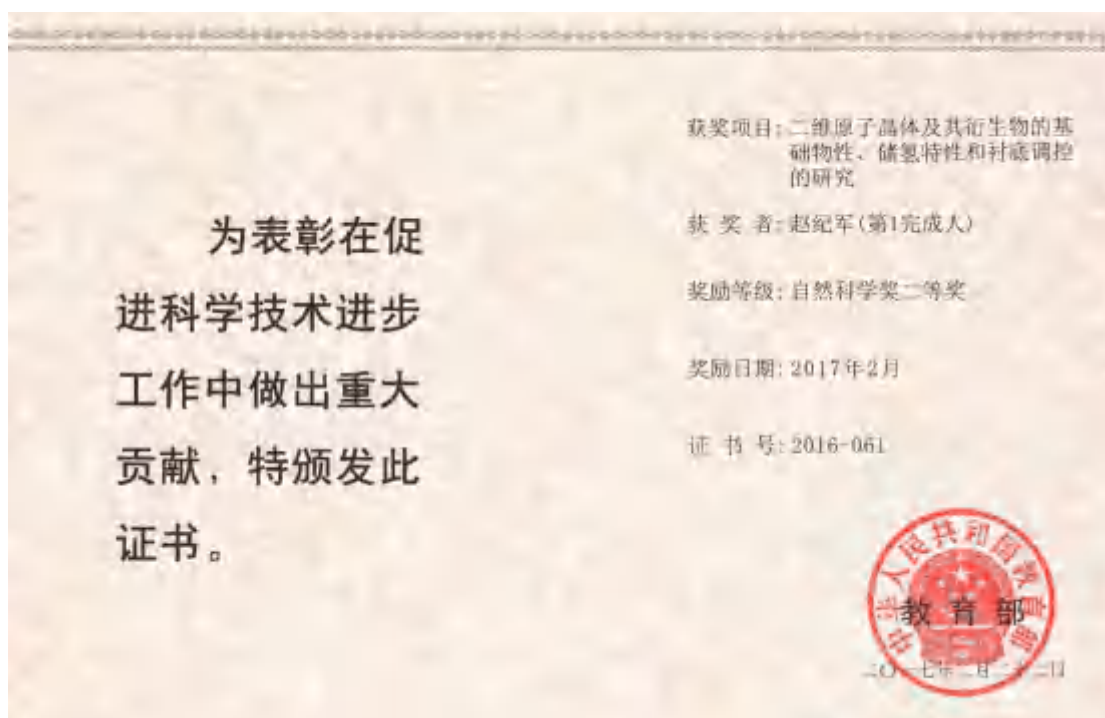
(48) 黑龙江省高等学校课程思政示范课程获奖证明



(49) 黑龙江省高等学校物理基础课程青年教师讲课比赛二等奖证明



(50) 教育部自然科学奖二等奖获奖证明



(51) 河北省自然科学奖二等奖获奖证明



(52) 辽宁省自然科学奖二等奖获奖证明



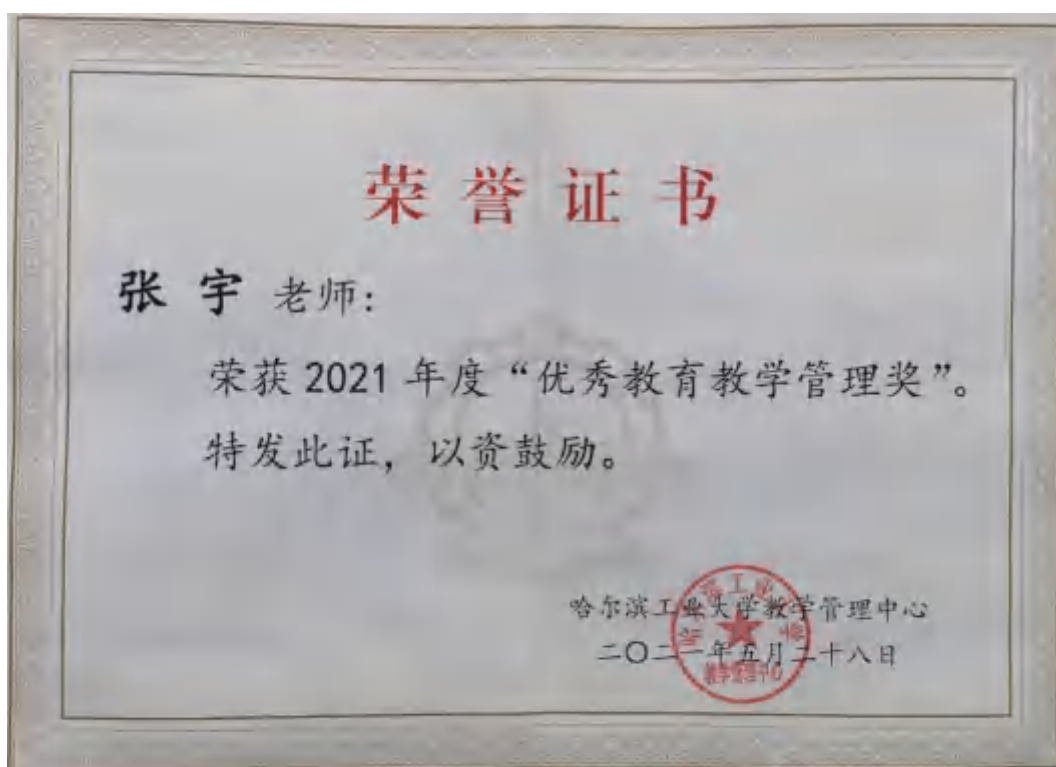
(53) 山西省自然科学奖二等奖获奖证明



(54) 哈尔滨工业大学第二届课程思政教学竞赛一等奖获奖证明



(55) 哈尔滨工业大学优秀教育教学管理奖获奖证明



(56) 哈尔滨工业大学 2020 年教学贡献奖励一等奖获奖证明



(57) 基于分类人才培养的物理分层次教学模式研究校级一等奖证明

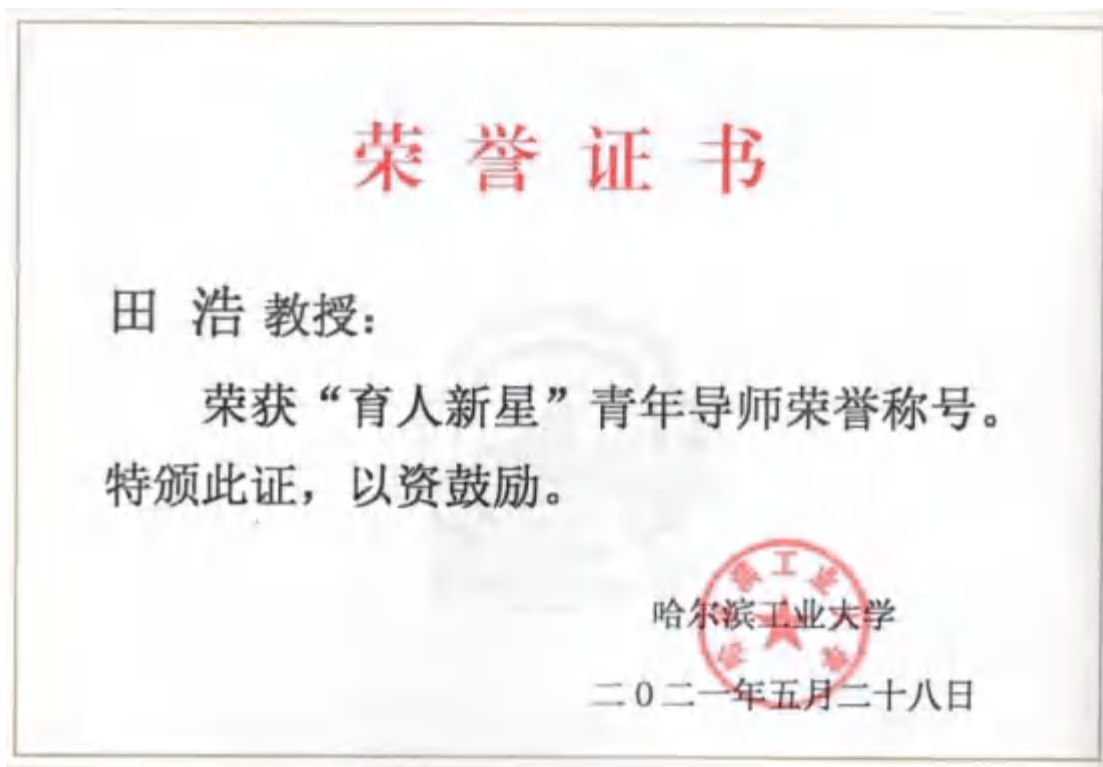


(58) 固体光学课程研究型教学理论与实践或研究生教育成果
 奖证明

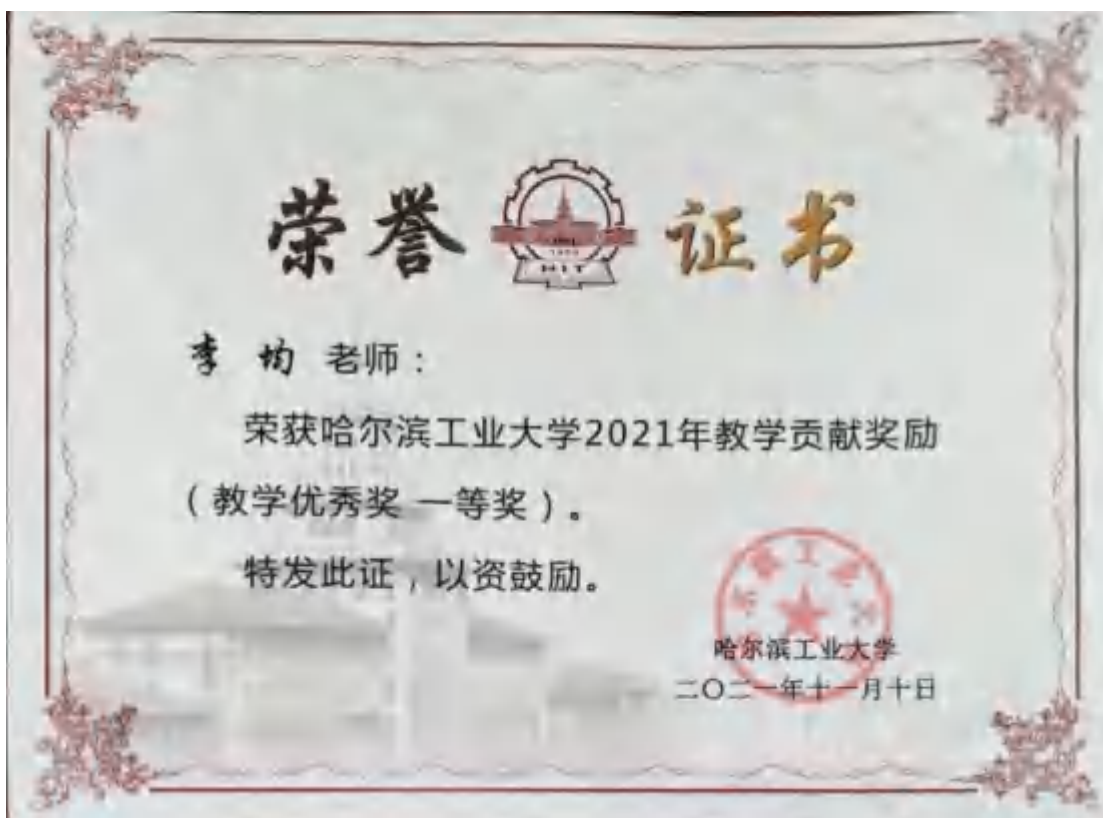
2017 年度研究生教育成果奖评审结果

| 序号 | 成果名称 | 获（第） | 评审结果 |
|----|-----------------------------------|---------|------|
| 1 | 全新发展与改进研究生教育实践的方法与实践——以合肥工业大学为例 | 研究生院 | 一等奖 |
| 2 | 基于工学在通用工业创新工场——航天类高层次人才培养模式的探索与实践 | 航天学院 | 一等奖 |
| 3 | “融合经济，四类培养”——工科化学类本研贯通式培养人才培养模式研究 | 化工与化学学院 | 一等奖 |
| 4 | 多元文化背景下研究生课程教学模式创新，理论与实践 | 外国语学院 | 一等奖 |
| 5 | 研究生国际化高水平课程建设的探索与实践 | 研究生院 | 一等奖 |
| 6 | 基于 COE 的“EDA/SCPC 应用实践”课程改革与实践 | 电子学院 | 一等奖 |
| 7 | “互联网+”下交通运输工程学科的新人才培养模式的实践 | 交通学院 | 一等奖 |
| 8 | 教师队伍建设与研究生培养模式改革与实践 | 化学学院 | 一等奖 |
| 9 | 增量式教育模式在技术大学研究生培养模式研究及借鉴 | 电气学院 | 一等奖 |
| 10 | 固体光学课程研究型教学理论与实践 | 物理学院 | 一等奖 |
| 11 | 基于混合教学的研究型课程建设与实践——以《应用物理过程》为例 | 物理学院 | 一等奖 |
| 12 | 专业学位硕士实践型人才培养模式——以风景园林学科为例 | 建筑学院 | 二等奖 |
| 13 | 《最优化方法》研究生精品课程建设 | 数学学院 | 二等奖 |
| 14 | 以学科内容为基础培养专业学位研究生的国际化沟通能力之研究与实践 | 外国语学院 | 二等奖 |
| 15 | 课程案例库建设实践，提高第七课文生实践能力及实践能力 | 外国语学院 | 二等奖 |
| 16 | 工商管理专业研究生教学模式改革研究 | 管理学院 | 二等奖 |
| 17 | 理工类基础日语基础教材及拓展翻译材料开发 | 外国语学院 | 二等奖 |
| 18 | 航空宇航科学与技术专业硕士学位培养模式学科基础课程建设研究 | 航空学院 | 二等奖 |
| 19 | 清华两校联合培养模式与建筑类学术研究生教育实践 | 建筑学院 | 二等奖 |
| 20 | 研究生新工科课程与培养模式创新与实践 | 航天学院 | 二等奖 |

(59) 首届哈尔滨工业大学“育人新星”青年导师证明



(60) 哈尔滨工业大学 2021 年教学贡献奖获奖证明



(61) 哈尔滨工业大学首届课程思政教学竞赛校级特等奖获奖证明



(62) 第七届校青年教师教学基本功竞赛一等奖证明



(63) 在线教学“我最喜爱的课程”奖证明

